

Topological Field Theory

Algebraic structure axiomatizing
coarse features of quantum field theory.

$\{M = \text{spacetime}\}$: smooth,
oriented n -dimensional manifold
[+ possible extra structures: spin etc.]

An n -dim TFT (with one tier)
is a multiplicative invariant Z (^{partition}_{function})
on manifolds M as above:

$$M \longmapsto Z(M) \in \mathbb{C} \text{ (say)}$$

$$\emptyset \longmapsto 1, \quad \sqcup \longmapsto \cdot$$

functorial: oriented diffeo. invariant.

Caricature: $Z(M) = \int e^{-S(\varphi)} D\varphi$

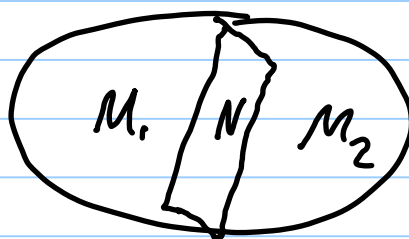
$\{\varphi\}$: fields - local expressions
on M (scalars, spinors,
sections of a bundle, maps to a
target space X , principal bundles,
connections, ...)

$S(\text{action})$: integral of local functional of φ .

Express locality of Z in M :

cut & paste formula

$$M = M_1 \amalg_N M_2$$



$$\Rightarrow Z(M) = \langle Z(M_1), Z(M_2) \rangle_{Z(N)}.$$

A 2-tier TFT assigns in addition

$$\bullet \quad Z: \begin{array}{c} (n-1)\text{-manifolds} \\ \text{compact, smooth, oriented} \end{array} \longrightarrow \text{Vect}_{\mathbb{C}}$$

$$N \longmapsto Z(N)$$

functorial & multiplicative:

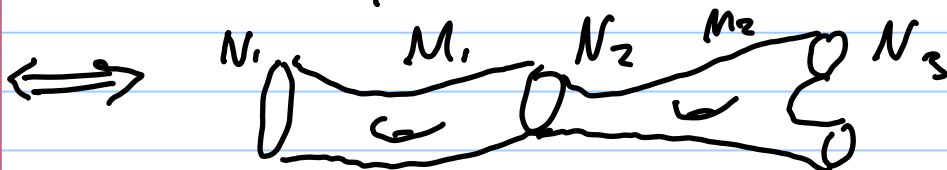
$$\emptyset \mapsto \mathbb{C}, \quad ()^{\text{op}} \mapsto ()^*, \quad \amalg \mapsto \otimes$$

$$\bullet \quad Z: M^n \text{ w/ boundary} \longmapsto Z(M) \in Z(\partial M)$$

so that above gluing rule holds:

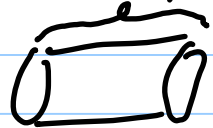
$$Z(M_1) \in Z(N), \quad Z(M_2) \in Z(N)^*$$

$\Rightarrow$ pair to give $Z(M)$.



$$Z(M) \in Z(\partial M_{\text{in}} \amalg \partial M_{\text{out}}^{\text{op}})$$

$$= \text{Hom}(Z(\partial M_{\text{in}}), Z(\partial M_{\text{out}}))$$

If we think of N as a fixed-time slice of $M \Rightarrow Z(N)$ is the Hilbert space of the field theory & this is just the Hamiltonian formulation of the theory ... except we set the Hamiltonian to be 0: 

$$Z(N \times I) = \text{id} \in \text{Hom}(Z(N), Z(N))$$

Construction of TFTs:

$Z(N)$ = linearization/quantization of (ie some kind of functions on) space of boundary values of fields $\varphi|_M$ (or jets)



$$Z(M, \varphi_2)$$

$$= \int_{\varphi|_N = \varphi_2} e^{-S(\varphi)} D\varphi$$

Evolution operators :
push-pull operators

$$\begin{array}{ccc}
 & \text{Fields}(M) & \\
 & \swarrow \quad \searrow & \\
 \text{Fields}(N_1) & & \text{Fields}(N_2) \\
 v \in \mathbb{Z}(N_1) & \longrightarrow & \mathbb{Z}(N_2)
 \end{array}$$

$$\mathbb{Z}(M)(v)(\varphi_2) = \int_{\varphi|_{N_2} = \varphi_2} v(\varphi|_{N_1}) e^{-S(\varphi)} \mathcal{D}\varphi$$

→ Review: Function Theory 101

X finite set $\rightsquigarrow \text{Fun}(X)$ ($\mathbb{C}$ -valued)

- Carries commutative multiplication
- $\pi : X \rightarrow Y$ map $\Rightarrow$

$$\begin{array}{ccc}
 \pi_X : \text{Fun}(X) & \xleftrightarrow{\quad} & \text{Fun}(Y) : \pi^* \\
 \text{sum along fibers} & & \text{pullback}
 \end{array}$$

- All linear operators $\text{Fun}(X) \rightarrow \text{Fun}(Y)$ are integral transforms, with a kernel function $K = K(x, y) \in \text{Fun}(X \times Y)$

$$\begin{array}{c}
 K \\
 X * Y \\
 \swarrow \pi_1 \quad \searrow \pi_2 \\
 X \quad Y
 \end{array}
 \quad
 K * f(Y) = \int f(x) K(x, y) dx$$

$$= \pi_{2*} (\pi_1^* f \cdot K)$$

$$\dots \text{Fun}(X) \cong \mathbb{C}^n, \quad \text{Fun}(Y) \cong \mathbb{C}^m,$$

$$\text{Fun}(X * Y) = \text{Mat}_{m \times n} :$$

this is just matrix representation of maps..

Convolution algebras :

$$\begin{array}{c}
 X * X * X \\
 \swarrow \pi_{12} \quad \searrow \pi_{23} \\
 X * X \quad X * X \\
 \downarrow \pi_{13} \\
 X * X
 \end{array}$$

Matrix multiplication

$$A * B = \pi_{13*} (\pi_{12}^* A \cdot \pi_{23}^* B)$$

$\text{Fun}(X * Y)$ likewise: block-diagonal matrices

$$\begin{pmatrix}
 \boxed{y_1} & & \\
 & \boxed{y_2} & \\
 & & \boxed{y_3}
 \end{pmatrix}$$

G group $\rightsquigarrow$ group algebra $\text{Fun}(G) = \mathbb{C}G$

$$\mu: G * G \rightarrow G, \quad f * g = \mu_* (\pi_1^* f \cdot \pi_2^* g)$$

Example 2d topological gauge theory:
 G a finite gauge group

Fields : principal G -bundles on M
 $\iff G$ -covering spaces of M
 $\iff$ (flat) G connections on M
 $\iff \{ \pi, (M) \longrightarrow G \} / \text{conj.}$
 $\mathcal{M}_G(X)$

$$Z(\Sigma) = \# \text{ fields} = \# \{ \pi, (\Sigma) \longrightarrow G \} / \sim$$

counted w/o auto.

$$= \# \{ A_1, \dots, A_g, B_1, \dots, B_g \in G : \prod [A_i, B_i] = 1 \} / \sim$$

How do we calculate this? (Frobenius, Schur)
 - cut & paste:

$$Z(S') : \text{linearization of Fields}(S')$$

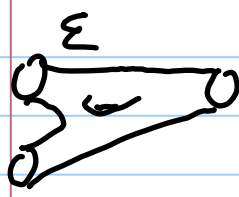
$$= \frac{G}{|G|} \quad \text{adjoint quotient :}$$

conjugacy classes

- clear meaning for functions on a finite set

$$\Rightarrow Z(S') = \text{Fun}\left(\frac{G}{G}\right) = (|G|)^{|G|}$$

class functions.



$$Z(\Sigma) : Z(S')^{\otimes 2} \longrightarrow Z(S')$$

push-pull along fields on Σ .

Theorem $2d$ TFT (2 tiers) $\iff$
 commutative Frobenius algebras
 (i.e. unital commutative associative algebra
 with nondegenerate trace)

How? $\mathbb{Z} \mapsto H = \mathbb{Z}(S')$

$$\mathbb{Z}(\bigcirc) = 1 \in H \quad \mathbb{Z}(\bigcirc) = \text{tr} : H \rightarrow \mathbb{C}$$

$$\mathbb{Z}(\text{cup}) : H \otimes H \xrightarrow{\mu} H.$$

check $\mathbb{Z}(\text{cup}) = \text{tr}(\mu) : H \otimes H \rightarrow \mathbb{C}$
 nondegenerate. Moreover any Σ^2 can be
 decomposed into such building blocks, so recover
 $\mathbb{Z}(\Sigma)$ from Frobenius structure.

Commutativity $\text{cup} \leftrightarrow \text{cap}$

Gauge theory example:

$\mathbb{C}G$ is a Frobenius algebra:
 $\text{tr}(f) = f(1)$ nondegenerate trace
 (dim of a character..)

$\mathbb{C}G^G$ is its center, a commutative Frobenius algebra
 $\iff$ local operators.

Geometric basis: conjugacy classes.

To solve the system: want to simultaneously diagonalize the local operators $\longleftrightarrow$ spectral decomposition

$\Rightarrow$ calculate Spec $\mathbb{C}G^G$

$= \{ \text{homomorphisms } \rightarrow \mathbb{C} \}$

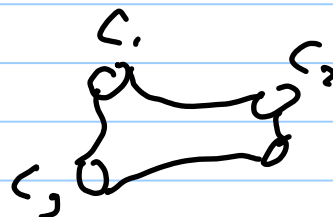
$= \{ \text{joint eigenvalues} \} = \hat{G}$

$= \{ \text{irred characters of } G \}$

"Quantum mechanics of the vacuum $[H=0]$ is completely integrable" - simultaneously diagonalize all local operators with nondegenerate spectrum!
 [Gln: get quantum many-body systems...]
 — in this basis easier to calculate:

$$\approx \mathbb{Z}(\mathbb{Z}) = \sum_{\substack{\chi \in \text{irred} \\ \text{chars of } G}} \left(\frac{|G|}{\dim \chi} \right)^{2g-2}$$

Similar formula for



for conjugacy classes

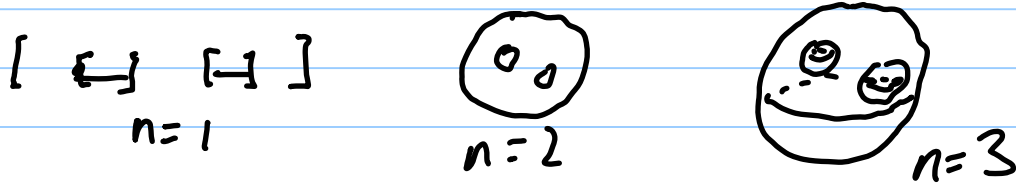
at punctures $\longleftrightarrow$ insert corresponding

operators $\# \{ \prod [A_i, B_i] \prod C_i = 1 \}$

$C_i \in \{C_i\}$

Local operators

In any TFT, $Z(S^{n-1})$ is an algebra, given by "pair of pants"



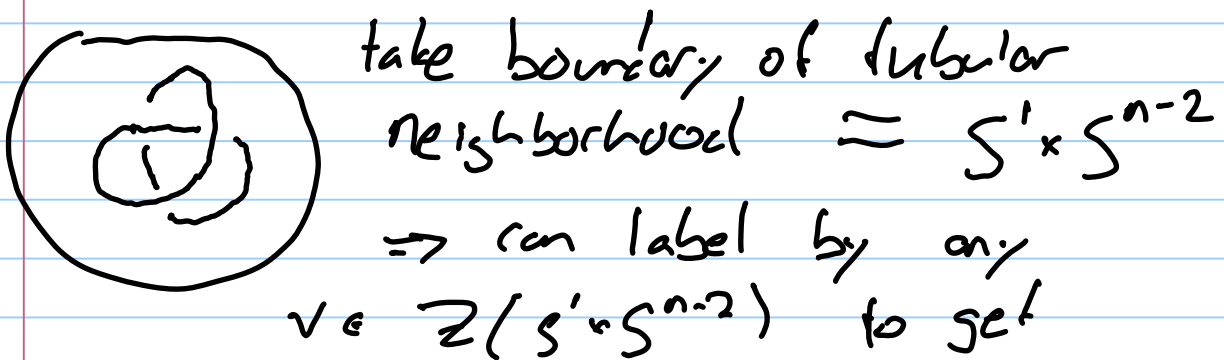
topological OPE : associative ($n=1$)
commutative ($n \geq 1$)

$\downarrow Z(S^{n-1}) \hookrightarrow Z(N)$ for any N as local operators :

Diagram of a cylinder $N \times I$ with a loop γ on the boundary. The loop is labeled $\gamma \in Z(S^{n-1})$.

$$\Rightarrow \int \mathcal{O}_{\gamma,t}(\varphi) e^{-S(\varphi)} D\varphi$$

Other labels : e.g. given loop in n -manifold



$\Rightarrow$ can label by any $v \in Z(S^1 \times S^{n-2})$ to get invariant of the pair, or loop operators.

Order & Disorder operators:

Note that the operators we insert are labelled by functionals on fields on the link of our the point x of insertion.

$\Rightarrow$ in addition to local functionals of fields can have functionals supported on fields with prescribed singularities at x

$\Leftrightarrow$ changing the path integral from regular fields to singular fields!

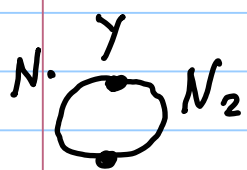
eg in gauge theory: point operators in 2d are all disorder operators, introduce monodromy around this point: [t Haff!] but loop operators can calculate trace of monodromy around loop in some representation, Wilson operator.

Extended TFT

Express locality of the vector space
 $Z(N)$ with respect to cut & paste in N

ie $N = N_1 \amalg_Y N_2 \Rightarrow$

$$Z(N) = \langle Z(N_1), Z(N_2) \rangle_{Z(Y)}$$



To get a vector space as
a pairing $\rightsquigarrow$ let $Z(Y)$ be
a linear category & $Z(N) \in \langle \rangle_{Z(\partial N)}$.

Put another way: before for 
we got a function ($\mathbb{C}$ -valued) on space of fields on N from path
integral on M : fixing boundary value
makes M into a closed n -manifold.

Now  fixing values on ∂N should
get a vector space

$\rightsquigarrow$ a vect-valued function, or sheaf,
on fields on Y $\therefore Z(Y) = \text{category}$
of sheaves on $\text{Fields}(Y)$.
a categorified linearization / quantization.

Slightly more formally: 2-category

Objects: $n-2$ manifolds Y •

1-morphisms: $n-1$ manifolds N w/ boundary 

2-morphisms: n manifolds M w/ corners 

An n -dim TFT with 3 tiers is
a functor to some target 2-category,
e.g. Ob: linear categories

1-Mor: Functors

2-Mor: natural transformations

Explored by Freed, Baez-Dolan,

2d: Moore-Segal, Costello: complete picture

Hopkins-Lurie in general: go
down all the way to a point: i.e.
totally local. All topology disappears, get
only algebra.

Theorem (Hopkins-Lurie) An n -dim TFT
extended down to a point is completely
& freely determined by the **self-dual**
object assigned to a point.

Example, 2d topological gauge theory
(G finite)

Linear category: {vector bundles on
gauge fields on a point}

$$= \text{Vect}(BG) = \{\text{linear reps of } G\}$$

$\cong$ $\mathbb{C}G$ -mod, modules over
the group algebra of G .

Note that vector bundles (or sheaves)
on finite sets also have $\otimes, \pi^*, \pi_*$
like functions $\Rightarrow$ define operations
for cobordisms
by push-pull.

$\text{Fields}(N)$
↙ ↘
 $\text{Fields}(Y, 1)$ $\text{Fields}(1/2)$

Structures:

V, W

$$V, W \in \text{Rep } G \Rightarrow \text{Hom}_G(V, W) = \text{Vect}$$



composition

End V is an associative algebra.

What's special about this category?

Modules over a Frobenius algebra ($\mathbb{C}G$)
form a Calabi-Yau category:

have nondegenerate traces on $\text{End } V$
making them into NC Frobenius algebras:

$$\text{cylinder} : \mathbb{C} \rightarrow Z(S') = \mathbb{C}G^e \rightarrow \text{End } V$$

$= Z(\mathbb{C}G)$

action of center

$$\text{cylinder}^* = \text{cylinder} \quad \text{adjoint:}$$

$$V \simeq \bigoplus V_i \otimes \mathbb{C}^{n_i} \quad V_i \text{ irreducible} \Rightarrow$$

$$\text{End } V \longrightarrow Z(S') = \mathbb{C}G^e \longrightarrow \mathbb{C}$$

$\downarrow \qquad \qquad \qquad \downarrow \qquad \qquad \qquad \downarrow$

$$\bigoplus \text{Mat}_{n_i, n_i} \ni 1; \longmapsto \text{char } V \longmapsto \dim V$$

character map into class functions

Moore, Segal, Costello, Hopkins-Lurie:

$Z(\cdot)$ in a 2d TFT is a CY category

$Z(S')$ is its center (HH^*): acts centrally
on every object $\text{---} = \text{End}(\text{Id})$.

Example σ -model on finite orbifold

$$X = \coprod \mathbb{C}^n / H_i$$

$$\begin{aligned} Z(\Sigma) &= \# \text{ Map}(\Sigma \rightarrow X) \\ &= \sum_i (\# H_i\text{-bundles}) \end{aligned}$$

$$\begin{aligned} Z(S^1) &= \text{Fun}(\text{Map}(S^1 \rightarrow X)) \\ &= \bigoplus_i \mathbb{C} \left[\frac{H_i}{h-1} \right] = \mathbb{C} H^H \end{aligned}$$

$$\begin{aligned} Z(\bullet) &= \text{Vect}(\text{Map}(\bullet \rightarrow X)) \\ &= \bigoplus_i \text{Rep } H_i \end{aligned}$$

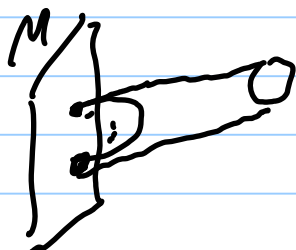
The character map (D-brane $\Rightarrow$ D-brane stack)

$$M \in Z(\bullet) \mapsto Z(\text{torus}) (1_M) \in Z(S^1)$$

is the Chern character of a vector bundle.

$$\dots \quad \text{Fun}(LX)^{S^1} = \text{cyclic homology} \\ (\text{roughly de Rham cohomology})$$

Stringy picture



open-closed
string transition

Local operators in extended TFT

In any TFT, $Z(S^{n-2})$ is an algebra (monoidal category)

$$[\text{---}] \quad [\text{---}]$$

$n=1$

associative
(monoidal)



$n=2$

braided



$n=3$

symmetric

$n=4$

$V \otimes W \xrightarrow{\sim} W \otimes V$
but square is not identity!
 $\Rightarrow$ braid group action on $V^{\otimes n}$

Note now in codimension 2 have
3 different possibilities for the
type of algebra we get!

Structure of 4d TFT:

- Σ closed surface $\mapsto$ category $Z(\Sigma)$.

In fact $\Sigma \mapsto$ 2d TFT Z_Σ with

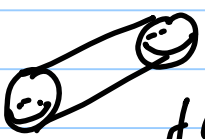
$$Z_\Sigma(C) = Z(\Sigma * C) \Rightarrow$$

$Z(\Sigma)$ is a Cobbi-Yau category

- Fields on $\Sigma * C$ are some kind of maps $(\rightarrow \text{Fields}(\Sigma))$:
- ie Z_Σ is some kind of σ -model on $\text{Fields}(\Sigma)$.

- $S^2 \mapsto$ tensor (symmetric monoidal) category

- $\forall x \in X$ get action $Z(S^2) \hookrightarrow Z(X)$
from , all commute

- $x, y \in X$ pick path $x \rightsquigarrow y$
 $\Rightarrow$  get isomorphism of
the x & y actions of $Z(S^2)$
depending only on the topology of
the corresponding cobordism

$\Rightarrow$ flat connection on the actions

- categorified quantum mechanical system

"Hilbert space" $Z(X)$, operator algebra $\supset$

$\mathcal{H}(X) := \{ Z(S^2) \}_{x \in X}$ large commuting family.

Try to solve - simultaneously diagonalize

- & hope for complete integrability:
non degeneracy of eigenstates!

Example 4d gauge theory, G finite.

$$\text{Fields}(S^2) = BG, \quad Z(S^2) = \text{Rep } G.$$

X
Surface

$$\text{Fields}(X) = \mathcal{M}_G(X)$$

$$= \{ G\text{-bundles} \}$$

$$= \coprod_i \bullet / \text{Aut } P_i = \coprod_i B(\text{Aut } P_i)$$

$\text{Aut } P \subset G$ centralizer of monodromy

$$Z(X) = \text{Vect}(\text{Fields}(X))$$

$$= \bigoplus_i \text{Rep}(\text{Aut } P_i)$$

$$x \in X \quad \forall v \in \text{Rep } G = Z(S^2)$$

$$\Rightarrow \text{Wilson operator } W_{x,v} \hookrightarrow Z(X)$$

$$W_{x,v}(F) = F \circ \underline{W_{x,v}}$$

$$\underline{W_{x,v}}|_P = \text{fiber at } x \text{ of bundle associated to } P \text{ in rep. } V$$

$$\cong V|_{\text{Aut } P} : \text{Fibers at } x \rightsquigarrow y$$

isomorphic but for loop $x \rightsquigarrow y$ get monodromy of associated bundle!

Function-sheaf correspondence

Compare X & $X \times S'$

$$Z(X \times S') = K(Z(X)), \text{ where } \text{traces lie.}$$

concretely, $\text{Fields}(X \times S') = G\text{-bundles on } X \times S' = \coprod \frac{H_i}{H_i}$

$$\begin{aligned} Z(X \times S') &= \text{Fun} \left(\coprod \frac{H_i}{H_i} \right) = \bigoplus \mathbb{C}[H_i^{H_i}] \\ &= \bigoplus \text{Rep ring of } H_i \end{aligned}$$

Wilson loops: $x \in X$ get loop l_x in $X \times S'$. (analog of Frobenius!)

$V \in \text{Rep } G \Rightarrow W_{x,V} \in Z(X \times S')$
multiplication by a function

$W_{x,V}(p) = \text{trace of holonomy along } l_x \text{ in representation } V$

Dimensional reduction $\longleftrightarrow$
de categorification:

categories $\xrightarrow{K} \text{Vect}$, Sheaves $\xrightarrow{\text{tr}} \text{functions}$

NC Hamiltonian System:

To "solve" the TFT, we try to simultaneously diagonalize all the local operators $\mathcal{H}(X) := \{Z(S^2)_x\}_{x \in X}$ i.e. perform spectral decomposition.

So need to find joint spectrum of the commutative algebra [category] acting:

- R ring $\Rightarrow$
 $\text{Spec } R(k) = \text{Ham}(R, k)$
& $R =$ functions on $\text{Spec } R$
 R -modules localize as sheaves on $\text{Spec } R$
- $\mathcal{C}$ tensor category $\Rightarrow$
"Spec" $\mathcal{C}(k) = \text{Tensor functors}(\mathcal{C}, \text{Vect}_k)$
& $\mathcal{C} =$ sheaves on "Spec" $\mathcal{C}$ (usually...)
[$\mathcal{C}$ -modules localize on "Spec" $\mathcal{C}$]
- the tensor functors become evaluation
e.g. $\mathcal{C} = \text{Rep } G^\vee \Rightarrow$
"Spec" $\mathcal{C} = BG$ & $\text{Vect}(BG) = \mathcal{C}$.
Tannakian reconstruction!

Examples: • $\mathcal{C} = \{ \text{graded vect} \}$
 $\Rightarrow G^\vee$ (Tannaka/Galois group)
is $\mathbb{C}^*$

• $\mathcal{C} = \{ \text{flat vector bundles on } X \}$
 $\Rightarrow G^\vee = \pi_1(X)$.

TFT: Suppose $Z(S^2)$ is Tannakian,
ie $Z(S^2) \simeq \text{Rep } G^\vee$ for some
group $G^\vee \iff \text{"Spec" } Z(S^2) = BG^\vee$

$\Rightarrow \text{"Spec" } \mathcal{H}(X) = \text{Loc}_{G^\vee} X$
flat $G^\vee$ -bundles on X !

So $Z(X)$ is a module over
"functions" (Vect) on $\text{Loc}_{G^\vee} X$

— ie our theory is trying to be
 $G^\vee$ gauge theory, local operators $\leadsto$ Wilson!

[G finite $\leadsto G^\vee = G$, boring]

it is when our quantum system is integrable!

Geometric Langlands conjecture

asserts that a particular 4d TFT

$$\mathbb{Z}_G \quad ["N=4 \text{ SYM in GL twist} \\ \text{with } \psi=0"]$$

is completely integrable in this sense

with "Galois" group $G^\vee = \text{Langlands dual group}$:

in our theory $\mathbb{Z}_\Sigma = \text{A-model on } T^* \text{Bun}_G \Sigma$

$\mathbb{Z}(\Sigma) = \text{A-branes on } T^* \text{Bun}_G \Sigma$

- eg Fukaya category (Lagrangians)

$\longleftrightarrow$ D-modules on $\text{Bun}_G \Sigma$
Kupersin-Witten
Nakler-Zaslav

- eg holonomic

... singular bundles with flat connection on Bun_G

$$\mathbb{Z}(S^2) = \text{Shv}(\text{Bun}_G S^2) \simeq \text{Rep } G^\vee$$

- Geometric Satake Theorem

Mirkovic-Vilonen, (Lusztig, Ginzburg, Drinfeld)

$\Rightarrow \text{Shv}(\text{Bun}_G \Sigma)$ has action of vector bundles on $\text{Loc}_{G^\vee} \Sigma$.

Conjecture: multiplicity one:

$$\text{Shv}(\text{Bun}_G \Sigma) \simeq D(\text{Loc}_{G^\vee} \Sigma)$$

[RHS get "span" of "Wilson" vector bundles = coherent sheaves]

$$\leadsto \mathbb{Z}_\Sigma \simeq \text{B-model on } \text{Loc}_G \Sigma$$

Motivations for the homological/dg world

- leaving the toy models!

- Physics TFTs in nature usually come as topological twists of Supersymmetric QFTs ("Witten type" or cohomological field theories - as opposed to "Schwarz type" like Chern-Simons, with manifestly topological Lagrangian)

- have bosons & fermions $\Rightarrow$ super vector spaces

- often have charge operator (ghost $\#$) $\Rightarrow \mathbb{Z}$ -graded space of states

- TFT arises as Q -cohomology of space of states of underlying QFT $Q^2=0$ (BRST operator - conserved supercharge)

- $\Rightarrow$ underlying dg vector space whose cohomology is ordinary ($\mathbb{Z}$ -graded) TFT.

Likewise the space of open string states with given boundary conditions is naturally a complex with BRST differential

So entire TFT picture is naturally refined to a dg version - which is better behaved! ("off shell")

• Math Nice function theories beyond
 finite sets : want $\circ, \pi^*, \pi_*$
 eg for compact oriented manifolds
 take $H^*(X)$ with $\cup, \pi^*$ & $\int_\pi$.
 (or K -theory, or, ...)

Poincaré + Künneth:

$$\begin{aligned} \text{Hom}(H^*(X), H^*(Y)) &= H^*(X) \otimes H^*(Y) \\ &= H^*(X \times Y) \ni K \text{ integral kernels.} \\ f &\mapsto \int_{\pi_2} K \cup \pi_1^* f \end{aligned}$$

For G compact Lie group can use
 this instead of Fun

[Frenkel-Hopkins-Teleman: similar approach
 to Chern Simons using twisted K -theory]

Cohomology is the shadow of a complex / quism
 (de Rham / singular / simplicial / Čech)

which is smarter & better behaved:

can't expect to get all the functoriality,
 & gluing we want - especially once we
 go to codim 2 - without consistently
 working on c/s level...

Geometric Function Theory

A new kind of functional & harmonic analysis in algebraic geometry, where we replace function spaces (L^2 , C^∞ , distributions, ...)

by various dg categories - forms of the derived category - of sheaves on algebraic varieties or stacks [orbifolds]

ie $\text{Fun} \rightsquigarrow \text{Shv}$

radim 1 $\rightsquigarrow$ radim 2

spaces of states $\rightsquigarrow$ categories of branes

[es Fourier-Mukai]

[Note: triangulated categories are bad.
need to keep open string BRST
complexes, not just their cohomology!]

Examples:

bounded } derived category of { perfect complexes
unbounded } { coherent sheaves
quasicoherent sheaves
D-modules }

Largest, best behaved: unbounded quasicoherent
[∞ -rank version of coherent] -

analog of distributions or L^2

- easiest to prove theorems there, then prove
"regularity" results. D-mods: harder!

Properties

Kontsevich's NC geometry

- Refined invariant: $D(X)$ detects deformation theory of X detects X up to discrete ambiguity (Toën-Vagstad)

Can recover $H_{\text{cl}}^*(X)$ ($\mathbb{Z}/2$ graded)

& H_{cl}^* & Hodge structure (notion of smooth & proper) & K-theory

CY case - equivalent to 2d dg TFT

- Interesting invariant: nontrivial dualities like T-duality $\longleftrightarrow$ Fourier-Mukai

- Has a "commutative ring" structure given by $\otimes$.

Moreover $X = \text{"Spec"} D_{\text{qcoh}}(X)$

(BZ- Frank's-Nadler Tannakian theorem)

$\Rightarrow$ theory of spectral decomposition for $D(X)$ -modules

- $f: X \rightarrow Y$ have

$$f_* : D(X) \rightleftarrows D(Y) : f^*$$

- don't make sense e.g. on vector bundles, need derived versions to get well behaved operations

- "Schwarz kernel theorem":

All "continuous" operators $D(X) \rightarrow D(Y)$ are given by integral transforms

$$K \in D(X \times Y) \quad X \times Y$$

$$K * F = T_{2*}(\pi_1^* F \otimes K) \quad \begin{array}{ccc} & \swarrow \pi_1 & \searrow \pi_2 \\ X & & Y \end{array}$$

[Toën for varieties,

BZ-Francis-Nadler for stacks, eg Bm_2]

- Künneth & Poincaré: [BZ-Francis-Nadler]

$$D(X \times Y) = D(X) \otimes D(Y)$$

$$D(X) \simeq \text{Hom}(D(X), D(\bullet))$$

All these tools mean we can build TFTs by push-pull operations for varieties & stacks just like we did for finite sets & orbifolds

[Castello construction of B-model,

Hopkins-Lurie construction of string

topology: M manifold,

$$Z_M(\text{pt}) = D_{\text{loc. cond}}(M), \quad Z_M(S) = C_*(LX)$$

- need one outgoing boundary compact]

Chomological (or dg) TFT

2 tiers $n\text{-Cob}$: category [or homotopical / version thereof...]
objects = closed $n-1$ manifolds

$\text{Hom}(N_1, N_2) = \text{Space}$ of n -manifolds M
giving cobordisms $N_1 \begin{array}{c} \text{---} \\ \text{---} \end{array} \text{---} N_2$
 M

Very versatile: spaces give $\text{Vect}(H_0)$
graded $\text{Vect}(H_*)$, or chain
complexes / quasi-isom (C_*) $\Rightarrow$
can easily generalize TFT:

dg TFT: functor

$\text{Chains}(n\text{-Cob}) \longrightarrow \text{dg Vect}$

ie $\bullet \quad N \longmapsto Z(N)$ complex

$\bullet \quad N_1 \begin{array}{c} \text{---} \\ \text{---} \end{array} \text{---} N_2 \text{ is a } 0\text{-cycle of } n\text{Cob}(N_1, N_2)$

$\Rightarrow$ get chain map of degree 0

$$Z(M): Z(N_1) \longrightarrow Z(N_2)$$

$\bullet$ Given k -dim family $\{M(x)\}: N_1 \rightarrow N_2$

- ie k -chain on $n\text{Cob}(N_1, N_2)$ -

get chain map of degree k $Z(N_1) \rightarrow Z(N_2)$

-eg homotopic M 's induce chain homotopic maps...

$$\Rightarrow C_*(\alpha(\partial_1(N_1, N_2))) \rightarrow \text{Hom}(Z(N_1), Z(N_2))$$

map of complexes.

Can pass to H_* or H^* $\Rightarrow$ graded or used TFT

Eg $n=2$ spaces of 2-manifolds with given boundary (eg empty) are homotopic to $\mathcal{M}_{g,n}$ moduli of Riemann surfaces $\simeq \text{BDiff } \Sigma_{g,n}$ (since Teichmüller space is contractible)

$\Rightarrow$ we are constructing cochains / cohomology classes on moduli of Riemann surface:

Topological conformal field theory:
this is precisely the structure of twisted $N=2$ SCFT (coupled to 2d gravity), $\rightsquigarrow$ topological strings.

What about branes?

Branes in topological string theories naturally form dg categories [of which triangulated categories are a pale shadow]: spaces of open string states come as graded (ghost #) vector spaces with BRST operator (differential): ie get category where Hom's are complexes (if everything understood up to quasi-isom)

3 tiers: Same story: 2-category

$Ob = n-2$ manifolds

$Hom(\gamma_1, \gamma_2) = \text{topological category of cobordisms } \gamma_1 \xrightarrow{N} \gamma_2$

$2Hom(N_1, N_2) = \text{space of cobordisms w/ corners } N_1 \xrightarrow{M} N_2$

[cf. Lurie: Morse Lectures, IAS '08 notes on my webpage]

dgTFT: $Z(\gamma)$ dg category

$Z(N)$ dg functors

$Z(M)$ natural transformations

Classification in two dimensions

Castella, Kontsevich-Siebertman classify

2d dg TFTs going down to a point (3 tiers)

Generalized by Hopkins-Lurie (ongoing)
to nclim TFTs going down to a point
valued in anything for any n .

Theorem 2d dgTFT $\iff$
NC Calabi-Yau

[strictly get $\dim=0$ CYs unless we twist
TFT definition by determinant line
 $\iff$ central charge]

Noncommutative Calabi-Yau of dim. n :
dg category with a nondegenerate trace
 $\mathrm{Hom}(A, A) \longrightarrow \mathbb{C}[n]$

$$[\mathrm{Hom}(A, B) \simeq \mathrm{Hom}(B, A[n])^*]$$

Serre duality is given by shift by n , $w=0$

Geometric construction: $X \subset Y \Rightarrow D_{\mathrm{cl}}^{\mathrm{cl}}(X)$
NC CY, "immersion" of space of

CYs into space of NC CYs.

So 2d TFTs are essentially Calabi-Yaus.

4d TFT gives in particular

- Σ surface $\Rightarrow Z_\Sigma$ NC CY
...roughly a CY variety, which
is $\text{Fields}(\Sigma)$. The 2d TFT
 Z_Σ is then a σ -model into this
NC CY.
- $S^2 \Rightarrow Z(S^2)$ has an algebra
structure, & acts on Z_Σ
for every $x \in X$. "t Hooft / Wilson loops"

Geometric Langlands : our 4d TFT
is a particular 4d gauge theory:

$$\text{Fields}(\Sigma) = T^* \text{Bun}_G(\Sigma) \text{ roughly}$$

Calabi-Yau

Z_Σ = some dg category
of sheaves on this
(D-branes in σ -model).

$Z(S^2)$ is [roughly] $\text{Rep } G^\vee$ dual group

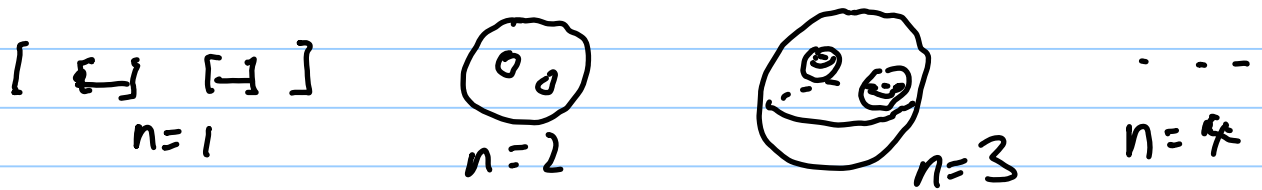
$$\text{"Spec"} \{ Z(S^2) \}_{x \in \Sigma} = \text{Loc}_{G^\vee} \Sigma$$

Geometric Langlands Conjecture (strong form)

These operators have simple spectrum: $Z(\Sigma) \simeq D(\text{Loc}_G \Sigma)$

Local operators

What structure does $Z(S^{n-1})$ have?



E_n -algebra: multiplication parametrized by space of discs ($\Leftrightarrow$ points) in n -disc: structure of n -dim OPE.

Even on level of graded Vect, all n are different: binary operations given by $H_*(2\text{-pt configs in } \mathbb{R}^n)$
$$= H_*(S^{n-1}) = \begin{cases} \mathbb{C} & \text{deg } n-1 \\ \mathbb{C} & \text{deg } 0 \end{cases}$$

$\Rightarrow$ An E_n graded vector space $V^\bullet$ ($n \geq 2$) has a commutative product (deg 0) & a Poisson bracket of deg $n-1$!

$n=2$ this is why vertex algebras look both like comm & Lie algebras!
- vertex algebras are a conformal refinement of E_2 algebras, where OPE can depend holomorphically, not just homotopically, on position.

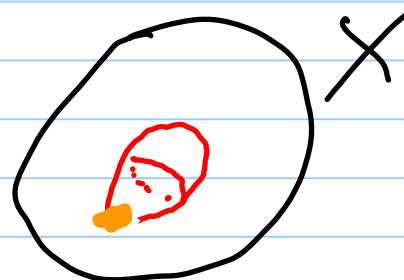
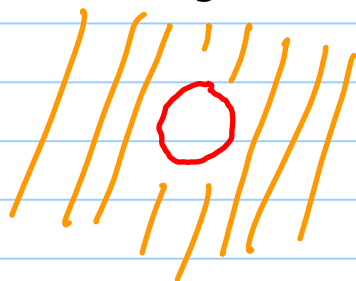
	Vect	Cal	gr Vect	dg Vect	dg Cal	Space
1	Assoc	Monoid	Assoc	A_∞	E_1	loop space
2	Comm	Braid	Gerstenhaber (odd Poisson)	E_2	E_2	double loop space
3		Symplectic	even Poisson	E_3	E_3	triple loop space
			odd Poisson	$\vdots$	$\vdots$	
			$\vdots$	$\vdots$	$\vdots$	
∞	$\downarrow$	$\downarrow$	Comm	E_∞	E_∞	infinite loop space. [connective] spectrum

Commutativity of n -dimensional OPE $\Leftrightarrow$
Commutativity of π_n of spaces :

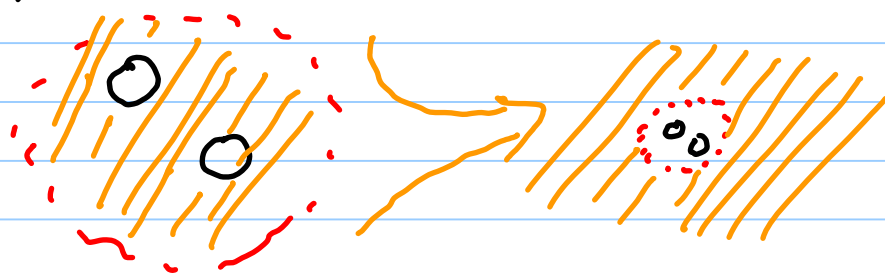
e.g $n=2$. $\pi_2(X)$ = components of $\Omega^2 X$

= iterated based loops in X

= maps from $D^2 \rightarrow X$ with boundary
going to a basepoint



Group structure on Π_2 comes from composition on iterated loops $\Omega^2 X$, which is an E_2 multiplication: labels! by little 2-discs:



rotate

up to homotopy: commutative since we can one around the other!

Vertex algebras: E_2 with holomorphic rather than homotopical dependence on positions
 $\leadsto$ hence vertex operators appear both as lie & as commutative algebras
 [Beilinson-Drinfeld, Gaiitsgory-Lurie]

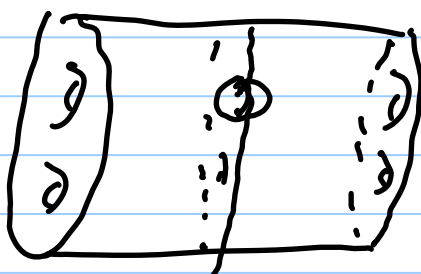
Theorem (May, Segal, ...)

A space Y is an n -fold loop space
 $Y \sim \Omega^n X$ for some $X \iff$

Y has an " n -dimensional OPE" structure (May version: E_n space, Segal version: iterably n -dim OPE)

Example: Geometric Satake theorem, or 't Hooft operators

Study local operators on G -bundles
(G complex or compact) on a Riemann
surface

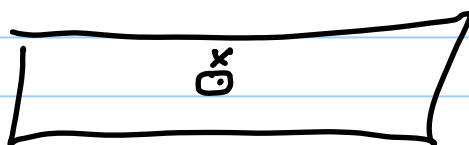


Solutions of
Bogomolnyi
equations
with singularity

zoom & slice



(lose some
symmetry)



local operators
on bundles
on a surface

(singularities of gauge fields)

$\mathcal{G} =$ G -bundles with trivialization
outside of x

$\cong LG/LG_+$

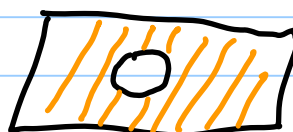


- clutching functions up to reparametrization

$\cong \Omega K$ based loops in the
compact form of G

$\cong \Omega^2(BK)$ double based loops in
the classifying space

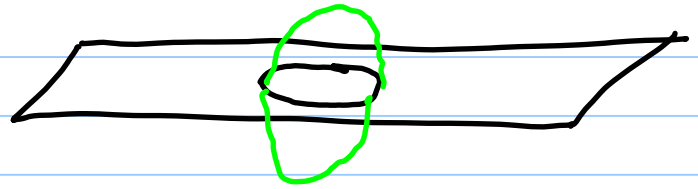
$=$ maps $D^2 \rightarrow BK$ with boundary
going to basepoint



$=$ K -bundles trivialized outside disc

[For 3d picture:

our modifications
are labelled by



$LG \backslash LG / LG$ double cosets = relative
positions of pairs of bundles $\Leftrightarrow$
instantons on the two-sphere.

These have an E_3 , not just E_2 , product.]

Hint at the CFT-Langlands connection?

the Kac-Moody vertex algebra arises
by linearizing the affine Grassmannian:

$G_r \rightsquigarrow \delta\text{-functions at the basepoint on } G_r$
 $\cong$ Vacuum K-M vertex
algebra

Construction of spaces of conformal
blocks $\longleftrightarrow$ insertion of "objects"
of the trivial surface operator



$$Z(\Sigma) = \langle Z(\Sigma, p), Z(0) \rangle_{Z(\Sigma)}$$

(Chern-Simons (3d analogy) :

$Z(\Sigma)$ = vector space of conformal blocks
= coinvariants of some representation
insertions (eg all vacuum)

$Z(S)$ = integrable level k reps of LG

Can construct vectors in $Z(\Sigma)$
from vectors in basic representation.

Likewise in our 4d TFT construct
objects of $Z(\Sigma)$ from objects
in $Z(D)$ - which are reps of
Kar-Moody?