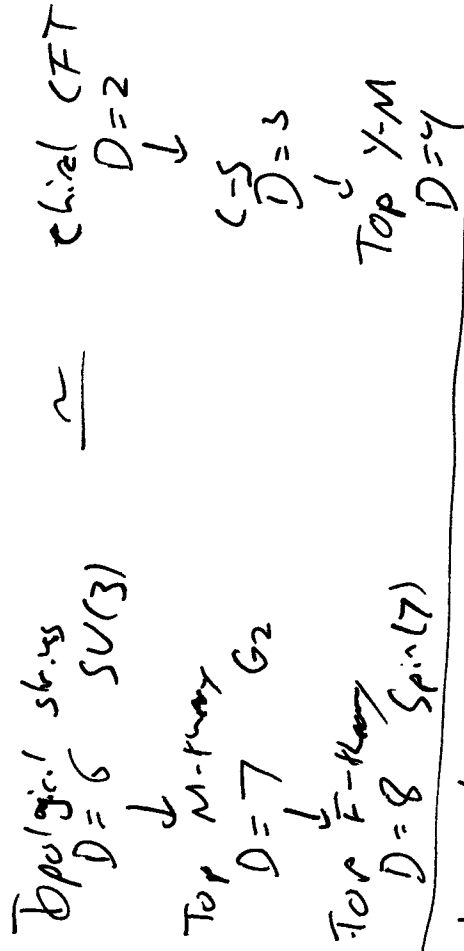


R. Dijkgraaf: Topological M-Theory
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Hints

Topological string partition for:

$Z_{top}(x)$ hol fn of $x \in M_{sp}$ moduli of vacua states
 well, hol anomaly $\implies Z_{top}$ wave function
 in Hilbert space of obtained by quantizing $H^3(X)$
 ... maybe think of $X = \mathcal{Z}(M^?)$,
 get particular state?

• 4dim Black hole partition: $Z_{BH} \sim Z_{top} \times Z_{top}^* \quad (\text{Ogura...})$
 $\implies$ theory on $X \times S^1 \dashrightarrow$ 7dim partition fn.

• 7dim geometry mixes Kähler & complex structure ---
 e.g. Swaminality $\longleftrightarrow$ β model

Topological string field theory: (Spectral POV)

BCOV X CY_3 : CX structure on X
 β -model: local coords $z^i, \bar{z}^i$. $\Omega = dz^i \wedge d\bar{z}^i \wedge d\bar{z}^3$
 (3,0) form
 Variations of CX structure: of form $\bar{\partial}$ by
 Keldner-Spencer field $\bar{\partial}_A = \bar{\partial} + A \bar{\partial}$

in comparison $\partial_j = A_j^i (z, \bar{z})$?

closed "field" $A' = \Omega \cdot A \in \Lambda^{2,1}$ closed string field for B-field
 ... gives info. defs of Ω , $\int \Omega = A'$.

Gauge condition $\partial A' = 0$

Write $A' = x + 2\phi$

$x \in H^{2,1}$ massless field, gives moduli
 (false as harmonic form)
 $A' = 2\phi$ massive modes (off-shell (fields))
 $\phi \in \Lambda^{1,1}$

$$\Rightarrow \text{KS action } S_{\text{KS}} = \int \frac{1}{2} \partial \phi \bar{\partial} \phi + \frac{g}{3} A' \wedge A'^2$$

$$Z_{\text{top}}(x) = \int D\phi e^{-S_{\text{KS}}}$$



Reduce KS to 2 dimensions: chiral theory
 similar to 2D chiral CFT

$$\int 2\phi \bar{\partial} \phi + g (2\phi)^3$$

... analog of only left moving, conformal field
 [holomorphic theory ... need to choose contours ...]
"integral on contour in p-h space"

A-model version -- Kähler gravity
 fundamental field is Kähler $(1,1)$ form $k = k_{i\bar{j}} dz^i d\bar{z}^{\bar{j}}$
 integrability $\Rightarrow dk = 0$

(Bershtinsky-Solov) Gauge $d_c^\dagger k = 0$ ($d_c = \partial - \bar{\partial}$)
 massless piece: $k = x$ $x \in H^{1,1}$ (moduli)
 massive off-shell fields: $k = d_c^\dagger \gamma$ $\gamma \in \Lambda^3$ 3-form
 Action $S_{\text{KG}} = \int \frac{1}{2} d\gamma \wedge d_c^\dagger \gamma + \frac{1}{2} \int k \wedge k \wedge k$
 ($k = x + d_c^\dagger \gamma \dots$)

→ gives $\mathbb{Z}^A$ in limit $\alpha \rightarrow 0$: ignores
 worksheet instructions, of form $\sum e - \int_2 (6 \cdot i \cdot g)$
 $\Sigma \rightarrow X$ RS.

Hitchin: p-forms in d-dimensions for special pt
 ... Stable forms: in $\Lambda^p(M^d)$ all forms
 are $GL(d, \mathbb{R})$ -equivariant ...
 typically $p = 3, 4$ $d = 6, 7, 8$

① 3-form p in $d=6$.
 $\dim \Lambda^3 = \binom{6}{3} = 20$. $G = GL(6, \mathbb{R})$ dim 36
 Stabilizer G_p has dim $36 - 20 = 16$.

Claim: $\exists$ open orbit of G of form

$$p = \alpha \wedge \beta. \quad \alpha, \beta \text{ decomposable: } \alpha = \omega_1 \wedge \omega_2 \wedge \omega_3$$

Case a. α complex, $\beta = \alpha^*$ conj-sab:

write as $p = \Omega + \bar{\Omega}$. $\Omega = e^1 \wedge e^2 \wedge e^3$, with
 e^i basis of $(T^* \otimes \mathbb{C})^{1,0}$. Stabilizer

$$G_p = SL(3, \mathbb{C}) \quad \dots \text{ "Fuchsian case"}$$

Case b. (Ninkowski case) α, β real
 $p = e^1 \wedge e^2 \wedge e^3 + e^4 \wedge e^5 \wedge e^6$ basis
 $G_p = SL_3 \mathbb{R} \times SL_3 \mathbb{R}$.

Case c: $p = \Omega + \bar{\Omega}$. Let $\hat{p} = \frac{1}{i}(\Omega - \bar{\Omega})$
 $\frac{1}{2}(p + i\hat{p}) = \Omega$... think of $\hat{p} = \text{Im} \Omega$
 $p = \text{Re} \Omega$.

If $d\rho = 0$ & $d\hat{p} = 0 \Rightarrow d\Omega = 0$
 & Ω defines integrable complex structure

$$\text{Action: } S_H(p) = \int \frac{1}{2} p \wedge \hat{p} = i \int \Omega \wedge \bar{\Omega}$$

$$\text{Variation } \delta S_H(p) = \int \delta p \wedge \hat{p}$$

$\Leftrightarrow S_H$ homogy of degree 2 in p , but not quadratic

$$K_p : \frac{\text{target}}{T} \xrightarrow{\text{S.G.M}} \Lambda^5 = \Lambda^6 \otimes T$$

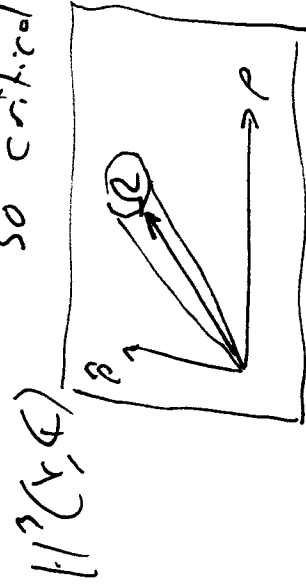
$$K_p(u) = i v p \wedge p$$

$$\text{Tr } K_p = 0, \quad \sqrt{-\frac{1}{6} \text{Tr } K_p^2} = p \wedge p \dots$$

get square root of quadratic functional in $p \dots$

Assume $dp=0$, vary p in cohomology class $X=[p] \in H^3(X_0)$
 $p: p_0 + d\beta \Rightarrow \delta S_H = \int d\beta \wedge p = - \int \beta \wedge dp$

$\Rightarrow d\hat{p}=0$ at critical points of S_H
 so critical pts give complex structures!



get one one moduli of complex structures as graph of $\hat{p}$ as function of p .

$$|1^3(x, R)$$

Quantum Hilbert theory:

$$Z_H(x) = \int D\beta e^{-\frac{1}{2} \int H(x + \beta p)}$$

Claim: $Z_H \propto Z_{\text{top}} \bar{Z}_{\text{top}}$ for topological B-model

(2) $\sigma \in \Lambda^4$ on X^6 : $\dim \Lambda^4 = \dim \Lambda^2 = \binom{6}{2} = 15$.
 G has $\dim 36 - 15 = 21$. Generically $G_\sigma = \text{Sp}(6, \mathbb{R})$
 $\leadsto$ get symplectic structures!
 Stable $\sigma \Rightarrow \sigma = \frac{1}{2} k \wedge k, k \in \Lambda^2$

Action $V = \int \frac{1}{3!} k \wedge k \wedge k = \int \frac{1}{3} \sigma \wedge k$
 has $\deg 3/2$ as function of σ
 $\int V = \frac{1}{2} \int \sigma \wedge k$.

Assume $d\sigma = 0$ $\sigma = \sigma_0 + \delta\sigma(\sigma)$

Critical points: $d\sigma = 0$ - independent symplectic structure

$$\int \sigma = \int \omega \wedge \sigma = 0$$

Use left-invariant operators $L = \tau \wedge -$ $\sigma \wedge 1 = \tau$

$$[L, d] = d_c^*$$

$$\Rightarrow \int \tau = d_c^* \tau + d(1 \wedge \tau)$$

$$\int \tau = \int \frac{1}{2} d\tau \wedge d_c^* \tau + \frac{1}{2} \int \tau^3 \Rightarrow \text{Kähler geometry}$$

(2)

3-form Φ or 4-form G on M^7

$$\Phi = \Phi_{ijk} dx^i \wedge dx^j \wedge dx^k \quad \dim 1^2 = 35 \quad \dim GL_7 \mathbb{R} = 49$$

$$\dim G\Phi = 14, \text{ geometrically get } G\Phi = G_2$$

$\Rightarrow \Phi = \sum_{ijk} e^i \wedge e^j \wedge e^k$ with C ∞ structure coeffs of octonions
Equivalently write $R^7 = \mathbb{R}^3 \times \mathbb{R}^4$

$$\varphi = \sum_a e^a \wedge \omega_a + e^1 \wedge e^2 \wedge e^3 \quad e^a, \omega_a \text{ self-dual 2-forms}$$

$$\Rightarrow \text{metric } ds^2 = \sum_{ij} c^{ij} dx^i dx^j = \sum e^i \wedge e^i$$

$$*\Phi = G, \quad *G = \Phi \quad 1^2 \ni \Phi_i = \Phi_{ijk} dx^j \wedge dx^k$$

$$\Rightarrow \text{metric } \beta_{ij} = \Phi_i \wedge \Phi_j \wedge \Phi$$

$$(\det B)^{1/9} \in 1^7 \text{ is volume form}$$

$$\mathcal{L} = \int G = (\det B)^{1/9} = \Phi \wedge \Phi \wedge \Phi = G \wedge *G$$

$$\int_{M^7} = \int_{S^3} = \int G \wedge *G \text{ has degree } 7/4 \text{ is } G, 7/3 \text{ is } \Phi$$

$$\int_{S^4} = \frac{7}{4} \int_{S^3} \int_{S^1} G \wedge *G$$

$$\forall \omega \in G \text{ in fixed cohomology class: } G = G + dC,$$

$$\int_{G_0} \omega = \int_{S^4} \omega$$

$$\int_{S^4} \omega = 0 \Rightarrow dG = 0, d * G = 0 \Rightarrow \text{metric with } G_2 \text{ holonomy}$$

$$\text{Quantum theory: } Z_H(\omega) = \int DC \, e^{-\frac{1}{2} \int G \wedge *G}$$

$$G = x + dC$$

Quantum capacity: $\mathcal{H}^4 \ni [\frac{12}{5}]$

Triffing Conclusions: $M \cong X \times \mathbb{R}$

H/He: 1.0
3-4 mag
on $\chi^2 \Rightarrow$

$\frac{\partial \mathcal{L}}{\partial \mathbf{c}} = 0$ variation $\mathbf{I}^T \mathbf{x} = \mathbf{x}$ eigen

$$\gamma^2 = \frac{\gamma^p}{\gamma^p \vee \gamma} = \frac{\gamma^c}{\sigma^c}$$

$$\Rightarrow \phi = \phi + \phi$$

fine independent, grand state
 X is CY.
 $\Rightarrow d'_{\mathcal{C}} = 0, d'_{\mathcal{P}} = 0$ so

Cooper, G. B. : 1952, $\sqrt[3]{V} \propto \sqrt[3]{V_{enc}}$.

[illegible]

$$cp_2 > \int_{\nu^2}^{\nu^2 + 1}$$

H. K. L. S. G. = $\int \rho \, r^2 \, dr$

$$H_{\text{tail}} H = T + V$$

$$\frac{1}{2} \int_{-\infty}^{\infty} \rho^2 \left(-\frac{1}{x} \right) dx$$

B model
complex value

Japan-1

$$\text{max } \frac{1}{2} = \frac{1}{3} \quad \text{min}$$

$$\sqrt[3]{\frac{16}{27}} = \sqrt[3]{\frac{2^4}{3^3}} = \frac{2}{3} \sqrt[3]{2}$$

$$Z_{g,h} = \frac{Z = \alpha x}{\alpha - x} \left(\frac{x}{h^3} - 2 \right) \cdot 1 = 46 Z$$

10-10-20

So this is classical but

- ...-factors: A, B $\mathbb{Z}_{H_1} \Rightarrow \mathbb{Z}_{H_2}^B(x)$ KS when.
- Limit $\bar{g}_k \rightarrow 0$

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