# On Cayley conditions for billiards inside ellipsoids ${ }^{\Delta \pi}$ 

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#### Abstract

All the segments (or their continuations) of a billiard trajectory inside an ellipsoid of $\mathbb{R}^{n}$ are tangent to $n-1$ quadrics of the pencil of confocal quadrics determined by the ellipsoid. The quadrics associated to periodic billiard trajectories verify certain algebraic conditions. Cayley found them in the planar case. Dragović and Radnović generalized them to any dimension. We rewrite the original matrix formulation of these generalized Cayley conditions as a simpler polynomial one. We find several remarkable algebraic relations between caustic parameters and ellipsoidal parameters that give rise to nonsingular periodic trajectories.


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## 1. Introduction

One of the best known discrete integrable system is the billiard inside ellipsoids. All the segments (or their continuations) of a billiard trajectory inside an ellipsoid of $\mathbb{R}^{n}$ are tangent to $n-1$ quadrics of the pencil of confocal quadrics determined by the ellipsoid [1-3]. This situation is fairly exceptional. Quadrics are the only smooth hypersurfaces of $\mathbb{R}^{n}, n \geq 3$, that have caustics [4,5]. A caustic is a smooth hypersurface with the property that a billiard trajectory, once tangent to it, stays tangent after every reflection. Caustics are a geometric manifestation of the integrability of billiards inside ellipsoids.

Periodic trajectories are the most distinctive trajectories, so their study is the first task. There exist two remarkable results about periodic billiard trajectories inside ellipsoids: the generalized Poncelet theorem and the generalized Cayley conditions.

A classical geometric theorem of Poncelet [6, 7] implies that if a billiard trajectory inside an ellipse is periodic, then all the trajectories sharing its caustic are also periodic. Its generalization to the spatial case was proved by Darboux [8]. The extension of this result to arbitrary dimensions can be found in [9-12]. The generalized Poncelet theorem can be stated as follows. If a billiard trajectory inside an ellipsoid is closed after $m_{0}$ bounces and has length $L_{0}$, then all trajectories sharing the same caustics are also closed after $m_{0}$ bounces and have length $L_{0}$. Thus, a natural

[^0]question arises. What caustics do give rise to periodic trajectories? The planar case was solved by Cayley [13, 14] in the XIX century. Dragović and Radnović [15, 16] found some generalized Cayley conditions for billiards inside ellipsoids fifteen years ago. They have also stated similar conditions in other billiard frameworks; see [17-21].

For simplicity, let us focus on the spatial case. Let $Q: x^{2} / a+y^{2} / b+z^{2} / c=1$ be the triaxial ellipsoid with ellipsoidal parameters $0<c<b<a$. Any billiard trajectory inside $Q$ has as caustics two elements of the family of confocal quadrics

$$
Q_{\lambda}=\left\{(x, y, z) \in \mathbb{R}^{3}: \frac{x^{2}}{a-\lambda}+\frac{y^{2}}{b-\lambda}+\frac{z^{2}}{c-\lambda}=1\right\} .
$$

We restrict our attention to nonsingular trajectories. That is, trajectories with two different caustics which are ellipsoids: $0<\lambda<c$, hyperboloids of one sheet: $c<\lambda<b$, or hyperboloids of two sheets: $b<\lambda<a$. The singular values $\lambda \in\{a, b, c\}$ are discarded. There exist some restrictions on the caustics $Q_{\lambda_{1}}$ and $Q_{\lambda_{2}}$. It is known that $\mathrm{EH} 1, \mathrm{H} 1 \mathrm{H} 1, \mathrm{EH} 2$, and H 1 H 2 are the only feasible caustic types; see [22,23]. The meaning of these notations is evident.

The generalized Cayley condition in this context can be expressed as follows. The billiard trajectories inside the triaxial ellipsoid $Q$ sharing the caustics $Q_{\lambda_{1}}$ and $Q_{\lambda_{2}}$ are periodic with elliptic period $m \geq 3$ if and only if

$$
\operatorname{rank}\left(\begin{array}{ccc}
f_{m+1} & \cdots & f_{4} \\
\vdots & & \vdots \\
f_{2 m-1} & \cdots & f_{m+2}
\end{array}\right)<m-2,
$$

where $f(t)=\sum_{l \geq 0} f_{l} t^{l}:=\sqrt{(1-t / a)(1-t / b)(1-t / c)\left(1-t / \lambda_{1}\right)\left(1-t / \lambda_{2}\right)}$. We claim that most results related to generalized Cayley conditions become simpler when expressed in terms of the inverse quantities $1 / a, 1 / b, 1 / c, 1 / \lambda_{1}$, and $1 / \lambda_{2}$. Proposition 12 is a paradigmatic example.

The elliptic period is defined in Section 2. Roughly speaking, the difference between the period $m_{0}$ and the elliptic period $m$ of the periodic billiard trajectories sharing two given caustics is that all those trajectories close in Cartesian (respectively, elliptic) coordinates after exactly $m_{0}$ (respectively, $m$ ) bounces. We will see that either $m=m_{0} / 2$ or $m=m_{0}$.

The previous matrix formulation is very nice from a theoretical point of view, but it has strong limitations from a computational point of view. We will see in Section 3 that it can be written as a system of two homogeneous symmetric polynomial equations with rational coefficients of degrees $m^{2}-2$ and $m^{2}-1$ in the variables $1 / a, 1 / b, 1 / c, 1 / \lambda_{1}$, and $1 / \lambda_{2}$. Thus, both degrees grow quadratically with the elliptic period $m$, which turns this approach into a tough challenge. In particular, to our knowledge, the caustic parameters $\lambda_{1}$ and $\lambda_{2}$ have never been explicitly expressed in terms of the ellipsoidal parameters $a, b$, and $c$ for any $m \geq 3$.

We will rewrite this matrix formulation as a computationally more appealing one which gives rise to (non-symmetric) homogeneous polynomial equations whose degrees are smaller than the elliptic period $m$. We will find the following remarkable algebraic relations between caustic and ellipsoidal parameters using the new formulation. The billiard trajectories inside the ellipsoid $Q$ sharing the caustics $Q_{\lambda_{1}}$ and $Q_{\lambda_{2}}$ are periodic with:

- Elliptic period $m=3$ if the roots of $t^{3}-(t-c)(t-b)(t-a)$ are the caustic parameters;
- Elliptic period $m=4$ if there exists $d \in \mathbb{R}$ such that the roots of $t^{4}-(t-d)^{2}\left(t-\lambda_{1}\right)\left(t-\lambda_{2}\right)$ are the ellipsoidal parameters; and
- Elliptic period $m=5$ if the roots of $t^{5}-(t-c)(t-b)(t-a)\left(t-\lambda_{1}\right)\left(t-\lambda_{2}\right)$ are double ones.

Let us compare both formulations in the third case. On the one hand, the two homogeneous symmetric polynomial equations obtained from the matrix formulation have degrees 23 and 24 in the variables $1 / a, 1 / b, 1 / c, 1 / \lambda_{1}$, and $1 / \lambda_{2}$. On the other hand, $t^{5}-(t-c)(t-b)(t-a)\left(t-\lambda_{1}\right)\left(t-\lambda_{2}\right)$ is a polynomial of degree four in a single variable. Clearly, the polynomial formulation leads to a much simpler problem. Nevertheless, it should be stressed that the matrix formulation determines all periodic billiard trajectories with elliptic period $m=5$. On the contrary, we find just a subset of such trajectories using the polynomial $t^{5}-(t-c)(t-b)(t-a)\left(t-\lambda_{1}\right)\left(t-\lambda_{2}\right)$.

Another natural question about periodic billiard trajectories is the following one. Which are the triaxial ellipsoids of $\mathbb{R}^{3}$ that display periodic billiard trajectories with a fixed caustic type and a fixed (elliptic) period? A numerical approach to that question was considered in [24], where the authors computed several bifurcations in the space of ellipsoidal parameters. We will find the algebraic relations that define the bifurcations associated to small elliptic periods. For instance, we will see that there exist periodic billiard trajectories with elliptic period $m=3$ and caustic type EH1 if and only if $c<a b /(a+b+\sqrt{a b})$.

For brevity, we will not depict billiard trajectories inside triaxial ellipsoids of $\mathbb{R}^{3}$. The reader interested in 3D graphical visualizations is referred to [25], where several periodic billiard trajectories with small periods are displayed from different perspectives.

We complete this introduction with a note on the organization of the article. In Section 2 we review briefly some well-known results about billiards inside ellipsoids, recalling the matrix formulation of the generalized Cayley conditions obtained by Dragović and Radnović. We also introduce the concept of elliptic period. The practical limitations of the matrix formulation are exposed in Section 3. Next, we present the polynomial formulation in Section 4. In Section 5 we carry out a detailed analysis for minimal elliptic periods, whereas the study of more general elliptic periods is postponed to Section 6. The previous results are adapted to billiards inside ellipses of $\mathbb{R}^{2}$ and inside triaxial ellipsoids of $\mathbb{R}^{3}$ in sections 7 and 8 , respectively.

## 2. Preliminaries

In this section we recall several classical results and their modern generalizations about billiards inside ellipsoids that go back to Jacobi, Chasles, Poncelet, Darboux, and Cayley.

We consider the billiard dynamics inside the ellipsoid

$$
\begin{equation*}
Q=\left\{x=\left(x_{1}, \ldots, x_{n}\right) \in \mathbb{R}^{n}: \sum_{i=1}^{n} \frac{x_{i}^{2}}{a_{i}}=1\right\}, \quad 0<a_{1}<\cdots<a_{n} . \tag{1}
\end{equation*}
$$

The degenerate cases in which the ellipsoid has some symmetry of revolution are not considered here. This ellipsoid is an element of the family of confocal quadrics

$$
Q_{\lambda}=\left\{x=\left(x_{1}, \ldots, x_{n}\right) \in \mathbb{R}^{n}: \sum_{i=1}^{n} \frac{x_{i}^{2}}{a_{i}-\lambda}=1\right\}, \quad \lambda \in \mathbb{R}
$$

We note that $Q_{\lambda}=\emptyset$ for $\lambda>a_{n}$. Thus, there are exactly $n$ different geometric types of nonsingular quadrics in this family, which correspond to the cases

$$
\begin{gathered}
\lambda \in\left(-\infty, a_{1}\right), \quad \lambda \in\left(a_{1}, a_{2}\right), \quad \ldots, \quad \lambda \in\left(a_{n-1}, a_{n}\right) . \\
3
\end{gathered}
$$

For instance, the confocal quadric $Q_{\lambda}$ is an ellipsoid if and only if $\lambda \in\left(-\infty, a_{1}\right)$. On the other hand, the meaning of $Q_{\lambda}$ in the singular cases $\lambda \in\left\{a_{1}, \ldots, a_{n}\right\}$ is

$$
Q_{a_{j}}=H_{j}=\left\{x=\left(x_{1}, \ldots, x_{n}\right) \in \mathbb{R}^{n}: x_{j}=0\right\} .
$$

The following theorems of Jacobi and Chasles can be found in [1-3].
Theorem 1 (Jacobi). Any generic point $x \in \mathbb{R}^{n}$ belongs to exactly $n$ distinct nonsingular quadrics $Q_{\mu_{0}}, \ldots, Q_{\mu_{n-1}}$ such that $\mu_{0}<a_{1}<\mu_{1}<a_{2}<\cdots<a_{n-1}<\mu_{n-1}<a_{n}$.

We denote by $\mu=\left(\mu_{0}, \ldots, \mu_{n-1}\right) \in \mathbb{R}^{n}$ the Jacobi elliptic coordinates of the point $x=$ $\left(x_{1}, \ldots, x_{n}\right)$. Cartesian and elliptic coordinates are linked by relations

$$
x_{j}^{2}=\frac{\prod_{i=0}^{n-1}\left(a_{j}-\mu_{i}\right)}{\prod_{i \neq j}\left(a_{j}-a_{i}\right)}, \quad j=1, \ldots, n .
$$

Hence, a point has the same elliptic coordinates that its orthogonal reflections with respect to the coordinate subspaces of $\mathbb{R}^{n}$. A point is generic, in the sense of Theorem 1, if and only if it is outside all coordinate hyperplanes. When a point tends to the coordinate hyperplane $H_{j}$, some of its elliptic coordinates tends to $a_{j}$.
Theorem 2 (Chasles). Any line in $\mathbb{R}^{n}$ is tangent to exactly $n-1$ confocal quadrics $Q_{\lambda_{1}}, \ldots, Q_{\lambda_{n-1}}$.
It is known that if two lines obey the reflection law at a point $x \in Q$, then both lines are tangent to the same confocal quadrics. Thus, all lines of a billiard trajectory inside the ellipsoid $Q$ are tangent to exactly $n-1$ confocal quadrics $Q_{\lambda_{1}}, \ldots, Q_{\lambda_{n-1}}$, which are called caustics of the trajectory, whereas $\lambda_{1}, \ldots, \lambda_{n-1}$ are the caustic parameters of the trajectory. We will say that a billiard trajectory inside $Q$ is nonsingular when it has $n-1$ distinct nonsingular caustics. We mostly deal with nonsingular billiard trajectories in this paper.

The caustic parameters cannot take arbitrary values. For instance, a line cannot be tangent to two different confocal ellipsoids, and all caustics parameters must be positive. The following complete characterization was given in [22,23].

Proposition 3. Let $\lambda_{1}<\cdots<\lambda_{n-1}$ be some real numbers such that

$$
\left\{a_{1}, \ldots, a_{n}\right\} \cap\left\{\lambda_{1}, \ldots, \lambda_{n-1}\right\}=\emptyset .
$$

Set $a_{0}=0$. Then there exist nonsingular billiard trajectories inside the ellipsoid $Q$ sharing the caustics $Q_{\lambda_{1}}, \ldots, Q_{\lambda_{n-1}}$ if and only if

$$
\begin{equation*}
\lambda_{k} \in\left(a_{k-1}, a_{k}\right) \cup\left(a_{k}, a_{k+1}\right), \quad k=1, \ldots, n-1 . \tag{2}
\end{equation*}
$$

Definition 1. The caustic type of a nonsingular trajectory is the vector $\varsigma=\left(\varsigma_{1}, \ldots, \varsigma_{n-1}\right) \in \mathbb{Z}^{n-1}$ such that

$$
\lambda_{k} \in\left(a_{\mathcal{S}_{k}}, a_{\mathcal{S}_{k}+1}\right), \quad k=1, \ldots, n-1 .
$$

We know from Proposition 3 that $\varsigma_{k} \in\{k-1, k\}$ for $k=1, \ldots, n-1$. Hence, there are exactly $2^{n-1}$ different caustic types. The two caustic types in the planar case correspond to ellipses: $\varsigma_{1}=0$, and hyperbolas: $\varsigma_{1}=1$. The four caustic types in the spatial case correspond to EH1: $\varsigma=(0,1), \mathrm{H} 1 \mathrm{H} 1: \varsigma=(1,1), \mathrm{EH} 2: \varsigma=(0,2)$, and H1H2: $\varsigma=(1,2)$. The notations EH1, H1H1, EH 2 , and H 1 H 2 were described in the introduction.

Next, we recall a result about periodic billiard trajectories inside ellipsoids.

Theorem 4 (Generalized Poncelet theorem). If a nonsingular billiard trajectory is closed after $m_{0}$ bounces and has length $L_{0}$, then all trajectories sharing the same caustics are also closed after $m_{0}$ bounces and have length $L_{0}$.

Poncelet proved this theorem for conics [6]. Darboux generalized it to triaxial ellipsoids of $\mathbb{R}^{3}$; see [8]. Later on, this result was generalized to any dimension in [9-12].

The periodic billiard trajectories sharing the same caustics also have the same winding numbers. In order to introduce these numbers, we set

$$
\begin{equation*}
\left\{c_{1}, \ldots, c_{2 n-1}\right\}=\left\{a_{1}, \ldots, a_{n}\right\} \cup\left\{\lambda_{1}, \ldots, \lambda_{n-1}\right\} \tag{3}
\end{equation*}
$$

and $f(t)=\sqrt{\prod_{i=1}^{2 n-1}\left(1-t / c_{i}\right)}$. We deal with nonsingular billiard trajectories inside ellipsoids without symmetries of revolution, so the parameters $c_{1}, \ldots, c_{2 n-1}$ are pairwise distinct, and we can assume that $c_{0}:=0<c_{1}<\cdots<c_{2 n-1}$.

Theorem 5 (Winding numbers). The nonsingular billiard trajectories inside the ellipsoid $Q$ sharing the caustics $Q_{\lambda_{1}}, \ldots, Q_{\lambda_{n-1}}$ are periodic with period $m_{0}$ if and only if there exist some positive integer numbers $m_{1}, \ldots, m_{n-1}$ such that

$$
\begin{equation*}
\sum_{j=0}^{n-1}(-1)^{j} m_{j} \int_{c_{2 j}}^{c_{2 j+1}} \frac{t^{i}}{f(t)} \mathrm{d} t=0, \quad \forall i=0, \ldots, n-2 \tag{4}
\end{equation*}
$$

Each of these periodic billiard trajectories has $m_{j}$ points at $Q_{c_{2 j}}$ and $m_{j}$ points at $Q_{c_{2 j+1}}$. Besides, $\left\{c_{2 j}, c_{2 j+1}\right\} \cap\left\{a_{1}, \ldots, a_{n}\right\} \neq \emptyset \Rightarrow m_{j}$ even. Finally, $\operatorname{gcd}\left(m_{0}, \ldots, m_{n-1}\right) \in\{1,2\}$.

Let $L_{0}$ be the common length of these periodic billiard trajectories. Let $x(t)$ be an arc-length parametrization of any of these trajectories. Let $\mu(t)=\left(\mu_{0}(t), \ldots, \mu_{n-1}(t)\right)$ be the corresponding parametrization in elliptic coordinates. Then:

1. $c_{2 j} \leq \mu_{j}(t) \leq c_{2 j+1}$ for all $t \in \mathbb{R}$.
2. Functions $\mu_{j}(t)$ are smooth everywhere, except $\mu_{0}(t)$ which is non-smooth at impact points -that is, when $x\left(t_{\star}\right) \in Q$-, in which case $\mu_{0}^{\prime}\left(t_{\star}+\right)=-\mu_{0}^{\prime}\left(t_{\star}-\right) \neq 0$.
3. If $\mu_{j}(t)$ is smooth at $t=t_{\star}$, then $\mu_{j}^{\prime}\left(t_{\star}\right)=0 \Leftrightarrow \mu_{j}\left(t_{\star}\right) \in\left\{c_{2 j}, c_{2 j+1}\right\}$.
4. $\mu_{j}(t)$ makes exactly $m_{j}$ complete oscillations (round trips) inside the interval $\left[c_{2 j}, c_{2 j+1}\right]$ along one period $0 \leq t \leq L_{0}$.
5. $\mu(t)$ has period $L=L_{0} / \operatorname{gcd}\left(m_{0}, \ldots, m_{n-1}\right)$.

Definition 2. The numbers $m_{0}, \ldots, m_{n-1}$ are called winding numbers. Theorem 5 contains three equivalent definitions for them: by means of the property regarding hyperelliptic integrals given in (4), as a geometric description of how the periodic billiard trajectories fold in $\mathbb{R}^{n}$, and as the number of oscillations of the elliptic coordinates along one period.

Most of the statements of Theorem 5 can be found in [18, 20], but the one about the even character of some winding numbers and the ones regarding $\operatorname{gcd}\left(m_{0}, \ldots, m_{n-1}\right)$. The first statement is trivial; suffice it to realize that a periodic billiard trajectory can only have an even number of crossings with any coordinate hyperplane. The second ones follow from the oscillating behaviour of elliptic coordinates along billiard trajectories described in Theorem 5; suffice it to note that all elliptic coordinates make an integer number of complete oscillations inside their corresponding intervals along one half-period $L_{0} / 2$ when $\operatorname{gcd}\left(m_{0}, \ldots, m_{n-1}\right)=2$.

The following conjecture was stated in [24], where it was numerically tested.

Conjecture 1. Winding numbers are always ordered in a strict decreasing way; that is,

$$
2 \leq m_{n-1}<\cdots<m_{1}<m_{0} .
$$

It is known that the conjecture holds in the planar case. If this conjecture holds, then any nonsingular periodic billiard trajectory inside $Q$ has period at least $n+1$. By the way, there are periodic billiard trajectories of smaller periods, but all of them are singular - they are contained in some coordinate hyperplane or in some ruled quadric of the confocal family.

In light of the last item of Theorem 5, we present the following definitions.
Definition 3. The elliptic period $m$ and the elliptic winding numbers $\widetilde{m}_{0}, \ldots, \widetilde{m}_{n-1}$ of the nonsingular periodic billiard trajectories with period $m_{0}$ and winding numbers $m_{0}, \ldots, m_{n-1}$ are

$$
m=m_{0} / d, \quad \widetilde{m}_{j}=m_{j} / d
$$

where $d=\operatorname{gcd}\left(m_{0}, \ldots, m_{n-1}\right)$.
Roughly speaking, the difference between the period $m_{0}$ and the elliptic period $m$ is that periodic billiard trajectories close in Cartesian (respectively, elliptic) coordinates after exactly $m_{0}$ (respectively, $m$ ) bounces. In order to clarify this difference, let us consider the six planar periodic trajectories shown in Figure 1; see Section 7. Only the trajectory in Figure 1(c) verifies that $m=m_{0}$. On the contrary, the trajectories in figures 1(a), 1(b), and 1(e) (respectively, Figure 1(d)) (respectively, Figure 1(f)) have even period $m_{0}$ and any of their impact points becomes its reflection with respect to the origin (respectively, the vertical axis) (respectively, the horizontal axis) after $m_{0} / 2$ bounces, so they have elliptic period $m=m_{0} / 2$.

It turns out that given any ellipsoid of the form (1) and any proper coordinate subspace of $\mathbb{R}^{n}$, there exist infinitely many sets of $n-1$ distinct nonsingular caustics such that their tangent trajectories are periodic with even period, say $m_{0}$, and any of their impact points becomes its reflection with respect to that coordinate subspace after $m_{0} / 2$ bounces. We will not prove this claim, since the proof requires some convoluted ideas developed in [24, 25].

It is natural to look for caustics giving rise to periodic billiard trajectories inside that ellipsoid. Such caustics can be found by means of certain algebraic conditions, called generalized Cayley conditions. They are found by working in elliptic coordinates, so they depend on the elliptic period $m$, not on the (Cartesian) period $m_{0}$.
Theorem 6 (Generalized Cayley conditions). The nonsingular billiard trajectories inside the ellipsoid $Q$ sharing the caustics $Q_{\lambda_{1}}, \ldots, Q_{\lambda_{n-1}}$ are periodic with elliptic period $m$ if and only if $m \geq n$ and

$$
\operatorname{rank}\left(\begin{array}{ccc}
f_{m+1} & \cdots & f_{n+1} \\
\vdots & & \vdots \\
f_{2 m-1} & \cdots & f_{m+n-1}
\end{array}\right)<m-n+1
$$

where $f(t)=\sum_{l \geq 0} f_{l} t^{l}:=\sqrt{\prod_{i=1}^{2 n-1}\left(1-t / c_{i}\right)}$.
Cayley proved this theorem for conics [13]. Later on, this result was generalized to any dimension by Dragović and Radnović in [15, 16]. These authors have also given similar Cayley conditions in many other billiard frameworks; see [17-21].

Definition 4. $C(m, n)$ denotes the generalized Cayley condition that characterizes billiard trajectories of elliptic period $m$ inside ellipsoids of $\mathbb{R}^{n}$ given in Theorem 6.

## 3. On the matrix formulation of the generalized Cayley conditions

The matrix formulation of the generalized Cayley condition $C(m, n)$ stated in Theorem 6 is very nice from a theoretical point of view, but has strong limitations from a practical point of view. Let us describe them.

The function $f(t)$ is symmetric in the inverse quantities $\gamma_{i}=1 / c_{i}$. In order to exploit it, we introduce some notations about symmetric polynomials. Let $\mathbb{Q}_{l}^{\text {hom,sym }}\left[x_{1}, \ldots, x_{s}\right]$ be the vectorial space over $\mathbb{Q}$ of all homogeneous symmetric polynomials with rational coefficients of degree $l$ in the variables $x_{1}, \ldots, x_{s}$. Let $\mathrm{e}_{l}\left(x_{1}, \ldots, x_{s}\right)$ be the elementary symmetric polynomial of degree $l$ in the variables $x_{1}, \ldots, x_{s}$. That is, $\prod_{i=1}^{s}\left(1+x_{i} t\right)=\sum_{l \geq 0} \mathrm{e}_{l}\left(x_{1}, \ldots, x_{s}\right) t^{l}$, so $\mathrm{e}_{l}\left(x_{1}, \ldots, x_{s}\right)=0$ for all $l>s$. Clearly, $\mathrm{e}_{l}=\mathrm{e}_{l}\left(x_{1}, \ldots, x_{s}\right) \in \mathbb{Q}_{l}^{\text {hom,sym }}\left[x_{1}, \ldots, x_{s}\right]$.

We stress that $f_{l}=f_{l}\left(\gamma_{1}, \ldots, \gamma_{2 n-1}\right) \in \mathbb{Q}_{l}^{\text {hom,sym }}\left[\gamma_{1}, \ldots, \gamma_{2 n-1}\right]$, which is one of the reasons for the introduction of the inverse quantities $\gamma_{i}=1 / c_{i}$. Indeed, using that $f^{2}(t)=\prod_{i=1}^{2 n-1}\left(1-\gamma_{i} t\right)$, we get the recursive relations

$$
f_{0}=1, \quad 2 f_{l}=(-1)^{l} \mathrm{e}_{l}\left(\gamma_{1}, \ldots, \gamma_{2 n-1}\right)-\sum_{k=1}^{l-1} f_{k} f_{l-k}, \quad \forall l \geq 1 .
$$

Hence, it is possible to compute recursively all Taylor coefficients $f_{l}$, although their expressions are rather complicated when $l$ is big. Nevertheless, the computation of the Taylor coefficients $f_{n+1}, \ldots, f_{2 m-1}$ is the simplest step in the practical implementation of the generalized Cayley condition $C(m, n)$. Next, we must impose that all $(m-n+1) \times(m-n+1)$ minors of the matrix that appear in Theorem 6 vanish. For simplicity, let us consider the minors formed by the first $m-n$ rows and the ( $m-n+l$ )-th row of that matrix, for $l=1, \ldots, n-1$. Then the Cayley condition $C(m, n)$ can be written as the system of $n-1$ polynomial equations

$$
M_{m, n, l}=M_{m, n, l}\left(\gamma_{1}, \ldots, \gamma_{2 n-1}\right):=\left|\begin{array}{ccc}
f_{m+1} & \cdots & f_{n+1}  \tag{5}\\
\vdots & & \vdots \\
f_{2 m-n} & \cdots & f_{m} \\
f_{2 m-n+l} & \cdots & f_{m+l}
\end{array}\right|=0, \quad 1 \leq l \leq n-1
$$

One can check that $M_{m, n, l}\left(\gamma_{1}, \ldots, \gamma_{2 n-1}\right) \in \mathbb{Q}_{(m-n+2) m-n+l}^{\text {hom,sym }}\left[\gamma_{1}, \ldots, \gamma_{2 n-1}\right]$ from the Leibniz formula for determinants. This implies that the resolution of system (5) is a formidable challenge, even from a purely numerical point of view and for relatively small values of $m$.

We want to write down the solutions of system (5) in an explicit algebraic way. Let us focus on the planar case $n=2$, when condition $C(m, 2)$ becomes a single homogeneous symmetric polynomial equation of degree $m^{2}-1$ in three unknowns; namely,

$$
M_{m}=M_{m}\left(\gamma_{1}, \gamma_{2}, \gamma_{3}\right):=\left|\begin{array}{ccc}
f_{m+1} & \cdots & f_{3} \\
\vdots & & \vdots \\
f_{2 m-1} & \cdots & f_{m+1}
\end{array}\right|=0
$$

For instance, condition $C(2,2)$ can be easily solved, since

$$
\begin{aligned}
-16 M_{2} & =\gamma_{1}^{3}+\gamma_{2}^{3}+\gamma_{3}^{3}-\gamma_{1}^{2} \gamma_{2}-\gamma_{1}^{2} \gamma_{3}-\gamma_{2}^{2} \gamma_{1}-\gamma_{2}^{2} \gamma_{3}-\gamma_{3}^{2} \gamma_{1}-\gamma_{3}^{2} \gamma_{2}+2 \gamma_{1} \gamma_{2} \gamma_{3} \\
& =\left(\gamma_{1}-\gamma_{2}-\gamma_{3}\right)\left(\gamma_{3}-\gamma_{1}-\gamma_{2}\right)\left(\gamma_{2}-\gamma_{3}-\gamma_{1}\right) .
\end{aligned}
$$

The inverse quantities $\gamma_{i}=1 / c_{i}$ verify that $0<\gamma_{3}<\gamma_{2}<\gamma_{1}$, since $0<c_{1}<c_{2}<c_{3}$. Therefore, only the first factor of the above formula provides a feasible solution, and so condition $C(2,2)$ has a unique solution:

$$
\gamma_{1}=\gamma_{2}+\gamma_{3} .
$$

The computations for condition $C(3,2)$ are much harder, so we have implemented them using a computer algebra system. We got the factorization $-16384 M_{3}=q_{0} q_{1} q_{2} q_{3}$, where

$$
\begin{aligned}
& q_{0}=\gamma_{1}^{2}+\gamma_{2}^{2}+\gamma_{3}^{2}-2 \gamma_{1} \gamma_{2}-2 \gamma_{1} \gamma_{3}-2 \gamma_{2} \gamma_{3}, \\
& q_{k}=3 \gamma_{k}^{2}-2\left(\gamma_{i}+\gamma_{j}\right) \gamma_{k}-\left(\gamma_{i}-\gamma_{j}\right)^{2}, \quad\{i, j, k\}=\{1,2,3\} .
\end{aligned}
$$

The first factor $q_{0}$ can, in its turn, be factored as

$$
q_{0}=\left(\sqrt{\gamma_{1}}-\sqrt{\gamma_{2}}-\sqrt{\gamma_{3}}\right)\left(\sqrt{\gamma_{1}}+\sqrt{\gamma_{2}}-\sqrt{\gamma_{3}}\right)\left(\sqrt{\gamma_{1}}-\sqrt{\gamma_{2}}+\sqrt{\gamma_{3}}\right)\left(\sqrt{\gamma_{1}}+\sqrt{\gamma_{2}}+\sqrt{\gamma_{3}}\right) .
$$

The factor $q_{0}$ provides a unique feasible solution: $\sqrt{\gamma_{1}}=\sqrt{\gamma_{2}}+\sqrt{\gamma_{3}}$, because $0<\gamma_{3}<\gamma_{2}<\gamma_{1}$. Next, we consider the factor $q_{k}$ as a second-order polynomial in the variable $\gamma_{k}$ with coefficients in $\mathbb{Z}^{\text {sym }}\left[\gamma_{i}, \gamma_{j}\right]$. Then we get the solutions

$$
\gamma_{k}=\gamma_{k}^{ \pm}\left(\gamma_{i}, \gamma_{j}\right):=\frac{\gamma_{i}+\gamma_{j}}{3} \pm \frac{2}{3} \sqrt{\gamma_{i}^{2}+\gamma_{j}^{2}-\gamma_{i} \gamma_{j}} .
$$

It turns out that $\gamma_{k}^{-} \leq 0<\max \left(\gamma_{i}, \gamma_{j}\right)<\gamma_{k}^{+}$, so only the factor $q_{1}$ gives a solution compatible with the ordering $0<\gamma_{3}<\gamma_{2}<\gamma_{1}$; namely, $\gamma_{1}=\gamma_{1}^{+}\left(\gamma_{2}, \gamma_{3}\right)$. Hence, $\mathcal{C}(3,2)$ has only two solutions:

$$
\begin{equation*}
\sqrt{\gamma_{1}}=\sqrt{\gamma_{2}}+\sqrt{\gamma_{3}} \quad \text { and } \quad 3 \gamma_{1}=\gamma_{2}+\gamma_{3}+2 \sqrt{\gamma_{2}^{2}+\gamma_{3}^{2}-\gamma_{2} \gamma_{3}} \tag{6}
\end{equation*}
$$

We have tried to write down explicitly the solutions of system (5) in other cases, but we did not succed, even after implementing the computations in a computer algebra system. This shows the limitations of the matrix formulation.
4. A polynomial formulation of the generalized Cayley conditions

Theorem 7. Let $r(t)=\prod_{i=1}^{2 n-1}\left(1-t / c_{i}\right)$ and $f(t)=\sqrt{r(t)}$. The generalized Cayley condition $C(m, n)$ is equivalent to any of the following two conditions:
(i) There exists a non-zero polynomial $s(t) \in \mathbb{R}_{m-n}[t]$ such that

$$
\begin{equation*}
\left.\frac{\mathrm{d}^{l}}{\mathrm{~d} t^{l}}\right|_{t=0}\{s(t) f(t)\}=0, \quad l=m+1, \ldots, 2 m-1 \tag{7}
\end{equation*}
$$

(ii) There exist $\alpha \neq 0, s(t) \in \mathbb{R}_{m-n}[t]$, and $q(t) \in \mathbb{R}_{m-1}[t]$ such that $s(0)=q(0)=1$ and

$$
\begin{equation*}
s^{2}(t) r(t)=\left(\alpha t^{m}+q(t)\right) q(t) \tag{8}
\end{equation*}
$$

Proof. We split the proof in three steps.
Step 1: $\mathcal{C}(m, n) \Leftrightarrow(i) . \mathcal{C}(m, n)$ means that the $m-n+1$ columns of the matrix given in Theorem 6 are linearly dependent, so there exist $s_{0}, \ldots, s_{m-n} \in \mathbb{R}$, not all zero, such that

$$
\begin{gathered}
s_{0} \times(\text { first column })+\cdots+s_{m-n} \times(\text { last column })=0, \\
8
\end{gathered}
$$

which is equivalent to condition (7) when $s(t)=\sum_{l=0}^{m-n} s_{l} t^{l} \in \mathbb{R}_{m-n}[t]$.
Step 2: $(i) \Rightarrow(i i)$. If $s(t) \in \mathbb{R}_{m-n}[t]$ verifies (7), then $g(t)=\sum_{l \geq 0} g_{l} t^{l}:=s(t) f(t)$ verifies that $g_{l}=0$ for $l=m+1, \ldots, 2 m-1$. Hence,

$$
g(t)=q(t)+\alpha t^{m} / 2+\mathrm{O}\left(t^{2 m}\right)
$$

where $q(t)=g_{0}+\cdots+g_{m-1} t^{m-1} \in \mathbb{R}_{m-1}[t]$ and $\alpha=2 g_{m}$. Therefore,

$$
s^{2} r=s^{2} f^{2}=g^{2}=q^{2}+\alpha t^{m} q+\mathrm{O}\left(t^{2 m}\right)=\left(q+\alpha t^{m}\right) q+\mathrm{O}\left(t^{2 m}\right),
$$

and so $s^{2} r=\left(q+\alpha t^{m}\right) q$, since $\operatorname{deg}\left[s^{2} r\right] \leq 2 m-1$ and $\operatorname{deg}\left[\left(q+\alpha t^{m}\right) q\right] \leq 2 m-1$. Besides, $\alpha \neq 0$, because $\operatorname{deg}[r]$ is odd and $s(t) \not \equiv 0$.

Let $h(t)=\sum_{l \geq 0} h_{l} t^{l}:=g^{2}(t)=s^{2}(t) r(t) \in \mathbb{R}_{2 m-1}[t]$. Then

$$
0=h_{2 m}=\sum_{l=0}^{2 m} g_{l} g_{2 m-l}=\left(g_{m}\right)^{2}+2 g_{0} g_{2 m} \Rightarrow g_{0} g_{2 m}=-\alpha^{2} / 8 \neq 0
$$

From this property, we deduce that $q(0)=g_{0} \neq 0$ and $s^{2}(0)=q^{2}(0) / r(0) \neq 0$. Thus, we can normalize $s(t)$ by imposing $s(0)=1$, since condition (7) only determines $s(t)$ up to a multiplicative constant. This implies that $q^{2}(0)=s^{2}(0) r(0)=1$, so $q(0)= \pm 1$. We can assume, without loss of generality, that $q(0)=1$. On the contrary, we substitute $q(t)$ and $\alpha$ in the identity $s^{2} r=\left(\alpha t^{m}+q\right) q$, by $-q(t)$ and $-\alpha$, respectively.

Step 3: (ii) $\Rightarrow(i)$. If there exist $\alpha \neq 0, s(t) \in \mathbb{R}_{m-n}[t]$, and $q(t) \in \mathbb{R}_{m-1}[t]$ such that $s(0)=$ $q(0)=1$ and relation (8) holds, we set $g(t)=\sum_{l \geq 0} g_{l} t^{l}:=s(t) f(t)$. Then,

$$
g^{2}=s^{2} f^{2}=s^{2} r=\left(q+\alpha t^{m}\right) q=\left(1+\alpha t^{m} / q\right) q^{2}
$$

The last operation is well-defined for small values of $|t|$, because $q(0) \neq 0$. Hence,

$$
g= \pm q \sqrt{1+\frac{\alpha t^{m}}{q}}= \pm q\left(1+\frac{\alpha t^{m}}{2 q}+\mathrm{O}\left(t^{2 m}\right)\right)= \pm q \pm \frac{\alpha t^{m}}{2}+\mathrm{O}\left(t^{2 m}\right)
$$

so $g_{l}=0$ for $l=m+1, \ldots, 2 m-1$.
Next, we present three examples of the results that can be obtained from this formulation.
Theorem 8. The nonsingular billiard trajectories inside the ellipsoid (1) sharing the caustics $Q_{\lambda_{1}}, \ldots, Q_{\lambda_{n-1}}$ are periodic with:

- Elliptic period $m=n$ if the roots of $t^{n}-\prod_{j=1}^{n}\left(t-a_{j}\right)$ are the caustic parameters;
- Elliptic period $m=n+1$ if there exists $d \in \mathbb{R}$ such that the roots of $t^{n+1}-(t-d)^{2} \prod_{k=1}^{n-1}\left(t-\lambda_{k}\right)$ are the ellipsoidal parameters; and
- Elliptic period $m=2 n-1$ if all the roots of $t^{2 n-1}-\prod_{j=1}^{n}\left(t-a_{j}\right) \prod_{k=1}^{n-1}\left(t-\lambda_{k}\right)$ are double.

Proof. Suffice it to find $\alpha \neq 0, s(t) \in \mathbb{R}_{m-n}[t]$, and $q(t) \in \mathbb{R}_{m-1}[t]$ such that

$$
s^{2}(t) r(t)=\left(\alpha t^{m}+q(t)\right) q(t), \quad s(0)=q(0)=1
$$

for $m=n, m=n+1$, and $m=2 n-1$, respectively; see Theorem 7. We recall that $r(t)=$ $\prod_{i=1}^{2 n-1}\left(1-t / c_{i}\right)$ with $\left\{c_{1}, \ldots, c_{2 n-1}\right\}=\left\{a_{1}, \ldots, a_{n}\right\} \cup\left\{\lambda_{1}, \ldots, \lambda_{n-1}\right\}$.

Case $m=n$. If the caustic parameters are the roots of $t^{n}-\prod_{j=1}^{n}\left(t-a_{j}\right)$, then

$$
t^{n}-\prod_{j=1}^{n}\left(t-a_{j}\right)=\kappa \prod_{k=1}^{n-1}\left(t-\lambda_{k}\right),
$$

for some factor $\kappa \in \mathbb{R}$. Indeed, $\kappa=\prod_{j=1}^{n} a_{j} \prod_{k=1}^{n-1} \lambda_{k}^{-1}$. We take $\alpha=(-1)^{n} \prod_{j=1}^{n} a_{j}^{-1}, s(t) \equiv 1$, and $q(t)=\prod_{k=1}^{n-1}\left(1-t / \lambda_{k}\right)$. Clearly, $\alpha \neq 0, s(t) \in \mathbb{R}_{0}[t], q(t) \in \mathbb{R}_{n-1}[t]$, and $s(0)=q(0)=1$. Besides,

$$
s^{2}(t) r(t)=\prod_{j=1}^{n}\left(1-t / a_{j}\right) \prod_{k=1}^{n-1}\left(1-t / \lambda_{k}\right)=\left(\alpha t^{n}+q(t)\right) q(t),
$$

since $\prod_{j=1}^{n}\left(1-t / a_{j}\right)=\alpha \prod_{j=1}^{n}\left(t-a_{j}\right)=\alpha t^{n}-\alpha \kappa \prod_{k=1}^{n-1}\left(t-\lambda_{k}\right)=\alpha t^{n}+q(t)$.
Case $m=n+1$. If $a_{1}, \ldots, a_{n}$ are the roots of $t^{n+1}-(t-d)^{2} \prod_{k=1}^{n-1}\left(t-\lambda_{k}\right)$, then

$$
t^{n+1}-(t-d)^{2} \prod_{k=1}^{n-1}\left(t-\lambda_{k}\right)=\kappa \prod_{j=1}^{n}\left(t-a_{j}\right)
$$

for some $\kappa \in \mathbb{R}$. Indeed, $\kappa=d^{2} \prod_{k=1}^{n-1} \lambda_{k} \prod_{j=1}^{n} a_{j}^{-1}$. We take $\alpha=(-1)^{n+1} d^{-2} \prod_{k=1}^{n-1} \lambda_{k}^{-1}, s(t)=$ $(1-t / d)$, and $q(t)=\prod_{j=1}^{n}\left(1-t / a_{j}\right)$. Clearly, $\alpha \neq 0, s(t) \in \mathbb{R}_{1}[t], q(t) \in \mathbb{R}_{n}[t]$, and $s(0)=q(0)=1$. Besides,

$$
s^{2}(t) r(t)=(1-t / d)^{2} \prod_{j=1}^{n}\left(1-t / a_{i}\right) \prod_{k=1}^{n-1}\left(1-t / \lambda_{k}\right)=\left(\alpha t^{n+1}+q(t)\right) q(t),
$$

since $(1-t / d)^{2} \prod_{k=1}^{n-1}\left(1-t / \lambda_{k}\right)=\alpha(t-d)^{2} \prod_{k=1}^{n-1}\left(1-t / \lambda_{k}\right)=\alpha t^{n+1}-\alpha \kappa \prod_{j=1}^{n}\left(t-a_{j}\right)=\alpha t^{n+1}+q(t)$.
Case $m=2 n-1$. If $t^{2 n-1}-\prod_{i=1}^{2 n-1}\left(t-c_{i}\right)$ has double roots $d_{1}, \ldots, d_{n-1}$, then

$$
t^{2 n-1}-\prod_{i=1}^{2 n-1}\left(t-c_{i}\right)=\kappa \prod_{l=1}^{n-1}\left(t-d_{l}\right)^{2},
$$

for some $\kappa \in \mathbb{R}$. Indeed, $\kappa=\prod_{i=1}^{2 n-1} c_{i} \prod_{l=1}^{n-1} d_{l}^{-2}$. We take $\alpha=-\prod_{i=1}^{2 n-1} c_{i}^{-1}, s(t)=\prod_{l=1}^{n-1}\left(1-t / d_{l}\right)$, and $q(t)=s^{2}(t)$. Clearly, $\alpha \neq 0, s(t) \in \mathbb{R}_{n-1}[t], q(t) \in \mathbb{R}_{2 n-2}[t]$, and $s(0)=q(0)=1$. Besides,

$$
s^{2}(t) r(t)=\prod_{l=1}^{n-1}\left(1-t / d_{l}\right)^{2} \prod_{i=1}^{2 n-1}\left(1-t / c_{i}\right)=\left(\alpha t^{2 n-1}+q(t)\right) q(t),
$$

since $\prod_{i=1}^{2 n-1}\left(1-t / c_{i}\right)=\alpha \prod_{i=1}^{2 n-1}\left(t-c_{i}\right)=\alpha t^{2 n-1}-\alpha \kappa \prod_{l=1}^{n-1}\left(t-d_{l}\right)^{2}=\alpha t^{2 n-1}+q(t)$.
Several questions arise about the periodic trajectories found in the previous theorem. Let us mention just three. Which are their caustic types, their (Cartesian) periods, and their winding numbers? Inside what ellipsoids exist them? Are there other nonsingular periodic billiard trajectories with elliptic period three, four or five?

We will give some partial answers in the next sections.
Some technicalities become simpler after the change of variables $t=1 / x$. Thus, we state another polynomial formulation of the generalized Cayley condition $C(m, n)$.

Proposition 9. Let $R(x)=x \prod_{i=1}^{2 n-1}\left(x-\gamma_{i}\right)$, where $\gamma_{i}=1 / c_{i}$. The generalized Cayley condition $C(m, n)$ holds if and only if there exist two monic polynomials $S(x), R(x) \in \mathbb{R}[x]$ such that $\operatorname{deg}[S]=m-n, \operatorname{deg}[P]=m, P(0) \neq 0$, and

$$
\begin{equation*}
S^{2}(x) R(x)=P(x)(P(x)-P(0)) \tag{9}
\end{equation*}
$$

Furthermore, if such polynomials $S(x)$ and $P(x)$ exist, the following properties hold:

1. $S(x)$ has no multiple roots;
2. All the real roots of $S(x)$ are contained in $\{x \in \mathbb{R}: R(x)<0\}$;
3. All the roots of $S(x)$ are real when $m \leq n+3$; and
4. $P(x)$ and $P(x)-P(0)$ have the same number of real roots (counted with multiplicity).

Proof. The "if and only if" follows directly from the change of variables $t=1 / x$. Concretely, the relation between the objects of identities (8) and (9) is

$$
R(x)=x^{2 n} r(1 / x), \quad S(x)=x^{m-n} s(1 / x), \quad P(x)=\alpha+x^{m} q(1 / x)
$$

Then $P(0) \neq 0$ if and only if $\alpha \neq 0, s(0)=1$ if and only if $S(x)$ is a monic polynomial of degree $m-n$, and $q(0)=1$ if and only if $P(x)$ is a monic polynomial of degree $m$.

To prove the first two properties, suffice it to prove that $\operatorname{gcd}\left[S, R S^{\prime}\right]=1$ and

$$
l_{+}:=\#\{x \in \mathbb{R}: S(x)=0<R(x)\}=0 .
$$

If $W(x)=P(x)(P(x)-P(0))$ and $T(x)=P(x)-P(0) / 2$, we get from (9) that

$$
\begin{aligned}
W(x) & =S^{2}(x) R(x)=T^{2}(x)-P^{2}(0) / 4 \\
W^{\prime}(x) & =S(x)\left(S(x) R^{\prime}(x)+2 R(x) S^{\prime}(x)\right)=2 T(x) P^{\prime}(x)
\end{aligned}
$$

We consider the factorization $W^{\prime}(x)=2 m W_{-}(x) W_{0}(x) W_{+}(x) W_{*}(x)$, where if $z \in \mathbb{C}$ is a root of multiplicity $\beta$ of $W^{\prime}(x)$ such that $W(z)<0, W(z)=0, W(z)>0$, or $W(z) \notin \mathbb{R}$, then $(x-z)^{\beta}$ is included in the monic factor $W_{-}(x), W_{0}(x), W_{+}(x)$, or $W_{*}(x)$, respectively. Next, we find some lower bounds of the degrees of these factors.

First, $T$ is divisor of $W_{-}$, because $W$ takes the negative value $-P^{2}(0) / 4$ at each root of $T$. Hence, $\operatorname{deg}\left[W_{-}\right] \geq \operatorname{deg}[T]=m$. Second, $S \operatorname{gcd}\left[S, R S^{\prime}\right]$ is a divisor of $W_{0}$, because $W$ vanishes at each root of $S$. Thus, $\operatorname{deg}\left[W_{0}\right] \geq m-n+l_{0}$, where $l_{0}$ denotes the degree of $\operatorname{gcd}\left[S, R S^{\prime}\right]$. Third,

$$
\operatorname{deg}\left[W_{+}\right] \geq \#\left\{\begin{array}{ll} 
& W(a)=W(b)=0 \\
(a, b) \subset \mathbb{R}: & R(x)>0 \text { for all } x \in(a, b) \\
& S(x) \neq 0 \text { for all } x \in(a, b)
\end{array}\right\}=n-1+l_{+} .
$$

To understand the above inequality, we realize that if $(a, b)$ is an open bounded interval that satisfies the above three properties, then $W(x)=S^{2}(x) R(x)>0$ for all $x \in(a, b)$, and $W^{\prime}(x)$ vanishes at some point $c \in(a, b)$, by Rolle's Theorem. Therefore, $\operatorname{deg}\left[W_{+}\right]$is at least the number of such intervals. We combine these three lower bounds:

$$
2 m-1=\operatorname{deg}\left[W^{\prime}\right] \geq \operatorname{deg}\left[W_{-}\right]+\operatorname{deg}\left[W_{0}\right]+\operatorname{deg}\left[W_{+}\right] \geq 2 m-1+l_{0}+l_{+}
$$

This implies that $l_{0}=l_{+}=0$. Indeed, $W_{-}=T, W_{0}=S, W_{*}=1$, and $\operatorname{gcd}\left[S, R S^{\prime}\right]=1$.
Next, we prove the property about the number of roots of $P(x)$ and $P(x)-P(0)$. Let $z$ be a root of the derivative $P^{\prime}$. Since $W^{\prime}=2 T P^{\prime}$ and $W_{-}=T$, we deduce that $W(z)$ cannot be a negative number. This implies that if $P(z)$ is a real value between 0 and $P(0)$, then $P^{\prime}(z) \neq 0$, since $W(z)=P(z)(P(z)-P(0))<0$. In particular, we deduce that the number of real roots (counted with multiplicity) of the polynomial $P(x)-\eta$ does not change when the constant $\eta \in \mathbb{R}$ moves from 0 to $P(0)$.

Finally, we prove that $S(x)$ has only real roots when $m \leq n+3$. Let us suppose that $z \notin \mathbb{R}$ is a root of $S(x)$. Then $\bar{z}$ is also a root of $S(x)$, so $(x-z)(x-\bar{z}) \mid S(x)$. Using the identity $S^{2}(x) R(x)=$
$P(x)(P(x)-P(0))$, we get that $(x-z)^{2}(x-\bar{z})^{2}$ is either a divisor of $P(x)$ or a divisor of $P(x)-P(0)$, since $P(x)$ and $P(x)-P(0)$ have no common factors. But $P(x)$ and $P(x)-P(0)$ have the same number of real roots, so there exists another $w \notin \mathbb{R} \cup\{z, \bar{z}\}$ such that $(x-w)^{2}(x-\bar{w})^{2}$ is a divisor of $P(x)(P(x)-P(0))$. This implies that $S(x)$ has, at least, four different complex roots, and so $m-n=\operatorname{deg}[S] \geq 4$.

There are some theoretical arguments against the existence of non-real roots of polynomial $S(x)$, although we have not been able to prove it.
Conjecture 2. Let $R(x)=x \prod_{i=1}^{2 n-1}\left(x-\gamma_{i}\right)$ with $0<\gamma_{2 n-1}<\cdots<\gamma_{1}$. If relation (9) holds for some polynomials $S(x), P(x) \in \mathbb{R}[x]$ such that $P(0) \neq 0$, then $S(x)$ has only real roots.

## 5. Generalized Cayley conditions in the minimal case

Let us consider the case of minimal elliptic periods; that is, $m=n$.
We begin with a technical lemma to describe how the roots of the polynomials of the form $P(x)(P(x)-P(0))$ with $P(0) \neq 0$ are ordered in the real line, assuming that all these roots are positive -but the trivial one-, and have multiplicity at most two.
Lemma 10. Let $P(x) \in \mathbb{R}[x]$ be a monic polynomial of degree $m$ such that $P(0) \neq 0$ and all the roots of $P(x)(P(x)-P(0))$ are positive -but a simple root at $x=0$-, and have multiplicity at most two. Let $\alpha_{m} \leq \cdots \leq \alpha_{1}$ be the positive roots of $P(x)$. Let $\beta_{m-1} \leq \cdots \leq \beta_{1}$ be the positive roots of $P(x)-P(0)$.

If $m$ is odd, then $\beta_{2 l-1}, \beta_{2 l} \in\left(\alpha_{2 l}, \alpha_{2 l-1}\right)$ for all $l=1, \ldots,(m-1) / 2$; so

$$
0<\alpha_{m} \leq \alpha_{m-1}<\beta_{m-1} \leq \beta_{m-2}<\alpha_{m-2} \leq \alpha_{m-3}<\cdots<\alpha_{3} \leq \alpha_{2}<\beta_{2} \leq \beta_{1}<\alpha_{1}
$$

If $m$ is even, then $\beta_{1}>\alpha_{1}$ and $\beta_{2 l}, \beta_{2 l+1} \in\left(\alpha_{2 l+1}, \alpha_{2 l}\right)$ for all $l=1, \ldots,(m-2) / 2$; so

$$
0<\alpha_{m} \leq \alpha_{m-1}<\beta_{m-1} \leq \beta_{m-2}<\alpha_{m-2} \leq \alpha_{m-3}<\cdots<\beta_{3} \leq \beta_{2}<\alpha_{2} \leq \alpha_{1}<\beta_{1}
$$

Proof. Let $\eta \in \mathbb{R}$. Using that the only critical points of $P(x)$ are non-degenerate local maxima or non-degenerate local minima, we deduce that the polynomial $P(x)-\eta$ has $m$ real roots (counted with multiplicity) if and only if $\underline{\eta} \leq \eta \leq \bar{\eta}$, where

$$
\begin{aligned}
\bar{\eta} & =\min \{P(\bar{x}): \bar{x} \text { is a non-degenerate local maximum of } P(x)\}, \\
\underline{\eta} & =\max \{P(\underline{x}): \underline{x} \text { is a non-degenerate local minimum of } P(x)\} .
\end{aligned}
$$

Therefore, $\underline{\eta} \leq \min (0, P(0))$ and $\bar{\eta} \geq \max (0, P(0))$.
We begin with the case $m$ odd, so $P(0)=(-1)^{m} \prod_{j=1}^{m} \alpha_{j}<0, \eta \leq P(0)$, and $\bar{\eta} \geq 0$. The roots of $P(x)$ and $P(x)-P(0)$ can be viewed as the abscissae of the intersections of the graph $\{y=P(x)\}$ with the horizontal line $\{y=0\}$ and $\{y=P(0)\}$, respectively. Double roots correspond to tangential intersections. We know that $P(\bar{x}) \geq \bar{\eta} \geq 0$ at the local maxima, and $P(\underline{x}) \leq \eta \leq P(0)$ at the local minima. This means that the intersections of the graph $\{y=P(x)\}$ with $\overline{\text { the lines }}$ $\{y=0\}$ and $\{y=P(0)\}$ have the following pattern from left to right. First, the graph crosses $\{y=P(0)\}$ at the abscissa $x=0$; second, it intersects $\{y=0\}$ at two abscissae $\alpha_{m}$ and $\alpha_{m-1}$, which may coincide giving rise to a double root of $P(x)$; third, it intersects $\{y=P(0)\}$ at two abscissae $\beta_{m-1}$ and $\beta_{m-2}$, which may coincide giving rise to a double root of $P(x)-P(0)$; fourth, it intersects $\{y=0\}$ at two abscissae $\alpha_{m-2}$ and $\alpha_{m-3}$, which may coincide giving rise to a double root of $P(x)$; and so on. The last intersection correspond to the abscissa $x=\alpha_{1}$.

The proof for $m$ even is similar. We skip the details.

We emphasize that ellipsoidal parameters $0<a_{1}<\cdots<a_{n}$ and nonsingular caustic parameters $\lambda_{1}<\cdots<\lambda_{n-1}$ verify restrictions (2); then the parameters $0<c_{1}<\cdots<c_{2 n-1}$ are defined in (3); next the inverse quantities $0<\gamma_{2 n-1}<\cdots<\gamma_{1}$ are given by $\gamma_{i}=1 / c_{i}$; and finally, $R(x)=x \prod_{i=1}^{2 n-1}\left(x-\gamma_{i}\right)$. We will make use of these notations, orderings, and conventions along the paper without any explicit mention.

Corollary 11. Let $\{1, \ldots, 2 n-1\}=J_{n} \cup K_{n}$ be the decomposition defined by

$$
J_{1}=\{1\}, \quad J_{2}=\{2,3\}, \quad J_{n}=J_{n-2} \cup\{2 n-2,2 n-1\}, \quad K_{n}=J_{n-1} .
$$

If $P(x) \in \mathbb{R}[x]$ is a monic polynomial of degree $n$ such that $P(0) \neq 0$ and

$$
R(x)=P(x)(P(x)-P(0)),
$$

then $P(x)=\prod_{j \in J_{n}}\left(x-\gamma_{j}\right)=P(0)+x \prod_{k \in K_{n}}\left(x-\gamma_{k}\right)$.
Proof. There exists a decomposition $\{1, \ldots, 2 n-1\}=J^{\prime} \cup K^{\prime}$ such that $\# J^{\prime}=n, \# K^{\prime}=n-1$, and $P(x)=\prod_{j \in J^{\prime}}\left(x-\gamma_{j}\right)=P(0)+x \prod_{k \in K^{\prime}}\left(x-\gamma_{k}\right)$. The polynomial $P(x)$ verifies the hypotheses stated in Lemma 10, so the roots $\left\{\alpha_{1}, \ldots, \alpha_{m}\right\}=\left\{\gamma_{j}: j \in J^{\prime}\right\}$ and $\left\{\beta_{1}, \ldots, \beta_{m-1}\right\}=\left\{\gamma_{k}: k \in K^{\prime}\right\}$ obey the ordering described in that lemma. Therefore, $J^{\prime}=J_{n}$ and $K^{\prime}=K_{n}$.

We rewrite now the generalized Cayley condition $C(n, n)$ using the previous results. For brevity, we omit the dependence of the decomposition $\{1, \ldots, 2 n-1\}=J \cup K$ on the index $n$. We note that $\# J=n$ and $\# K=n-1$. The symbol $\mathrm{e}_{l}$ ("a set of parameters") denotes the elementary symmetric polynomial of degree $l$ in those parameters.

Proposition 12. $C(n, n)$ is equivalent to any of the following four conditions:
(i) If $P(x)=\prod_{j \in J}\left(x-\gamma_{j}\right)$, then $P(x)-P(0)=x \prod_{k \in K}\left(x-\gamma_{k}\right)$.
(ii) $\prod_{j \in J}\left(\gamma_{j}-\gamma_{k}\right)=\prod_{j \in J} \gamma_{j}$, for all $k \in K$.
(iii) $\mathrm{e}_{l}\left(\left\{\gamma_{j}\right\}_{j \in J}\right)=\mathrm{e}_{l}\left(\left\{\gamma_{k}\right\}_{k \in K}\right)$, for all $l=1, \ldots, n-1$.
(iv) $\sum_{j \in J} \gamma_{j}^{l}=\sum_{k \in K} \gamma_{k}^{l}$, for all $l=1, \ldots, n-1$.

Proof. We split the proof in four steps.
Step 1: $\mathcal{C}(n, n) \Leftrightarrow(i)$. Let us assume that $\mathcal{C}(n, n)$ holds. Then there exist a monic polynomial $P(x) \in \mathbb{R}[x]$ of degree $n$ such that $P(0) \neq 0$ and

$$
R(x)=P(x)(P(x)-P(0)) .
$$

Thus, condition (i) follows from Corollary 11.
Reciprocally, if condition (i) holds, $P(x)(P(x)-P(0))=x \prod_{i=1}^{2 n-1}\left(x-\gamma_{i}\right)$, so $C(n, n)$ holds.
Step 2: (i) $\Leftrightarrow$ (ii). If $P(x)=\prod_{j \in J}\left(x-\gamma_{j}\right)$ and $Q(x)=x \prod_{k \in K}\left(x-\gamma_{k}\right)$, then

$$
Q(x)=P(x)-P(0) \Leftrightarrow P\left(\gamma_{k}\right)=P(0), \quad \forall k \in K \Leftrightarrow \prod_{j \in J}\left(\gamma_{j}-\gamma_{k}\right)=\prod_{j \in J} \gamma_{j}, \quad \forall k \in K .
$$

Step 3: $(i) \Leftrightarrow($ iii $)$. If $P(x)=\prod_{j \in J}\left(x-\gamma_{j}\right)$ and $Q(x)=x \prod_{k \in K}\left(x-\gamma_{k}\right)$, then

$$
Q(x)=x^{n}+\sum_{l=1}^{n-1}(-1)^{l} \mathrm{e}_{l}\left(\left\{\gamma_{k}\right\}_{k \in K}\right) x^{n-l}, \quad P(x)=x^{n}+\sum_{l=1}^{n-1}(-1)^{l} \mathrm{e}_{l}\left(\left\{\gamma_{j}\right\}_{j \in J}\right) x^{n-l}+P(0) .
$$

Step 4: (iii) $\Leftrightarrow$ (iv). It follows from the Newton's identities connecting the elementary symmetric polynomials and the power sum symmetric polynomials; see [26].

Example 1. The quantities $\gamma_{3}=1, \gamma_{2}=2$, and $\gamma_{1}=3$ verify $C(2,2)$, because $1+2=3$. This means that the billiard trajectories

- Inside the ellipse $Q: x^{2}+2 y^{2}=1$ with caustic parameter $\lambda=1 / 3$; or
- Inside the ellipse $Q: x^{2}+3 y^{2}=1$ with caustic parameter $\lambda=1 / 2$;
are periodic with elliptic periodic $m=2$. Their caustic types are $\varsigma=0$ and $\varsigma=1$, respectively.
Example 2. The quantities $\gamma_{5}=1, \gamma_{4}=2, \gamma_{3}=4, \gamma_{2}=5$, and $\gamma_{1}=6$ verify $C(3,3)$, because $1+2+6=4+5$ and $1^{2}+2^{2}+6^{2}=4^{2}+5^{2}$. Hence, the billiard trajectories
- Inside the ellipsoid $Q: x^{2}+2 y^{2}+5 z^{2}=1$ with caustic parameters $\lambda_{1}=\frac{1}{6}$ and $\lambda_{2}=\frac{1}{4}$; or
- Inside the ellipsoid $Q: x^{2}+2 y^{2}+6 z^{2}=1$ with caustic parameters $\lambda_{1}=\frac{1}{5}$ and $\lambda_{2}=\frac{1}{4}$; or
- Inside the ellipsoid $Q: x^{2}+4 y^{2}+5 z^{2}=1$ with caustic parameters $\lambda_{1}=\frac{1}{6}$ and $\lambda_{2}=\frac{1}{2}$; or
- Inside the ellipsoid $Q: x^{2}+4 y^{2}+6 z^{2}=1$ with caustic parameters $\lambda_{1}=\frac{1}{5}$ and $\lambda_{2}=\frac{1}{2}$;
are periodic with elliptic periodic $m=3$. Their caustic types are $\varsigma=(0,1), \varsigma=(1,1), \varsigma=(0,2)$, and $\varsigma=(1,2)$, respectively. We recall that these caustic types were denoted EH1, H1H1, EH2, and H 1 H 2 in the introduction.

Let us compare the system of homogeneous symmetric polynomial equations (5), which was obtained directly from the matrix formulation, with the system of homogeneous non-symmetric polynomial equations $\sum_{j \in J} \gamma_{j}^{l}=\sum_{k \in K} \gamma_{k}^{l}, 1 \leq l \leq n-1$, obtained in the previous proposition. We are dealing with the case $m=n$, so the $l$-th equation of the former system has degree $n+l$, whereas the $l$-th equation of the new system has degree $l$. Besides, the new system has a remarkably simple closed expression. This shows that the polynomial formulation simplifies the problem.

The beauty of the conditions regarding the elementary symmetric polynomials and the power sum symmetric polynomials given in Proposition 12 has been the motivation for the introduction of the inverse quantities $\gamma_{i}=1 / c_{i}$. Nevertheless, we find useful to state the following result in terms of the ellipsoidal parameters $a_{j}$, in order to answer some questions about the nonsingular periodic billiard trajectories found in the first item of Theorem 8.
Theorem 13. There exist nonsingular periodic billiard trajectories inside the ellipsoid (1) with elliptic period $m=n$ and caustic type

$$
\varsigma= \begin{cases}(1,1,3,3, \ldots, n-2, n-2) & \text { for } n \text { odd }  \tag{10}\\ (0,2,2,4,4, \ldots, n-2, n-2) & \text { for } n \text { even }\end{cases}
$$

if and only if all the roots of $t^{n}-\prod_{j=1}^{n}\left(t-a_{j}\right)$ are real and simple. These periodic billiard trajectories have the roots of $t^{n}-\prod_{j=1}^{n}\left(t-a_{j}\right)$ as caustic parameters, period $m_{0}=2 n$, and even winding numbers $m_{0}, \ldots, m_{n-1}$. Indeed,

$$
\begin{equation*}
m_{j}=2 \widetilde{m}_{j}=2(n-j), \quad j=0, \ldots, n-1 \tag{11}
\end{equation*}
$$

provided Conjecture 1 on the strict decreasing ordering of winding numbers holds.
Proof. Let us assume that there exist nonsingular periodic billiard trajectories with elliptic period $n$ and caustic type (10). By definition of caustic type, the caustic parameters $\lambda_{1}<\cdots<\lambda_{n-1}$ of such trajectories verify that

- If $n$ is odd, then $\lambda_{2 l-1}, \lambda_{2 l} \in\left(a_{2 l-1}, a_{2 l}\right)$, for $l=1, \ldots,(n-1) / 2$;
- If $n$ is even, then $\lambda_{1} \in\left(0, a_{1}\right)$, and $\lambda_{2 l}, \lambda_{2 l+1} \in\left(a_{2 l}, a_{2 l+1}\right)$, for $l=1, \ldots,(n-2) / 2$.

Hence, we can split the set $\left\{\gamma_{i}=1 / c_{i}: i=1, \ldots, 2 n-1\right\}$ as the disjoint union of the sets

$$
\left\{\gamma_{j}: j \in J\right\}=\left\{1 / a_{1}, \ldots, 1 / a_{n}\right\}, \quad\left\{\gamma_{k}: k \in K\right\}=\left\{1 / \lambda_{1}, \ldots, 1 / \lambda_{n-1}\right\}
$$

where $\{1, \ldots, 2 n-1\}=J \cup K$ is the decomposition described in Corollary 11. Thus, we know from condition iv of Proposition 12 that

$$
\prod_{j=1}^{n}\left(\frac{1}{a_{j}}-\frac{1}{\lambda_{k}}\right)=\prod_{j=1}^{n} \frac{1}{a_{j}}, \quad k=1, \ldots, n-1 .
$$

This identity can be written as $\lambda_{k}^{n}=\prod_{j=1}^{n}\left(\lambda_{k}-a_{j}\right)$, for all $k=1, \ldots, n-1$, which implies that the caustic parameters $\lambda_{1}, \ldots, \lambda_{n-1}$ are the roots of $t^{n}-\prod_{j=1}^{n}\left(t-a_{j}\right)$.

Reciprocally, let us assume that the roots of $q(t)=t^{n}-\prod_{j=1}^{n}\left(t-a_{j}\right)$ are real and simple. Let $\lambda_{1}<\cdots<\lambda_{n-1}$ be these roots. None of them is zero, since $q(0) \neq 0$. Besides, $\lambda_{k}$ is a root of $q(t)$ if and only if $\beta_{k}:=1 / \lambda_{k} \neq 0$ is a root of

$$
Q(x):=\frac{(-1)^{n-1} x^{n} q(1 / x)}{a_{1} \cdots a_{n}}=P(x)-P(0),
$$

where $P(x)=\prod_{j=1}^{n}\left(x-\alpha_{j}\right)$ with $\alpha_{j}=1 / a_{j}$. Therefore, the roots $\alpha_{j}=1 / a_{j}$ and $\beta_{k}=1 / \lambda_{k}$ are ordered as stated in Lemma 10. The consequences are two-fold. On the one hand, $\lambda_{k} \in\left(a_{S_{k}}, a_{S_{k}+1}\right)$, where $\varsigma=\left(\varsigma_{1}, \ldots, \varsigma_{n-1}\right)$ is the caustic type given in (10). On the other hand, there exist nonsingular billiard trajectories inside the ellipsoid $Q$ sharing the caustics $Q_{\lambda_{1}}, \ldots, Q_{\lambda_{n-1}}$, because the existence conditions (2) hold. Thus, the trajectories sharing the caustics $Q_{\lambda_{1}}, \ldots, Q_{\lambda_{n-1}}$ are periodic with elliptic period $n$ and caustic type $\varsigma$, since the generalized Cayley condition $C(n, n)$ holds; see Proposition 12.

Next, we prove the claims on the (Cartesian) period and the winding numbers. The caustic parameters are located in the intervals delimited by the ellipsoidal parameters given at the beginning of the proof, which implies that

$$
\left\{c_{2 j}, c_{2 j+1}\right\} \cap\left\{a_{1}, \ldots, a_{n}\right\} \neq \emptyset, \quad j=0, \ldots, n-1,
$$

where $c_{0}:=0<c_{1}<\cdots<c_{2 n-1}$ are defined in (3). Thus, all winding numbers are even -see Theorem 5- and so, by definition of elliptic period, $m_{0}=2 m=2 n$.

Finally, let us assume that winding numbers are ordered as stated in Conjecture 1 , so $2 \leq$ $m_{n-1}<\cdots<m_{0}=2 n$ with $m_{0}, \ldots, m_{n-1}$ even. Then $m_{j}=2 \widetilde{m}_{j}=2(n-j)$.

Remark 1. If all the roots of $t^{n}-\prod_{j=1}^{n}\left(t-a_{j}\right)$ are real, but some of them are double, then we get singular periodic billiard trajectories. In that case, there are only two possible scenarios. Either $n$ is odd and $\lambda_{2 l-1}=\lambda_{2 l}$ for some $l=1, \ldots,(n-1) / 2$; or $n$ is even and $\lambda_{2 l}=\lambda_{2 l+1}$ for some $l=1, \ldots,(n-2) / 2$. In all these cases, the singular periodic trajectories are formed by segments contained in some nonsingular ruled confocal quadrics.

Remark 2. All the periodic billiard trajectories mentioned in Theorem 13 have caustic type (10). One may establish similar theorems for other caustic types. For instance, the versions EH1, H1H1, EH2, and H1H2 of Theorem 13 in the spatial case will be listed in Table 2.

## 6. Cayley conditions in the general case

Once we have understood the minimal case $m=n$, we tackle the general case $m \geq n$.
Let us explain the fundamental question by means of an example. In Section 3 we saw that condition $C(3,2)$ becomes a single homogeneous symmetric polynomial equation of degree eight in the variables $\gamma_{1}, \gamma_{2}, \gamma_{3}$, with only two feasible solutions; namely, the ones given in (6). Thus, it is natural to ask whether can we rewrite $\mathcal{C}(3,2)$ as a set of two simpler conditions such that each one of them gives rise to one of the solutions given in (6).

By the way, we raise this question for any $m \geq n$. Can we rewrite $C(m, n)$ as a set of "simpler" conditions such that each one of them gives rise to just "one" solution of $C(m, n)$ ? We answer this question in the affirmative. Indeed, we parametrize these "simpler" conditions by the elements of the set

$$
\mathcal{T}(m, n)=\left\{\left(\tau_{1}, \ldots, \tau_{n}\right) \in \mathbb{Z}^{n}: \tau_{1}+\cdots+\tau_{n}=m-n, \tau_{1}, \ldots, \tau_{n} \geq 0\right\} .
$$

The cardinal of $\mathcal{T}(m, n)$ is the number of monomials of degree $m-n$ in $n$ variables. Thus, $\# \mathcal{T}(m, n)=\binom{m-1}{n-1}$, which gives a precise estimate of the complexity of the Cayley condition $\mathcal{C}(m, n)$ when $m$ grows. We will refer to the elements of $\mathcal{T}(m, n)$ as signatures. We set $\gamma_{2 n}=0$ in order to simplify some notations.

Definition 5. Given any signature $\tau=\left(\tau_{1}, \ldots, \tau_{n}\right) \in \mathcal{T}(m, n)$, we say that condition $C(m, n ; \tau)$ holds if and only if there exist two monic polynomials $S(x), P(x) \in \mathbb{R}[x]$ such that $\operatorname{deg}[S]=m-n$, $\operatorname{deg}[P]=m, P(0) \neq 0, S^{2}(x) R(x)=P(x)(P(x)-P(0))$, and $S(x)$ has $m-n$ simple real roots $\delta_{m-n}<\cdots<\delta_{1}$ such that

$$
\begin{equation*}
\#\left(\left\{\delta_{1}, \ldots, \delta_{m-n}\right\} \cap\left(\gamma_{2 r}, \gamma_{2 r-1}\right)\right)=\tau_{r}, \quad r=1, \ldots, n . \tag{12}
\end{equation*}
$$

Corollary 14. If there exists $\tau \in \mathcal{T}(m, n)$ such that $C(m, n ; \tau)$ holds, then $\mathcal{C}(m, n)$ also holds. The reciprocal implication is true for $m \leq n+3$ (or provided Conjecture 2 holds).

Proof. The first implication is obvious. For the reciprocal implication, we simply recall that $S(x)$ has only real roots when $m \leq n+3$ and all its real roots are contained in $\{x \in \mathbb{R}: R(x)<0\}=$ $\bigcup_{r=1}^{n}\left(\gamma_{2 r}, \gamma_{2 r-1}\right)$; see Proposition 9.

Definition 6. Given any signature $\tau \in \mathcal{T}(m, n)$, let $\{1, \ldots, 2 n-1\}=J_{\tau} \cup K_{\tau}$ and $\{1, \ldots, m-n\}=$ $V_{\tau} \cup W_{\tau}$ be the decompositions determined as follows. If $\delta_{m-n}<\cdots<\delta_{1}$ is any ordered sequence verifying (12), then the elements of the multisets

$$
\begin{aligned}
\left\{\alpha_{1}, \ldots, \alpha_{m}\right\} & =\left\{\gamma_{j}: j \in J_{\tau}\right\} \cup\left\{\delta_{v}, \delta_{v}: v \in V_{\tau}\right\}, \\
\left\{\beta_{1}, \ldots, \beta_{m-1}\right\} & =\left\{\gamma_{k}: k \in K_{\tau}\right\} \cup\left\{\delta_{w}, \delta_{w}: w \in W_{\tau}\right\},
\end{aligned}
$$

are ordered as in Lemma 10.
Multisets are a generalization of sets in which members are allowed to appear more than once; see [27]. In our case, the numbers $\delta_{1}, \ldots, \delta_{m-n}$ appear twice.

These decompositions are well-defined. That is, they only depend on the signature $\tau$, since any ordered sequence $\delta_{m-n}<\cdots<\delta_{1}$ verifying (12) gives rise to the same decomposition. The decomposition $\{1, \ldots, 2 n-1\}=J_{n} \cup K_{n}$ given in Corollary 11 correspond to the trivial signature $\tau=(0, \ldots, 0) \in \mathcal{T}(n, n)$.

Next, we generalize Corollary 11 and Proposition 12 to the case $m \geq n$.

Corollary 15. Let $\delta_{m-n}<\cdots<\delta_{1}$ be an ordered sequence verifying (12) for some signature $\tau \in \mathcal{T}(m, n)$. If $P(x)$ is a monic polynomial of degree $m$ such that $P(0) \neq 0$ and

$$
\prod_{u=1}^{m-n}\left(x-\delta_{u}\right)^{2} \cdot R(x)=P(x)(P(x)-P(0))
$$

then $P(x)=\prod_{j \in J_{T}}\left(x-\gamma_{j}\right) \prod_{v \in V_{T}}\left(x-\delta_{v}\right)^{2}=P(0)+x \prod_{k \in K_{T}}\left(x-\gamma_{k}\right) \prod_{w \in W_{\tau}}\left(x-\delta_{w}\right)^{2}$.
Proof. There exist two decompositions $\{1, \ldots, 2 n-1\}=J^{\prime} \cup K^{\prime}$ and $\{1, \ldots, m-n\}=V^{\prime} \cup W^{\prime}$ such that $P(x)=\prod_{j \in J^{\prime}}\left(x-\gamma_{j}\right) \prod_{v \in V^{\prime}}\left(x-\delta_{v}\right)^{2}=P(0)+x \prod_{k \in K^{\prime}}\left(x-\gamma_{k}\right) \prod_{w \in W^{\prime}}\left(x-\delta_{w}\right)^{2}$. The polynomial $P(x)$ verifies the hypotheses stated in Lemma 10, so the roots

$$
\begin{aligned}
\left\{\alpha_{1}, \ldots, \alpha_{m}\right\} & =\left\{\gamma_{j}: j \in J^{\prime}\right\} \cup\left\{\delta_{v}, \delta_{v}: v \in V^{\prime}\right\} \\
\left\{\beta_{1}, \ldots, \beta_{m-1}\right\} & =\left\{\gamma_{k}: k \in K^{\prime}\right\} \cup\left\{\delta_{w}, \delta_{w}: w \in W^{\prime}\right\}
\end{aligned}
$$

obey the ordering described in that lemma. Hence, $J^{\prime}=J_{\tau}, K^{\prime}=K_{\tau}, V^{\prime}=V_{\tau}$, and $W^{\prime}=W_{\tau}$.
Proposition 16. Condition $C(m, n ; \tau)$ holds if and only if there exists a sequence $\delta_{m-n}<\cdots<\delta_{1}$ verifying (12) such that the following three equivalent properties hold:
(i) If $P(x)=\prod_{j \in J_{T}}\left(x-\gamma_{j}\right) \prod_{v \in V_{T}}\left(x-\delta_{v}\right)^{2}$, then $P(x)-P(0)=x \prod_{k \in K_{T}}\left(x-\gamma_{k}\right) \prod_{w \in W_{T}}\left(x-\delta_{w}\right)^{2}$.
(ii) $\mathrm{e}_{l}\left(\left\{\gamma_{j}\right\}_{j \in J_{\tau}} \cup\left\{\delta_{v}, \delta_{v}\right\}_{v \in V_{\tau}}\right)=\mathrm{e}_{l}\left(\left\{\gamma_{k}\right\}_{k \in K_{\tau}} \cup\left\{\delta_{w}, \delta_{w}\right\}_{w \in W_{\tau}}\right)$, for all $l=1, \ldots, m-1$.
(iii) $\sum_{j \in J_{\tau}} \gamma_{j}^{l}+2 \sum_{v \in V_{\tau}} \delta_{v}^{l}=\sum_{k \in K_{\tau}} \gamma_{k}^{l}+2 \sum_{w \in W_{\tau}} \delta_{w}^{l}$, for all $l=1, \ldots, m-1$.

Proof. We simply repeat the steps of the proof of Proposition 12, but using Corollary 15 instead of Corollary 11.

Example 3. The quantities $\gamma_{3}=1, \gamma_{2}=4$, and $\gamma_{1}=9$ verify condition $C(3,2 ; \tau)$ with $\tau=(1,0)$, because $1+4+9=2 \cdot 7,1^{2}+4^{2}+9^{2}=2 \cdot 7^{2}$, and $7 \in(4,9)$. Hence, the billiard trajectories

- Inside the ellipse $Q: x^{2}+4 y^{2}=1$ with caustic parameter $\lambda=1 / 9$; or
- Inside the ellipse $Q: x^{2}+9 y^{2}=1$ with caustic parameter $\lambda=1 / 4$;
are periodic with elliptic periodic $m=3$. Their caustic types are $\varsigma=0$ and $\varsigma=1$, respectively.
All conditions $C(m, n ; \tau), \tau \in \mathcal{T}(m, n)$, give rise to nonsingular periodic billiard trajectories with elliptic period $m$, so we wondered which is the dynamical meaning of the signature $\tau$. We believe that there exists a one-to-one correspondence between the elliptic winding numbers $\widetilde{m}_{0}, \ldots, \widetilde{m}_{n-1}$-see Definition 3- and the signature $\tau=\left(\tau_{1}, \ldots, \tau_{n}\right)$.

Conjecture 3. Set $\widetilde{m}_{n}=0$. Then $\widetilde{m}_{j}=\widetilde{m}_{j+1}+\tau_{j+1}+1$ for all $j=0, \ldots, n-1$.
This conjecture follows from the interpretation of $C(m, n ; \tau)$ as a singular limit of $C(m, m)$ when $m-n$ couples of simple roots collide, so they become double roots. Unfortunately, we have not been able to transform this argument into a rigorous proof, although all our analytical and numerical computations agree with the conjecture.

To end this section, we stress that if conjectures 2 and 3 hold, then the elliptic winding numbers $\widetilde{m}_{0}, \ldots, \widetilde{m}_{n-1}$ of any nonsingular periodic billiard trajectory verify the above-mentioned relations for some signature $\tau=\left(\tau_{1}, \ldots, \tau_{n}\right)$ with non-negative entries, so the sequence $\widetilde{m}_{0}, \ldots, \widetilde{m}_{n-1}$ strictly decreases, and Conjecture 1 holds.

## 7. The planar case

We adapt the previous setting of billiards inside ellipsoids of $\mathbb{R}^{n}$ to the planar case $n=2$. To follow traditional conventions in the literature, we write the ellipse as

$$
\begin{equation*}
Q=\left\{(x, y) \in \mathbb{R}^{2}: \frac{x^{2}}{a}+\frac{y^{2}}{b}=1\right\}, \quad a>b>0 \tag{13}
\end{equation*}
$$

Any nonsingular billiard trajectory inside $Q$ is tangent to one confocal caustic

$$
Q_{\lambda}=\left\{(x, y) \in \mathbb{R}^{2}: \frac{x^{2}}{a-\lambda}+\frac{y^{2}}{b-\lambda}=1\right\},
$$

where $\lambda \in \Lambda=E \cup H$, with $E=(0, b)$ and $H=(b, a)$.
The names of the connected components of $\Lambda$ come from the fact that $Q_{\lambda}$ is a confocal ellipse for $\lambda \in E$ and a confocal hyperbola for $\lambda \in H$. The singular cases $\lambda=b$ and $\lambda=a$ correspond to the $x$-axis and $y$-axis, respectively. We say that the caustic type of a billiard trajectory is E or H , when its caustic is an ellipse or a hyperbola (compare with Definition 1). We also distinguish between E-caustics and H-caustics.

We recall some concepts related to periodic trajectories of billiards inside ellipses. These results can be found, for instance, in [9, 24]. To begin with, we introduce the function $\rho: \Lambda \rightarrow \mathbb{R}$ given by the quotient of elliptic integrals

$$
\begin{equation*}
\rho(\lambda)=\rho(\lambda ; b, a):=\frac{\int_{0}^{\min (b, \lambda)} \frac{\mathrm{d} t}{\sqrt{(\lambda-t)(b-t)(a-t)}}}{2 \int_{\max (b, \lambda)}^{a} \frac{\mathrm{~d} t}{\sqrt{(\lambda-t)(b-t)(a-t)}}} . \tag{14}
\end{equation*}
$$

It is called the rotation number and characterizes the caustic parameters that give rise to periodic trajectories. To be precise, the billiard trajectories with caustic $Q_{\lambda}$ are periodic if and only if

$$
\rho(\lambda)=m_{1} / 2 m_{0} \in \mathbb{Q}
$$

for some integers $2 \leq m_{1}<m_{0}$, which are the winding numbers. On the one hand, $m_{0}$ is the period, On the other hand, $m_{1}$ is twice the number of turns around the ellipse $Q_{\lambda}$ for E-caustics, and the number of crossings of the $y$-axis for H-caustics. Thus, $m_{1}$ is always even. Besides, all periodic trajectories with H -caustics have even period. (Compare with Theorem 5.)

Proposition 17. The winding numbers $2 \leq m_{1}<m_{0}$, rotation number $\rho=m_{1} / 2 m_{0}$, signature $\tau=\left(\tau_{1}, \tau_{2}\right) \in \mathcal{T}(m, 2)$, caustic type ( $E$ or $H$ ), and caustic parameter $\lambda$ of all nonsingular periodic billiard trajectories inside the ellipse (13) with elliptic period $m \in\{2,3\}$ are listed in Table 1. The ellipses where such trajectories take place are also listed.

Proof. We split the proof in four steps.
Step 1: To find the solutions of $C(m, 2)$ in terms of the inverse quantities $\gamma_{i}$. First, we saw in Proposition 12 that $C(2,2)$ holds if and only if $\gamma_{1}=\gamma_{2}+\gamma_{3}$.

Next, we focus on the case $m=3$. We note that $C(3,2)$ holds if and only if $C(3,2 ; \tau)$ holds for some $\tau=\left(\tau_{1}, \tau_{2}\right) \in \mathbb{Z}^{2}$ such that $\tau_{1}+\tau_{2}=1$ and $\tau_{1}, \tau_{2} \geq 0$; see Corollary 14 .

| $m$ | $m_{0}$ | $m_{1}$ | $\rho$ | $\tau$ | Type | Ellipses | Caustic parameter |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2 | 4 | 2 | $1 / 4$ | $(0,0)$ | E | any | $\frac{a b}{a+b}$ |
| 2 | 4 | 2 | $1 / 4$ | $(0,0)$ | H | $2 b<a$ | $\frac{a b}{a b b}$ |
| 3 | 6 | 2 | $1 / 6$ | $(1,0)$ | E | any | $\frac{a b}{a+b+2 \sqrt{a b}}$ |
| 3 | 6 | 2 | $1 / 6$ | $(1,0)$ | H | $4 b<a$ | $\frac{a b}{a+b-2 \sqrt{a b}}$ |
| 3 | 3 | 2 | $1 / 3$ | $(0,1)$ | E | any | $\frac{a b}{a+b+2 \sqrt{a^{2}-a b+b^{2}}}$ |
| 3 | 6 | 4 | $1 / 3$ | $(0,1)$ | H | $4 b<3 a$ | $\frac{a b}{2 \sqrt{a^{2}-a b}+b-a}$ |

Table 1: Algebraic formulas for the caustic parameter corresponding to nonsingular periodic billiard trajectories with elliptic period $m \in\{2,3\}$ in the planar case.

Let us begin with the signature $\tau=(1,0)$. After a straightforward check, we get that the decompositions presented in Definition 6 are $J_{\tau}=\{1,2,3\}, K_{\tau}=V_{\tau}=\emptyset$, and $W_{\tau}=\{1\}$. Thus, $\mathcal{C}(3,2 ; \tau)$ holds if and only if there exists some $\delta_{1} \in\left(\gamma_{2}, \gamma_{1}\right)$ such that

$$
P(x)=\left(x-\gamma_{1}\right)\left(x-\gamma_{2}\right)\left(x-\gamma_{3}\right)=P(0)+x\left(x-\delta_{1}\right)^{2}
$$

or, equivalently, if and only if the discriminant of the polynomial

$$
Q(x)=\frac{P(x)-P(0)}{x}=x^{2}-\mathrm{e}_{1}\left(\gamma_{1}, \gamma_{2}, \gamma_{3}\right) x+\mathrm{e}_{2}\left(\gamma_{1}, \gamma_{2}, \gamma_{3}\right)
$$

is equal to zero. The discriminant of $Q(x)$ is

$$
\Delta=\gamma_{1}^{2}+\gamma_{2}^{2}+\gamma_{3}^{2}-2 \gamma_{1} \gamma_{2}-2 \gamma_{1} \gamma_{3}-2 \gamma_{2} \gamma_{3} .
$$

We already saw in Section 3 that the only feasible solution of $\Delta=0$ is $\sqrt{\gamma_{1}}=\sqrt{\gamma_{2}}+\sqrt{\gamma_{3}}$.
When $\tau=(0,1)$, the decompositions are $J_{\tau}=\{1\}, K_{\tau}=\{2,3\}, V_{\tau}=\{1\}$, and $W_{\tau}=\emptyset$. Thus, $\mathcal{C}(3,2 ; \tau)$ holds if and only if there exists some $\delta_{1} \in\left(0, \gamma_{3}\right)$ such that

$$
\gamma_{1}+2 \delta_{1}=\gamma_{2}+\gamma_{3}, \quad \gamma_{1}^{2}+2 \delta_{1}^{2}=\gamma_{2}^{2}+\gamma_{3}^{2}
$$

or, equivalently, if and only if

$$
\begin{aligned}
3 \gamma_{1}^{2}-2\left(\gamma_{2}+\gamma_{3}\right) \gamma_{1}-\left(\gamma_{2}-\gamma_{3}\right)^{2} & =\left(\gamma_{2}+\gamma_{3}-\gamma_{1}\right)^{2}-2\left(\gamma_{2}^{2}+\gamma_{3}^{2}-\gamma_{1}^{2}\right) \\
& =\left(2 \delta_{1}\right)^{2}-4 \delta_{1}^{2}=0
\end{aligned}
$$

And we already saw in Section 3 that the only feasible solution of the above equation is

$$
3 \gamma_{1}=\gamma_{2}+\gamma_{3}+2 \sqrt{\gamma_{2}^{2}+\gamma_{3}^{2}-\gamma_{2} \gamma_{3}}
$$

Step 2: To express the above solutions in terms of $a, b$, and $\lambda$. If the caustic type is E , then $\lambda \in(0, b), \gamma_{1}=1 / \lambda, \gamma_{2}=1 / b$, and $\gamma_{3}=1 / a$. Thus,

$$
\begin{aligned}
\gamma_{1}=\gamma_{2}+\gamma_{3} & \Leftrightarrow \lambda=\frac{a b}{a+b}, \\
\sqrt{\gamma_{1}}=\sqrt{\gamma_{2}}+\sqrt{\gamma_{3}} & \Leftrightarrow \lambda=\frac{a b}{a+b+2 \sqrt{a b}} \\
3 \gamma_{1}=\gamma_{2}+\gamma_{3}+2 \sqrt{\gamma_{2}^{2}+\gamma_{3}^{2}-\gamma_{2} \gamma_{3}} & \Leftrightarrow \lambda=\frac{3 a b}{a+b+2 \sqrt{a^{2}-a b+b^{2}}} .
\end{aligned}
$$


(a) $\lambda=\frac{a b}{a+b}$

(d) $\lambda=\frac{a b}{a-b}$

(b) $\lambda=\frac{a b}{a+b+2 \sqrt{a b}}$

(e) $\lambda=\frac{a b}{a+b-2 \sqrt{a b}}$

(c) $\lambda=\frac{3 a b}{a+b+2 \sqrt{a^{2}-a b+b^{2}}}$

(f) $\lambda=\frac{a b}{2 \sqrt{a^{2}-a b}+b-a}$

Figure 1: Periodic trajectories corresponding to the caustic parameters listed in Table 1. The ellipse for $\lambda=\frac{a b}{a+b-2 \sqrt{a b}}$ is flatter, because it must satisfy condition $4 b<a$.

If the caustic type is H , then $\lambda \in(b, a), \gamma_{1}=1 / b, \gamma_{2}=1 / \lambda$, and $\gamma_{3}=1 / a$. Thus,

$$
\begin{aligned}
\gamma_{1}=\gamma_{2}+\gamma_{3} & \Leftrightarrow \lambda=\frac{a b}{a-b}, \\
\sqrt{\gamma_{1}}=\sqrt{\gamma_{2}}+\sqrt{\gamma_{3}} & \Leftrightarrow \lambda=\frac{a b}{a+b-2 \sqrt{a b}} \\
3 \gamma_{1}=\gamma_{2}+\gamma_{3}+2 \sqrt{\gamma_{2}^{2}+\gamma_{3}^{2}-\gamma_{2} \gamma_{3}} & \Leftrightarrow \lambda=\frac{a b}{2 \sqrt{a^{2}-a b}+b-a}
\end{aligned}
$$

Step 3: To determine the ellipses where such periodic billiard trajectories take place. We ask whether the caustic parameters found above belong to the interval $(0, b)$ for E-caustics, and to the interval $(b, a)$ for H -caustics. The caustic type E does not give any restriction, because

$$
0<b<a \text { and } \lambda \in\left\{\frac{a b}{a+b}, \frac{a b}{a+b+2 \sqrt{a b}}, \frac{3 a b}{a+b+2 \sqrt{a^{2}-a b+b^{2}}}\right\} \Rightarrow \lambda \in(0, b) .
$$

On the contrary, the caustic type H gives rise to some restrictions. Namely,

$$
\begin{aligned}
& b<\frac{a b}{a-b}<a \Leftrightarrow \quad 2 b<a \\
& b<\frac{a b}{a+b-2 \sqrt{a b}}<a \quad \Leftrightarrow \quad 4 b<a \\
& b<\frac{a b}{2 \sqrt{a^{2}-a b}+b-a}<a \quad \Leftrightarrow \quad 4 b<3 a
\end{aligned}
$$

Step 4: To find the winding numbers and the rotation number. The winding numbers $2 \leq$ $m_{1}<m_{0}$ and the rotation number $\rho(\lambda)=m_{1} / 2 m_{0}$ are obtained from geometric arguments. To be precise, we draw in Figure 1 a billiard trajectory tangent to $Q_{\lambda}$ for each of the caustic parameters listed in Table 1. Then we recall that $m_{0}$ is the period and $m_{1}$ is twice the number of turns around the ellipse $Q_{\lambda}$ for E-caustics, and the number of crossings of the $y$-axis for H -caustics.

Let $\rho_{*} \in\{1 / 3,1 / 4,1 / 6\}$. We have seen above that inside any ellipse (13) there exists a unique E-caustic whose tangent billiard trajectories have rotation number $\rho_{*}$. Besides, if $\lambda_{\mathrm{E}}\left(a, b ; \rho_{*}\right)$ and $\lambda_{\mathrm{H}}\left(a, b ; \rho_{*}\right)$ denote the caustic parameters associated to the E-caustic and H-caustic with rotation number $\rho_{*}$, we see that

$$
\begin{equation*}
\lambda_{\mathrm{E}}\left(b, a ; \rho_{*}\right)=\lambda_{\mathrm{E}}\left(a, b ; \rho_{*}\right), \quad b=\lambda_{\mathrm{E}}\left(a, \lambda_{\mathrm{H}}\left(a, b ; \rho_{*}\right) ; \rho_{*}\right) . \tag{15}
\end{equation*}
$$

These properties can be generalized. On the one hand, the rotation number (14) is an increasing function in the interval $(0, b)$ such that $\rho(0)=0$ and $\rho(b)=1 / 2$; see [24]. This means that given any $\rho_{*} \in(0,1 / 2)$, there exists a unique $\lambda_{*} \in(0, b)$ such that $\rho\left(\lambda_{*}\right)=\rho_{*}$, and so, there exists a unique E-caustic whose tangent billiard trajectories have rotation number $\rho_{*}$. On the other hand, relations (15) can be obtained by using that the rotation number (14) is symmetric in the three parameters $a, b$, and $\lambda$. Consequently, one can find the formula for $\lambda_{\mathrm{H}}$ (respectively, $\lambda_{\mathrm{E}}$ ) from the formula for $\lambda_{\mathrm{E}}$ (respectively, $\lambda_{\mathrm{H}}$ ) by using the second relation.

It is interesting to realize that the results in Table 1 agree with Conjecture 3.
In the planar case $n=2$, the caustic type (10) is $\varsigma=0$ or, equivalently, E. Hence, the planar version of Theorem 13 is shown in the first row of Table 1, because $\lambda=\lambda_{\mathrm{E}}(a, b ; 1 / 4)=a b /(a+b)$ is the root of $t^{2}-(t-a)(t-b)$. This naive observation was the germ of this paper.

## 8. The spatial case

In order to study the spatial case $n=3$, we consider the triaxial ellipsoid

$$
\begin{equation*}
Q=\left\{(x, y, z) \in \mathbb{R}^{3}: \frac{x^{2}}{a}+\frac{y^{2}}{b}+\frac{z^{2}}{c}=1\right\}, \quad a>b>c>0 . \tag{16}
\end{equation*}
$$

Any nonsingular billiard trajectory inside $Q$ is tangent to two distinct nonsingular caustics $Q_{\lambda_{1}}$ and $Q_{\lambda_{2}}$, with $\lambda_{1}<\lambda_{2}$, of the confocal family

$$
\begin{equation*}
Q_{\lambda}=\left\{(x, y, z) \in \mathbb{R}^{3}: \frac{x^{2}}{a-\lambda}+\frac{y^{2}}{b-\lambda}+\frac{z^{2}}{c-\lambda}=1\right\} \tag{17}
\end{equation*}
$$

The caustic $Q_{\lambda}$ is an ellipsoid for $\lambda \in(0, c)$, a hyperboloid of one sheet when $\lambda \in(c, b)$, and a hyperboloid of two sheets if $\lambda \in(b, a)$. Not all combinations of nonsingular caustics can take place, but only the four caustic types EH1, H1H1, EH2, and H1H2.

Proposition 18. The caustic type and caustic parameters of all nonsingular periodic billiard trajectories inside the triaxial ellipsoid (16) with elliptic period $m=3$ are listed in Table 2. The ellipsoids where such trajectories take place are also listed.

Proof. If $a, b$, and $c$ are the ellipsoidal parameters, and $\lambda_{1}$ and $\lambda_{2}$ are the caustic parameters, we set $\left\{c_{1}, c_{2}, c_{3}, c_{4}, c_{5}\right\}=\left\{a, b, c, \lambda_{1}, \lambda_{2}\right\}$, where $0<c_{1}<c_{2}<c_{3}<c_{4}<c_{5}$. We also set $\gamma_{i}=1 / c_{i}$. Let $\{1,2,3,4,5\}=J \cup K$, with $J=\{1,4,5\}$ and $K=\{2,3\}$, be the decomposition defined in Corollary 11 when $n=3$. From Proposition 12 we know that

$$
\begin{aligned}
C(3,3) & \Leftrightarrow \gamma_{2}+\gamma_{3}=\gamma_{1}+\gamma_{4}+\gamma_{5} \text { and } \gamma_{2}^{2}+\gamma_{3}^{2}=\gamma_{1}^{2}+\gamma_{4}^{2}+\gamma_{5}^{2} \\
& \Leftrightarrow\left(\gamma_{1}-\gamma_{k}\right)\left(\gamma_{4}-\gamma_{k}\right)\left(\gamma_{5}-\gamma_{k}\right)=\gamma_{1} \gamma_{4} \gamma_{5}, \text { for } k=2,3 \\
& \Leftrightarrow c_{k}^{3}=\left(c_{k}-c_{1}\right)\left(c_{4}-c_{k}\right)\left(c_{5}-c_{k}\right), \text { for } k=2,3 .
\end{aligned}
$$

\(\left.$$
\begin{array}{lll}\hline \text { Type } & \text { Ellipsoids } & \text { Caustic parameters } \\
\hline \text { EH1 } & c<\frac{a b}{a+b+\sqrt{a b}} & \left\{\begin{array}{c}c^{3}=\left(c-\lambda_{1}\right)(b-c)(a-c) \\
1 / \lambda_{2}+1 / c=1 / a+1 / b+1 / \lambda_{1}\end{array}
$$\right. <br>

H1H1 \& c<\frac{a b}{a+b+2 \sqrt{a b}} \& Roots of t^{3}-(t-a)(t-b)(t-c)\end{array}\right\}\)| EH2 $\left\{\begin{array}{l}c<\frac{a-2 b}{2 a-3 b} a \\ 2 b<a \\ c<\frac{a-2 b}{(a-b)^{2}} a b \\ b>\frac{a c}{a+c-\sqrt{a c}}\end{array}\right.$ |
| :--- |
| H1H2 $\left\{\begin{array}{c}b^{3}=(b-c)\left(\lambda_{2}-b\right)(a-b) \\ 1 / \lambda_{1}+1 / b=1 / a+1 / c+1 / \lambda_{2}\end{array}\right.$ |

Table 2: Algebraic formulas for the caustic parameters corresponding to the nonsingular periodic trajectories with elliptic period $m=3$ in the spatial case.

In the rest of the proof, we study each caustic type separately.
Caustic type EH1. In this case $0<\lambda_{1}<c<\lambda_{2}<b<a$, so

$$
c_{1}=\lambda_{1}, \quad c_{2}=c, \quad c_{3}=\lambda_{2}, \quad c_{4}=b, \quad c_{5}=a .
$$

Thus the formula for $\lambda_{1}$ follows from relation $c_{2}^{3}=\left(c_{2}-c_{1}\right)\left(c_{4}-c_{2}\right)\left(c_{5}-c_{2}\right)$, whereas the formula for $\lambda_{2}$ follows from relation $\gamma_{2}+\gamma_{3}=\gamma_{1}+\gamma_{4}+\gamma_{5}$. Next, we look for ellipsoidal parameters such that the caustic parameters computed using these two formulas are placed in the right intervals: $\lambda_{1} \in(0, c)$ and $\lambda_{2} \in(c, b)$.

To begin with, we note that $\lambda_{1}<c$, since $\left(c-\lambda_{1}\right)(b-c)(a-c)=c^{3}>0$. Besides,

$$
\lambda_{1}=\frac{a b-(a+b) c}{(b-c)(a-c)} c>0 \Leftrightarrow c<\frac{a b}{a+b} .
$$

On the other hand, if $\lambda_{1} \in(0, c)$, then

$$
1 / \lambda_{2}=1 / a+1 / b+\left(1 / \lambda_{1}-1 / c\right)>1 / a+1 / b>1 / b
$$

so $\lambda_{2}<b$. Finally,

$$
\begin{aligned}
\lambda_{2}>c & \Leftrightarrow \frac{1}{c}+\frac{c}{a b-(a+b) c}=\frac{1}{\lambda_{1}}=\frac{1}{\lambda_{2}}+\frac{1}{c}-\frac{1}{a}-\frac{1}{b}<\frac{2}{c}-\frac{1}{a}-\frac{1}{b} \\
& \Leftrightarrow c<\frac{a b}{a+b+\sqrt{a b}} .
\end{aligned}
$$

Therefore, $\lambda_{1} \in(0, c)$ and $\lambda_{2} \in(c, b)$ if and only if $c<a b /(a+b+\sqrt{a b})$.
Caustic type H 1 H 1 . If $n=3$, then the caustic type (10) is $\varsigma=(1,1)$ or, equivalently, H 1 H 1 . Hence, the study for the caustic type H1H1 was already carried out in Theorem 13. Suffice it to note that the polynomial

$$
t^{3}-(t-a)(t-b)(t-c)=(a+b+c) t^{2}-(a b+a c+b c) t+a b c
$$

has two real simple roots if and only if its discriminant

$$
\Delta=(a b+a c+b c)^{2}-4 a b c(a+b+c)=(a-b)^{2} c^{2}-2 a b(a+b) c+a^{2} b^{2}
$$

is positive. This discriminant is a second-degree polynomial in $c$ whose roots are

$$
c_{ \pm}=\frac{a b(a+b) \pm 2 a b \sqrt{a b}}{(a-b)^{2}}=\frac{a b}{a+b \mp 2 \sqrt{a b}} .
$$

We note that $0<c_{-}<b<c_{+}$. Thus, using that $0<c<b<a$, we get $\Delta>0 \Leftrightarrow c<c_{-}$.
Caustic type EH2. In this case $0<\lambda_{1}<c<b<\lambda_{2}<a$, so

$$
c_{1}=\lambda_{1}, \quad c_{2}=c, \quad c_{3}=b, \quad c_{4}=\lambda_{2}, \quad c_{5}=a .
$$

Using relations $\gamma_{2}^{l}+\gamma_{3}^{l}=\gamma_{1}^{l}+\gamma_{4}^{l}+\gamma_{5}^{l}$, with $l=1,2$, we know that

$$
s_{l}:=\frac{1}{\lambda_{1}^{l}}+\frac{1}{\lambda_{2}^{l}}=\frac{1}{c^{l}}+\frac{1}{b^{l}}-\frac{1}{a^{l}}, \quad l=1,2 .
$$

Hence, $1 / \lambda_{1}$ and $1 / \lambda_{2}$ are the roots of the polynomial

$$
\left(x-1 / \lambda_{1}\right)\left(x-1 / \lambda_{2}\right)=x^{2}-s_{1} x+\frac{s_{1}^{2}-s_{2}}{2}=x^{2}+\frac{b c-a(b+c)}{a b c} x+\frac{(a-b)(a-c)}{a^{2} b c} .
$$

Thus, using the change of variables $t=1 / x$, we get that $\lambda_{1}$ and $\lambda_{2}$ are the roots of

$$
Q(t)=(a-b)(a-c) t^{2}+(b c-a(b+c)) a t+a^{2} b c
$$

We look for ellipsoidal parameters such that $Q(t)$ has a root in $(0, c)$ and a root in $(b, a)$. The root in $(0, c)$ always exists, since $Q(0)=a^{2} b c>0$ and $Q(c)=-c^{3}(a-b)<0$. Besides, $Q(b)=-b^{3}(a-c)<0$ and $\lim _{t \rightarrow+\infty} Q(t)=+\infty$, so $Q(t)$ has a root in $(b, a)$ if and only if

$$
Q(a)=a^{2}\left(a^{2}-2 a(b+c)+3 b c\right)>0
$$

or, equivalently, if and only if $c<(a-2 b) a /(2 a-3 b)$ and $2 b<a$. We have used that $0<c<b<a$ in the last equivalence.

Caustic type H1H2. In this case $0<c<\lambda_{1}<b<\lambda_{2}<a$, so

$$
c_{1}=c, \quad c_{2}=\lambda_{1}, \quad c_{3}=b, \quad c_{4}=\lambda_{2}, \quad c_{5}=a .
$$

Thus the formula for $\lambda_{2}$ follows from relation $c_{2}^{3}=\left(c_{2}-c_{1}\right)\left(c_{4}-c_{2}\right)\left(c_{5}-c_{2}\right)$, whereas the formula for $\lambda_{1}$ follows from relation $\gamma_{2}+\gamma_{3}=\gamma_{1}+\gamma_{4}+\gamma_{5}$. Next, we look for conditions on the ellipsoidal parameters such that the caustic parameters computed from the previous formulas are placed in the right intervals: $\lambda_{1} \in(c, b)$ and $\lambda_{2} \in(b, a)$.

To begin with, we note that $\lambda_{2}>b$, because $(b-c)\left(\lambda_{2}-b\right)(a-b)=b^{3}>0$. Besides,

$$
\frac{(a+c) b-a c}{(a-b)(b-c)} b=\lambda_{2}<a \Leftrightarrow c<\frac{a-2 b}{(a-b)^{2}} a b .
$$

On the other hand, using that $0<c<b<a$, we get that

$$
\begin{aligned}
c<\lambda_{1}<b & \Leftrightarrow \frac{2}{b}-\frac{1}{a}-\frac{1}{c}<\frac{1}{\lambda_{2}}=\frac{1}{\lambda_{1}}+\frac{1}{b}-\frac{1}{a}-\frac{1}{c}<\frac{1}{b}-\frac{1}{a} \\
& \Leftrightarrow \frac{2}{b}-\frac{1}{a}-\frac{1}{c}<\frac{1}{b}-\frac{b}{(a+c) b-a c}<\frac{1}{b}-\frac{1}{a} \\
& \Leftrightarrow b>\frac{a c}{a+c-\sqrt{a c}} .
\end{aligned}
$$

Thus, $\lambda_{1} \in(0, c)$ and $\lambda_{2} \in(b, a)$ if and only if $c<(a-2 b) a b /(a-b)^{2}$ and $b>a c /(a+c-\sqrt{a c})$.

Let us look for the winding numbers of the trajectories described in the previous proposition. The winding numbers $m_{0}, m_{1}$, and $m_{2}$ describe how the periodic billiard trajectories fold in $\mathbb{R}^{3}$. The following results can be found in [24, table 1]. First, $m_{0}$ is the period. Second, $m_{1}$ is the number of $x y$-crossings and $m_{2}$ is twice the number of turns around the $z$-axis for EH1-caustics; $m_{1}$ is twice the number of turns around the $x$-axis and $m_{2}$ is the number of $y z$-crossings for EH2-caustics; $m_{1}$ is the number of tangential touches with each hyperboloid of one sheet caustic and $m_{2}$ is twice the number of turns around the $z$-axis for H 1 H 1 -caustics; $m_{1}$ is the number of $x z$-crossings and $m_{2}$ is the number of $y z$-crossings for H 1 H 2 -caustics. Besides, all periodic trajectories with H1H1-caustics or H1H2-caustics have even period. Several periodic billiard trajectories with elliptic period $m=3$ were depicted in [25, Table XV and Table XVII]. We conclude by direct inspection of those pictures that the nonsingular periodic billiard trajectories inside a triaxial ellipsoid with elliptic period $m=3$ have winding numbers

$$
m_{2}=2, \quad m_{1}=4, \quad m_{0}=6
$$

This agrees with the formulas (11) given in Theorem 13. We emphasize that those formulas were not rigorously proved, because their "proof" was based on Conjecture 1.

Next, we establish the algebraic formulas for the caustic parameters of other nonsingular periodic billiard trajectories. We begin with a technical lemma about four-degree polynomials.

Lemma 19. Let $Q(x)=\left(x-\alpha_{-}\right)\left(x-\beta_{-}\right)\left(x-\beta_{+}\right)\left(x-\alpha_{+}\right)$for some $\alpha_{-}<\beta_{-}<\beta_{+}<\alpha_{+}$. Let $\nu_{-} \in\left(\alpha_{-}, \beta_{-}\right), v \in\left(\beta_{-}, \beta_{+}\right)$, and $v_{+} \in\left(\beta_{+}, \alpha_{+}\right)$be the three roots of $Q^{\prime}(x)$. Then

$$
Q\left(v_{+}\right)<Q\left(v_{-}\right) \Leftrightarrow \alpha_{-}+\alpha_{+}>\beta_{-}+\beta_{+} .
$$

Proof. If we set $\eta=\left(\beta_{+}+\beta_{-}\right) / 2$ and $\xi=\left(\beta_{+}-\beta_{-}\right) / 2$, then

$$
Q(\eta+s)-Q(\eta-s)=2\left(s^{2}-\xi^{2}\right)\left(\beta_{-}+\beta_{+}-\alpha_{-}-\alpha_{+}\right) s, \quad \forall s \in \mathbb{R} .
$$

On the one hand, if $\alpha_{-}+\alpha_{+}>\beta_{-}+\beta_{+}$, then $Q(\eta+s)<Q(\eta-s)$ for all $s>\xi$, which implies that $Q\left(v_{+}\right)<Q\left(v_{-}\right)$. On the other hand, if $\alpha_{-}+\alpha_{+}<\beta_{-}+\beta_{+}$, then $Q(\eta+s)>Q(\eta-s)$ for all $s>\xi$, which implies that $Q\left(v_{+}\right)>Q\left(v_{-}\right)$. Finally, if $\alpha_{-}+\alpha_{+}=\beta_{-}+\beta_{+}$, then $Q(\eta+s)=Q(\eta-s)$ for all $s \in \mathbb{R}$, which implies that $Q\left(v_{+}\right)=Q\left(v_{-}\right)$.

We can now answer some questions about the nonsingular periodic billiard trajectories found in the second item of Theorem 8, although the study is restricted to the spatial case.

Proposition 20. There exist periodic billiard trajectories inside the triaxial ellipsoid (16) with elliptic period $m=4$, signature $\tau=(0,0,1)$, and caustic type H1H1 if and only if

$$
c<a b /(a+b)
$$

Besides, the caustic parameters $\lambda_{1}$ and $\lambda_{2}$ of such periodic billiard trajectories are the roots of the quadratic polynomial $\left(s_{1}^{2}-s_{2}\right) t^{2} / 2-s_{1} t+1$, where

$$
\begin{equation*}
s_{l}=1 / a^{l}+1 / b^{l}+1 / c^{l}-2 / d^{l}, \quad l=1,2 \tag{18}
\end{equation*}
$$

and $d$ is the only root of the cubic polynomial $t^{3}-2(a+b+c) t^{2}+3(a b+a c+b c) t-4 a b c$ in the interval $(a,+\infty)$.

Proof. If $\tau=(0,0,1)$, the decompositions presented in Definition 6 are $J_{\tau}=\{2,3\}, K_{\tau}=\{1,4,5\}$, $V_{\tau}=\{1\}$, and $W_{\tau}=\emptyset$. Thus, $C(4,3 ; \tau)$ holds if and only if there exists some $\delta_{1} \in\left(0, \gamma_{5}\right)$ such that the following two equivalent properties hold:
(i) $P(x)=\left(x-\delta_{1}\right)^{2}\left(x-\gamma_{2}\right)\left(x-\gamma_{3}\right) \Rightarrow Q(x):=x\left(x-\gamma_{1}\right)\left(x-\gamma_{4}\right)\left(x-\gamma_{5}\right)=P(x)-P(0)$.
(ii) $\gamma_{2}^{l}+\gamma_{3}^{l}+2 \delta_{1}^{l}=\gamma_{1}^{l}+\gamma_{4}^{l}+\gamma_{5}^{l}$, for $l=1,2,3$.

If the caustic type is H 1 H 1 , then $0<c<\lambda_{1}<\lambda_{2}<b<a$, so

$$
\gamma_{1}=1 / c, \quad \gamma_{2}=1 / \lambda_{1}, \quad \gamma_{3}=1 / \lambda_{2}, \quad \gamma_{4}=1 / b, \quad \gamma_{5}=1 / a .
$$

Let $\mathrm{e}_{l}=\mathrm{e}_{l}\left(\gamma_{1}, \gamma_{4}, \gamma_{5}\right)$ for $l=1,2,3$. We set $d=1 / \delta_{1}>a$. Then $Q(x)=x^{4}-\mathrm{e}_{1} x^{3}+\mathrm{e}_{2} x^{2}-\mathrm{e}_{1} x$ and $\delta_{1}$ is a root of $Q^{\prime}(x)=4 x^{3}-3 \mathrm{e}_{1} x^{2}+2 \mathrm{e}_{2} x-\mathrm{e}_{3}$. Hence, $d$ is a root of the cubic polynomial

$$
q(t)=-a b c t^{3} Q^{\prime}(1 / t)=t^{3}-2(a+b+c) t^{2}+3(a b+a c+b c) t-4 a b c .
$$

We note that $q(0)=-4 a b c<0, q(b)=-b(b-a)(b-c)>0, q(a)=-a(a-b)(a-c)<0$, and $\lim _{t \rightarrow+\infty} q(t)=+\infty$. This shows that $q(t)$ has just one root in the interval $(a,+\infty)$.

From property (ii) above, we deduce that the sums $s_{l}:=1 / \lambda_{1}^{l}+1 / \lambda_{2}^{l}$ verify relations (18). Besides, $1 / \lambda_{1}$ and $1 / \lambda_{2}$ are the roots of $\left(x-1 / \lambda_{1}\right)\left(x-1 / \lambda_{2}\right)=x^{2}-s_{1} x+\left(s_{1}^{2}-s_{2}\right) / 2$, so that $\lambda_{1}$ and $\lambda_{2}$ are the roots of the quadratic polynomial $\left(s_{1}^{2}-s_{2}\right) t^{2} / 2-s_{1} t+1$.

We look for ellipsoidal parameters such that the previous periodic trajectories exist. From property (i) above, we deduce that such ellipsoidal parameters exist if and only the graph $\{y=$ $Q(x)\}$ intersects the horizontal line $\left\{y=Q\left(\delta_{1}\right)\right\}$ at two different points $\gamma_{2}, \gamma_{3} \in\left(\gamma_{4}, \gamma_{5}\right)$ or, equivalently, if and only if $Q\left(\delta_{3}\right)<Q\left(\delta_{1}\right)$, where $\delta_{1}<\delta_{2}<\delta_{3}$ are the three ordered roots of the derivative of the polynomial $Q(x)=x\left(x-\gamma_{1}\right)\left(x-\gamma_{4}\right)\left(x-\gamma_{5}\right)$. But

$$
Q\left(\delta_{3}\right)<Q\left(\delta_{1}\right) \Leftrightarrow \gamma_{1}>\gamma_{4}+\gamma_{5} \Leftrightarrow c<a b /(a+b)
$$

according to Lemma 19.
As we have explained before, the period and winding numbers of any nonsingular periodic billiard trajectory can be determined by direct inspection of its corresponding figure. A periodic billiard trajectory with elliptic period $m=4$ and caustic type H 1 H 1 whose caustic parameters verify the relations given in Proposition 20 is displayed in [24, Figure 13]. That trajectory has (Cartesian) period $m_{0}=4$ and winding numbers

$$
m_{2}=2, \quad m_{1}=3, \quad m_{0}=4
$$

Hence, $\operatorname{gcd}\left(m_{0}, m_{1}, m_{2}\right)=1$, so the elliptic winding numbers are $\widetilde{m}_{2}=2, \widetilde{m}_{1}=3$, and $\widetilde{m}_{0}=4$. This result reinforces Conjecture 3. Besides, these nonsingular four-periodic billiard trajectories with caustic type H 1 H 1 are quite interesting, because they display the minimal period among all nonsingular periodic billiard trajectories; see [24, Theorem 1].

We end the study at this point. We just mention that there exist similar results when the signature or the caustic type do not coincide with the ones given in Proposition 20. Analogously, the case $m=5$ can be dealt with using the same techniques, although the final formulas become more complicated. For instance, it can be easily checked that the caustic parameters $\lambda_{1}$ and $\lambda_{2}$ of the billiard trajectories with elliptic period $m=5$ and signature $\tau=(1,1,0)$ verify the homogeneous symmetric polynomial equations

$$
8 s_{3}+s_{1}^{3}=6 s_{1} s_{2}, \quad 16 s_{4}+s_{1}^{4}=4 s_{1}^{2} s_{2}+4 s_{2}^{2}
$$

where $s_{l}=1 / a^{l}+1 / b^{l}+1 / c^{l}+1 / \lambda_{2}^{l}+1 / \lambda_{1}^{l}$ for $l=1,2,3$, 4 . Each of the other signatures $\tau \in \mathcal{T}(5,3)$ gives rise to similar homogeneous -although not symmetric- polynomial equations of degrees three and four in the variables $1 / a, 1 / b, 1 / c, 1 / \lambda_{1}$ and $1 / \lambda_{2}$. We left the details to the reader. Finally, we remember that the original matrix formulation of the generalized Cayley condition $C(5,3)$ gives rise to two homogeneous symmetric polynomial equations of degrees 23 and 24 in those five variables, as explained in Section 3. This confirms, once again, that the polynomial formulation offers great computational advantages over the matrix formulation.

## References

[1] V. Kozlov V and D. Treshchëv, Billiards: a Genetic Introduction to the Dynamics of Systems with Impacts, Transl. Math. Monographs, vol. 89, 1991.
[2] S. Tabachnikov, Billiards, Panoramas et Syntheses, vol. 1, Société Mathématique de France, 1995.
[3] S. Tabachnikov, Geometry and Billiards, Stud. Math. Libr., vol. 30, AMS, Providence, RI, 2005.
[4] M. Berger, Seules les quadriques admettent des caustiques, Bull. Soc. Math. France, 123 (1995) 107-116.
[5] P. M. Gruber, Only ellipsoids have caustics, Math. Ann., 303 (1995) 185-194.
[6] J. V. Poncelet, Traité des Propriétés Projectives des Figures, Mett-Paris, 1822.
[7] Ph. Griffiths and J. Harris, A Poncelet theorem in space, Comment. Math. Helvetici., 52 (1977) 145-160.
[8] G. Darboux, Leçons sur la théorie générale des surfaces et les applications géométriques du calcul infinitesimal, vol. 2 and 3, Gauthier-Villars, 1887, 1889.
[9] S.-J. Chang and R. Friedberg, Elliptical billiards and Poncelet's theorem, J. Math. Phys., 29 (1988) 1537-1550.
[10] S.-J. Chang, B. Crespi and K.-J. Shi, Elliptical billiard systems and the full Poncelet's theorem in $n$ dimensions, J. Math. Phys., 34 (1993) 2242-2256.
[11] B. Crespi, S.-J. Chang and K.-J. Shi, Elliptical billiards and hyperelliptic functions, J. Math. Phys., 34 (1993) 2257-2289.
[12] E. Previato, Poncelet theorem in space, Proc. Amer. Math. Soc., 127 (1999) 2547-2556.
[13] A. Cayley, Developments on the porism of the in-and-circumscribed polygon, Philosophical Magazine, 7 (1854) 339-345.
[14] Ph. Griffiths and J. Harris, On Cayley's explicit solution to Poncelet's porism, Enseign. Math., 24 (1978) 31-40.
[15] V. Dragović and M. Radnović, Conditions of Cayley's type for ellipsoidal billiard, J. Math. Phys., 39 (1998) 355362.
[16] V. Dragović and M. Radnović, On periodical trajectories of the billiard systems within an ellipsoid in $\mathbb{R}^{d}$ and generalized Cayley's condition, J. Math. Phys., 39 (1998) 5866-5869.
[17] V. Dragović and M. Radnović, Cayley-type conditions for billiards within $k$ quadrics in $\mathbb{R}^{d}$, J. Phys. A: Math. Gen., 37 (2004) 1269-1276.
[18] V. Dragović and M. Radnović, Geometry of integrable billiards and pencils of quadrics, J. Math. Pures Appl., 85 (2006) 758-790.
[19] V. Dragović and M. Radnović, A survey of the analytical description of periodic elliptical billiard trajectories, J. Math. Sciences, 135 (2006) 3244-3255.
[20] V. Dragović and M. Radnović, Hyperelliptic Jacobians as billiard algebra of pencils of quadrics: Beyond Poncelet porisms, Adv. Math., 219 (2008) 1577-1607.
[21] V. Dragović and M. Radnović, Ellipsoidal billiards in pseudo-Euclidean spaces and relativistic quadrics, Adv. Math., 231 (2012) 1173-1201.
[22] H. Knörrer, Geodesics on the ellipsoid, Inv. Math., 59 (1980) 119-143.
[23] M. Audin, Courbes algébriques et sistèmes intégrables: géodésiques des quadriques, Expo. Math., 12 (1994) 193226.
[24] P. S. Casas and R. Ramírez-Ros, The frequency map for billiards inside ellipsoids, SIAM J. Applied Dynamical Systems, 10 (2011) 278-324.
[25] P. S. Casas and R. Ramírez-Ros, Classification of symmetric periodic trajectories in ellipsoidal billiards, CHAOS, 22 (2012) 026110.
[26] D. G. Mead, Newton's identities, Amer. Math. Monthly, 99 (1992) 749-751.
[27] D. E. Knuth, The Art of Computer Programming - Vol. 2: Seminumerical Algorithms, third edition, AddisonWesley, 1998.
[28] A. Delshams, Yu. Fedorov and R. Ramírez-Ros, Homoclinic billiard orbits inside symmetrically perturbed ellipsoids, Nonlinearity, 14 (2001) 1141-95.
[29] Yu. Fedorov, Classical integrable systems and billiards related to generalized Jacobians, Acta Appl. Math., 55 (1999) 251-301.


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