

# Tom Bridgeland - Joyce's work on counting

Note Title

3/31/2007

invariants and wall-crossing

Work of Dominic Joyce

(today: joint discussions with Valerio Toledano Laredo)

D triangulated  
category



Stab D

complex manifold:

"space of stability conditions"

in particular  
construct holomorphic  
functions



should carry interesting  
geometric structure

Baby model:  $\mathcal{A}$  abelian category

$R$  finite dim  $k$ -algebra - eg path algebra  
of a quiver

$$\mathcal{A} = \text{Mod}_{\text{fd}}(R)$$

$\mathcal{A} \rightsquigarrow \text{Stab } \mathcal{A}$

Recall the Grothendieck group

$$K(\mathcal{A}) = \mathbb{Z} \cdot \text{Ob } \mathcal{A} / [B] = [A] + [C]$$

when  $0 \rightarrow A \rightarrow B \rightarrow C \rightarrow 0$

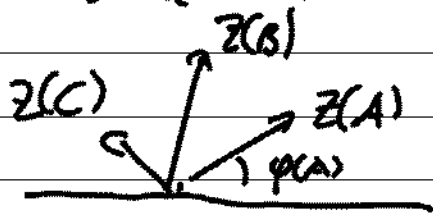
In our case  $K(\mathcal{A}) =$  free abelian group  
on simple  $R$ -modules  $\sim$ ,  
 $K(\mathcal{A}) = \mathbb{Z}^{\oplus n} = \bigoplus \mathbb{Z}[S_i]$

Positive cone  $K_{\geq 0}(\mathcal{A}) = \{ [A] : 0 \neq A \in \mathcal{A} \}$   
 $= \{ \sum k_i [S_i] : k_i \geq 0, \text{ not all } = 0 \}$

Def A stability condition on  $\mathcal{A}$  is a group  
homomorphism  $Z : K(\mathcal{A}) \rightarrow \mathbb{C}$   
s.t.  $Z(K_{\geq 0}(\mathcal{A})) \subset \mathbb{H}$  upper half plane

[general D: take t-structure & glue together  
such  $Z$ 's on heart...]

Today we're just studying a local model  
for  $\text{Stab}(\mathcal{D})$ , which is all we need to  
describe functors]



Write  $\text{Stab}(\mathcal{A})$  for the set of stability conditions  
 $\text{Stab } \mathcal{A} \simeq \mathbb{H}^n$  via  $Z \mapsto (Z(S_1) \dots Z(S_n))$

An object  $0 \neq E \in \mathcal{A}$  is semi-stable wrt  $Z \in \text{stab } \mathcal{A}$  if

$$0 \neq A \subsetneq E \Rightarrow \varphi(A) \leq \varphi(E)$$

where the phase  $\varphi(E) = \frac{1}{\pi} \arg Z(E) \in (0, 1)$

Example  $\mathcal{A} = \text{Rep}(\bullet \rightarrow \bullet)$

$\mathcal{A}$  has two simple objects  $S = (\mathbb{C} \rightarrow 0)$

$T = (0 \rightarrow \mathbb{C})$

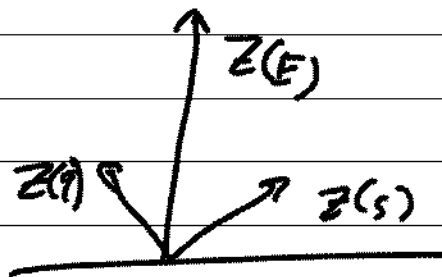
$E = (\mathbb{C} \xrightarrow{\lambda} \mathbb{C})$

↳ are other indecomposable

↳  $0 \rightarrow T \rightarrow E \rightarrow S \rightarrow 0$

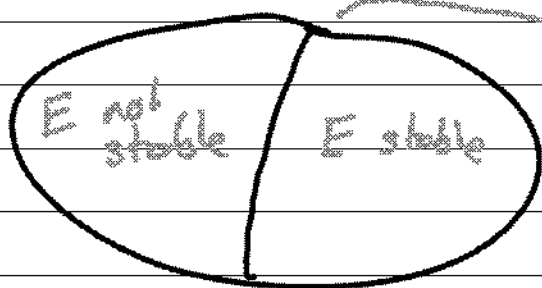
$$Z(E) = Z(S) + Z(T)$$

$S, T$  always semi-stable



$E$  semi-stable  $\Leftrightarrow \varphi(T) \leq \varphi(S)$

stab  $\mathcal{A}$   
is  
 $H^2$



real codim 1 wall:  
 $Z(S)/Z(T) \in \mathbb{R}_{>0}$

Aim: Want to make holomorphic functions  
Stab  $\mathcal{A} \rightarrow \mathbb{C}$  by "counting" semistable  
objects & forming a generating function.  
(weighted by  $\mathbb{Z}$ )

Obvious problem: Counting invariants are  
discontinuous due to wall crossing behaviour

Want discuss how to count (Donaldson-Thomas theory)  
but rather how to deal with wall crossing.

[ •  $\mathcal{A}$  = modules for a f.d algebra  $\Rightarrow$   
stability conditions do come from GIT  
stability ]

Notes i)  $\exists$  good moduli space  
 $M_{\mathcal{A}}(\alpha)$  parametrisng semistable objects  
of type  $\alpha \in K(\mathcal{A})$

ii) Harder-Narasimhan property:  
 $\forall 0 \neq E \in \mathcal{A} \quad \exists!$  filtration

$0 = E_0 \subset E_1 \subset \dots \subset E_n = E$  with  $F_i = E_i/E_{i-1}$   
semistable &  $\varphi(F_1) > \dots > \varphi(F_n)$ .

## Hall algebras

General idea:  $\mathcal{M}$  = stack of all objects of  $\mathcal{A}$

$\mathcal{M}'$  = stack of short exact sequences in  $\mathcal{A}$

$$\begin{array}{ccc} \{0 \rightarrow A \rightarrow B \rightarrow C \rightarrow 0\} \mathcal{M}' & \longrightarrow & \mathcal{M} \\ \downarrow & & \{B\} \\ \mathcal{M} = \mathcal{M} & & \text{correspondence.} \\ \{A\} \{C\} & & \end{array}$$

Applying a cohomology theory  
gives a convolution product

$$H^*(\mathcal{M}) \otimes H^*(\mathcal{M}) \longrightarrow H^*(\mathcal{M})$$

Second isomorphism theorem  $\Rightarrow$  this is associative.

One way to make this idea work:

Assume  $k$  finite field

$\mathcal{M}$  = set of isom classes of objects

$$= \bigsqcup_{\alpha \in K(k)} \mathcal{M}_\alpha.$$

Each  $\mathcal{M}_\alpha$  is a finite set!

Define  $H(\mathcal{A}) = \left\{ \begin{array}{l} \text{functions } f: M \rightarrow \mathbb{C} \text{ with} \\ \text{finite support} \end{array} \right\}$

$$= \bigoplus_{\lambda \in K(t)} \{ \text{functions } M_\lambda \rightarrow \mathbb{C} \}$$

Product:  $f * g(E) = \sum_{A \subseteq E} f(A) g(E/A)$  (Ringsel)

associative,  $f \circ g \circ h (E) = \sum_{A \subset B \subset E} f(A) g(B/A) h(E/B)$

Riesel:  $A = \text{Rep } Q$   $Q$  ruler with no loans

$$\Rightarrow H(A) = U_2 \eta_+ \quad \sigma_j = \eta_+ \oplus h \oplus \eta_-$$

$(\text{Re})_{\text{crit}}$  - critical  $\text{Re}$   $K_{\text{ex-Moody}}(Q)$

Define characteristic functions for  $\alpha \in K(A)$

$$1_\alpha \in H(A), \quad 1_\alpha(E) = \begin{cases} 1 & [E] = \alpha \\ 0 & \text{otherwise} \end{cases}$$

$$1_{\alpha}^{ss}(\bar{z}) \in H(X), \quad 1_{\alpha}^{ss}(\bar{z})(E) = \begin{cases} 1 & [E] = \alpha, \\ & \bar{z}\text{-semi-stable} \\ 0 & \text{otherwise} \end{cases}$$

ie  $1_{\alpha}^{ss}: \text{Stab } X \rightarrow H(X)$

Harder-Narasimhan property:

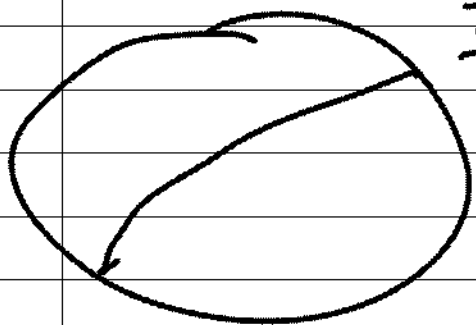
$$1_{\alpha} = \sum_{n \geq 1} \sum_{\substack{\alpha_1 + \dots + \alpha_n = \alpha \\ \varphi(\alpha_1) > \dots > \varphi(\alpha_n)}} 1_{\alpha_1}^{ss}(\bar{z}) * \dots * 1_{\alpha_n}^{ss}(\bar{z})$$

Reineke: inverted this recursion explicitly

$$1_{\alpha}^{ss}(\bar{z}) = \sum_{n \geq 1} \sum_{\substack{\alpha_1 + \dots + \alpha_n = \alpha \\ \varphi(\alpha_1 + \dots + \alpha_n) > \varphi(\alpha)}} (-1)^{n-1} 1_{\alpha_1} * \dots * 1_{\alpha_n}$$

So we know exactly how  $1_{\alpha}^{ss}(\bar{z})$  jumps as we cross walls in phase-space.

Consider  $1_{\alpha}^{ss}(\bar{z})$  on either side of a wall in  $\text{Stab } X$ :  
ie  $\alpha_1 + \alpha_2 = \alpha \in K(X)$



$$\frac{z(\alpha_1)}{z(\alpha_2)} \in \mathbb{R}_{>0} \quad [A_i] = \alpha_i, [A] = \alpha$$

$$0 \rightarrow A_1 \rightarrow A \rightarrow A_2 \rightarrow 0$$

we're looking at how  $A$   
switches from very stable  
to unstable.

$$\text{Wall crossing: } 1_{\alpha}^{ss} \mapsto 1_{\alpha}^{ss} \pm [1_{\alpha_1}^{ss}, 1_{\alpha_2}^{ss}]$$

commutator in  $H(A)$

- lose extensions going  $A_1 \rightarrow \_ \rightarrow A_2$

but gain extensions  $A_2 \rightarrow \_ \rightarrow A_1$  !

To make a holomorphic function

ex:  $f_{\alpha} : \text{stab } \mathcal{A} \rightarrow H_{\alpha}(\mathcal{A})$

should write

$$f_{\alpha}(z) = 1_{\alpha}^{ss}(z) + \frac{1}{2\pi i} \sum_{\substack{\alpha_1 + \alpha_2 \\ = \alpha}} \log\left(\frac{z(\alpha_1)}{z(\alpha_2)}\right) 1_{\alpha_1}^{ss}(z) \pm 1_{\alpha_2}^{ss}(z)$$

Correct for jump with a log which branches  
in same way. (branch of log jumping on  $\mathbb{R}_{>0}$ )



This just takes care of first order stuff  
— one term filtration! but have  
deeper intersection of walls to account for!

Joyce writes

$$f_\alpha(z) = \sum_{n \geq 1} \sum_{\substack{d_1 + \dots + d_n \\ = d \\ \alpha_i = 0}} J_n(z_{d_1}, \dots, z_{d_n}) I_{d_1}^{\text{ss}}(z) \times \dots \times I_{d_n}^{\text{ss}}(z)$$

for some functions with branch cuts!

$$J_1(z) = 1 \quad J_2(z_1, z_2) = \frac{1}{2\pi i} \log(z_1/z_2)$$

$J_n$ 's have branch loci chosen to balance the  
jumping behaviour of  $I_d^{\text{ss}}(z)$

— universal functions for any  $d$ .

Here in abelian category world get only finite sums  
 $\sum \alpha_i = \alpha$  ... in triangulated world gets  
much worse!

Surprising fact:  $J_n$  is essentially unique  
with this property,

& satisfy differential equation

$$dJ_n(z_1, \dots, z_n) = \sum_{i=1}^{n-1} J_i(z_1, \dots, z_i) J_{n-1}(z_{i+1}, \dots, z_n) \cdot d \log \left( \frac{z_1 + \dots + z_i}{z_{i+1} + \dots + z_n} \right)$$

(Bridgeland)

In fact if we insert the basis change

$$1_n^S(z) = \sum_{\substack{n \geq 1 \\ \lambda_1 + \dots + \lambda_n = 1}} K_n(z(z_1), \dots, z(z_n)) f_{\lambda_1}(z) \dots f_{\lambda_n}(z)$$

Then the  $K$ 's are given explicitly  
by path logarithms / iterated integrals

$$K(z_1, \dots, z_n) = \frac{1}{2\pi i} \int_{\gamma} \frac{dt}{t - s_1} \circ \dots \circ \frac{dt}{t - s_{n-1}}$$

$$s_n = z_1 + \dots + z_n$$

$$s_1 = z_1 + \dots + z_i$$

0

$$= \frac{1}{2\pi i} \int_{\gamma(\Delta^{n-1})} \frac{dt_1}{t_1 - s_1} \wedge \dots \wedge \frac{dt_{n-1}}{t_{n-1} - s_{n-1}}$$

integrated on a simplex

Analogy Consider a meromorphic connection

$$\nabla = d + \sum_{i=1}^n \frac{A_i dz}{z - u_i}$$

on a trivial  $\mathbb{C}^n$  bundle over  $\mathbb{P}^1$ , poles at  $u_i, \infty$ .

Riemann-Hilbert:  $A_1, \dots, A_n \mapsto M_1, \dots, M_n$   
monodromy matrices

$$\{ (u_1, \dots, u_n) : u_i \neq u_j \} \sim \text{Stab } \mathcal{A}$$

$$A_i \sim f_i$$

$$M_i \sim 1_i^{ss}$$

Analogy: isomonodromy: need to deform  $A$ 's in special way with parameters to keep  $M_i$  locally constant.  
(Schlesinger equation)

If fix paths rather than foliating then cross will get jumps ....

$\begin{matrix} B \\ \& \\ T-L \end{matrix}$ 
 Using  $f_a$  can write connection  

$$\nabla = d + \left( \frac{z}{t} + \frac{f}{t^2} \right) dt$$
  
 connection on  $H(t)$  bundle over  $\text{Stab } A$ .  
 Stokes matrices of this are the  $1^{ss}$ .

[ Dynkin case:  $\text{Stab } A$  is a complexified  
 Weyl chamber, has  $t$ -structures  
 corresponding to Weyl chambers,  
 need to glue flags together. ]

ie  $\nabla$  is really characterized by  
 the property that the  $1^{ss}$  are  
 its Stokes matrices!

Joyce was implicitly explicitly inventing  
 the Riemann-Hilbert correspondence!

Laplace transform for  $\nabla$  is a Fuchsian  
 connection on  $\mathbb{P}^1$ :

$$cI + \sum_{\alpha} \frac{\pi_{\alpha} \circ f}{t - z(\alpha)} dt \quad (\text{poles at } z(\alpha) \text{ many points!})$$

where  $\pi_{\alpha} : H(t) \rightarrow H_{\alpha}(t)$

The Stokes matrices become in this picture  
 certain transition functions:  
 take a solution near  $z(\alpha)$  which  
 lies in  $H_{\alpha}$  & see its poles upon  
 continuing to point  $z(\rho)$