

From Kapustin-Witten
to Extended Bogomolny
to Hitchin

Rafe Mazzeo

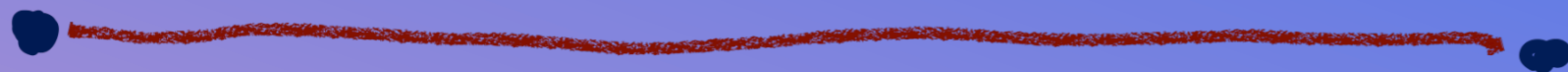
Stanford

Between

Topology

and

QFT



?

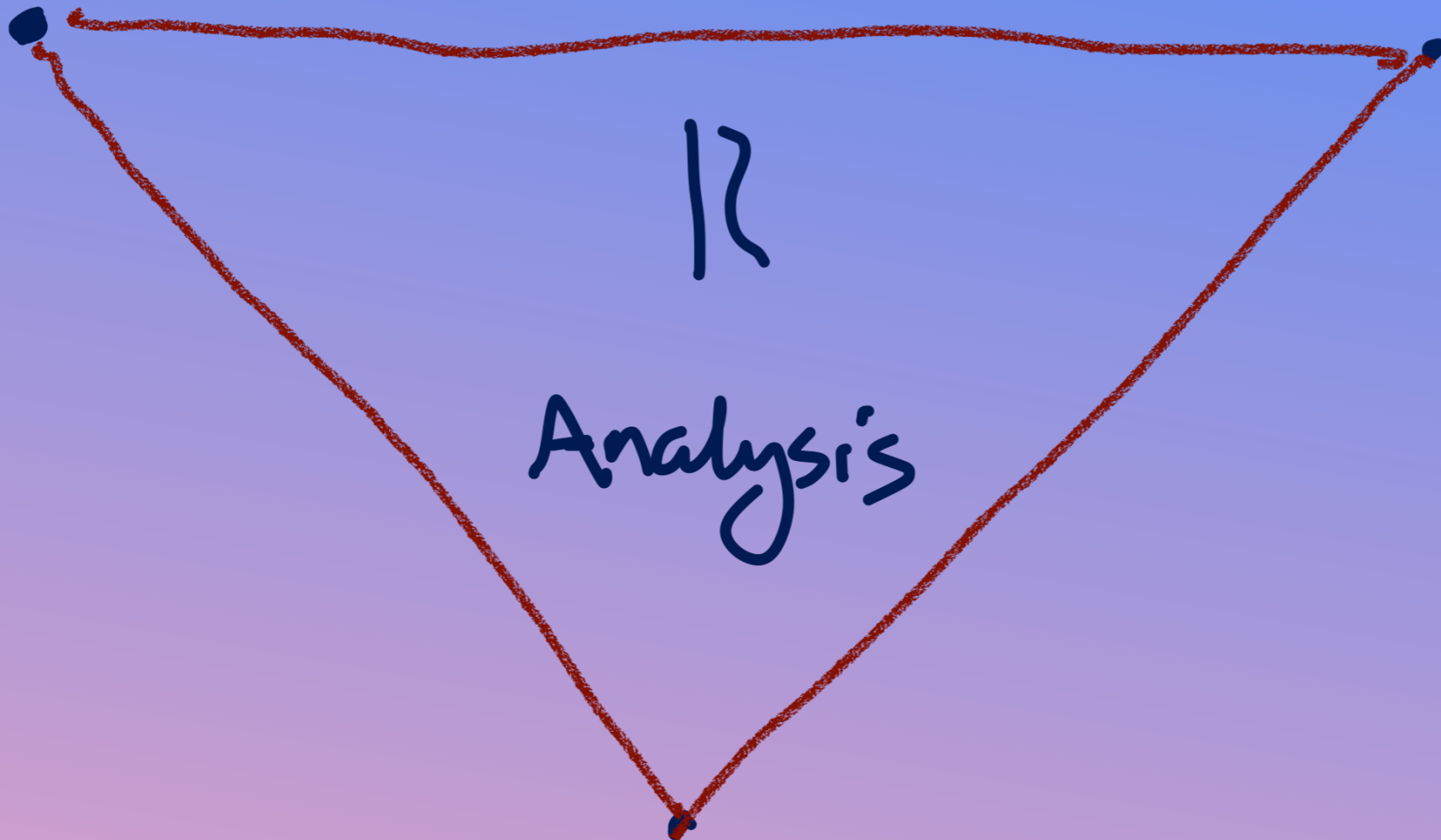
Topology

QFT

\mathbb{R}

Analysis

Geometry



Instantons and Four-Manifolds

Daniel S. Freed
Karen K. Uhlenbeck

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Start with a
"new" gauge theory

Kapustin - Witten
2005

(Geometric Langlands)

(X^4, g) \mathbb{F} G -bundle



X

(think $G = SU(2)$).

$A =$ connection on \mathbb{F}

$\phi \in \Omega^1(X, \text{ad } \mathbb{F})$

$$\left. \begin{aligned} F_A - \phi \wedge \phi + * d_A \phi &= 0 \\ d_A * \phi &= 0 \end{aligned} \right\} \text{KW}_1$$

cf. Gagliardo-Uhlenbeck

Specialization of 5D theory
(Mangstys, Witten)

Embedded in family of
equations

$$\left. \begin{aligned} & \left(F_A - \phi \wedge \phi + t \, d_A \phi \right)^+ = 0 \\ & \left(F_A - \phi \wedge \phi - t^{-1} \, d_A \phi \right)^- = 0 \\ & d_A * \phi = 0 \end{aligned} \right\} \text{KW}_t$$

Observation:

$$\mathcal{A} = A + i\phi \in \mathfrak{g}_{\mathbb{C}}$$

$$F_{\mathcal{A}} = dA + i d\phi + (A + i\phi) \wedge (A + i\phi)$$

$$= (F_A - \phi \wedge \phi) + i d_A \phi = 0$$

Include $d_A * \phi = 0$

Thus $F_A = 0$, $d_A * \phi = 0$

\rightsquigarrow solutions of KW_t for all t

$KW_t = 0$ only uses half of this.

$F_A = 0$ complex gauge invariant

$d_A * \phi = 0$ real moment map

X^4 closed?

$$KW_t(A, \phi) = 0$$

$$\Rightarrow F_A = 0$$

~~~~~> simpler  
topological theory

(PS: Weizenböck formula, ...)

Focus instead on



$X^Y$  with boundary  $W$

← What are the  
right boundary  
conditions?



$$X^4 = w^3 \mathbb{R}_y^+$$

$\mathbb{R}_y^+$

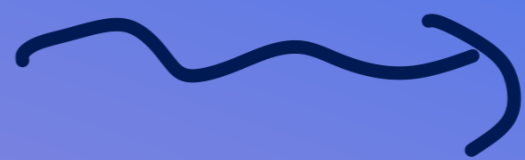
Boundary  
conditions at  $y=0$   
and as  $y \rightarrow \infty$

Demand that solutions become  
translation-invariant in  $y$

as  $y \rightarrow \infty$



$$A_y = 0, \quad \phi_y = 0$$



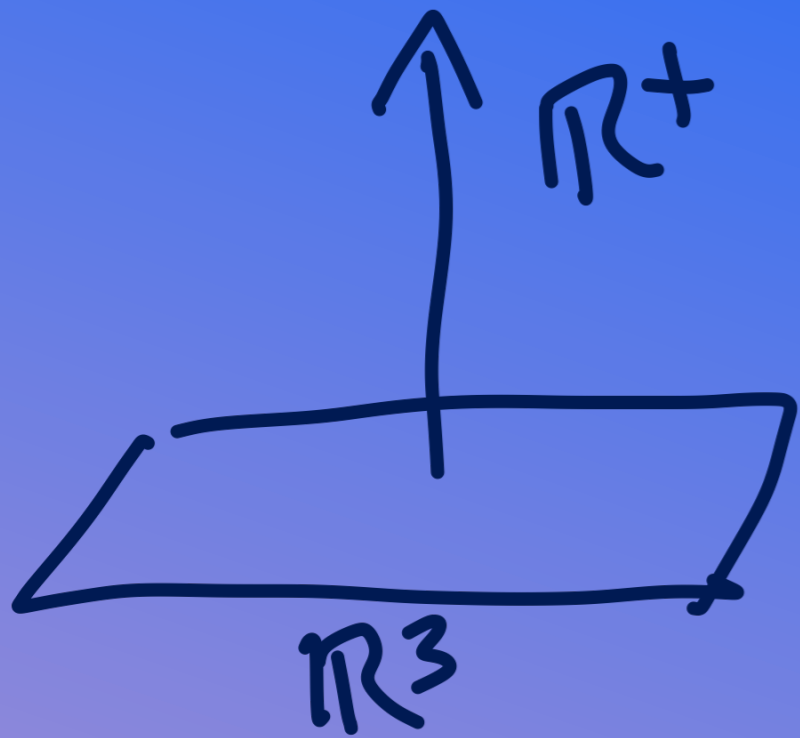
$$F_A - \phi \wedge \phi = 0$$

$$\left. \begin{array}{l} F_A = 0 \\ \end{array} \right\}$$

$$d_A \phi = 0, \quad d_A^* \phi = 0$$

$$\left. \begin{array}{l} d_A^* \phi = 0 \\ \end{array} \right\}$$

Note: asymptotics  $\rightsquigarrow$  decoupling



$$X \in \mathbb{R}^3, y \geq 0$$

Model:  $A_y = 0, \phi_y = 0$   
and can gauge away  
 $A_x$  too

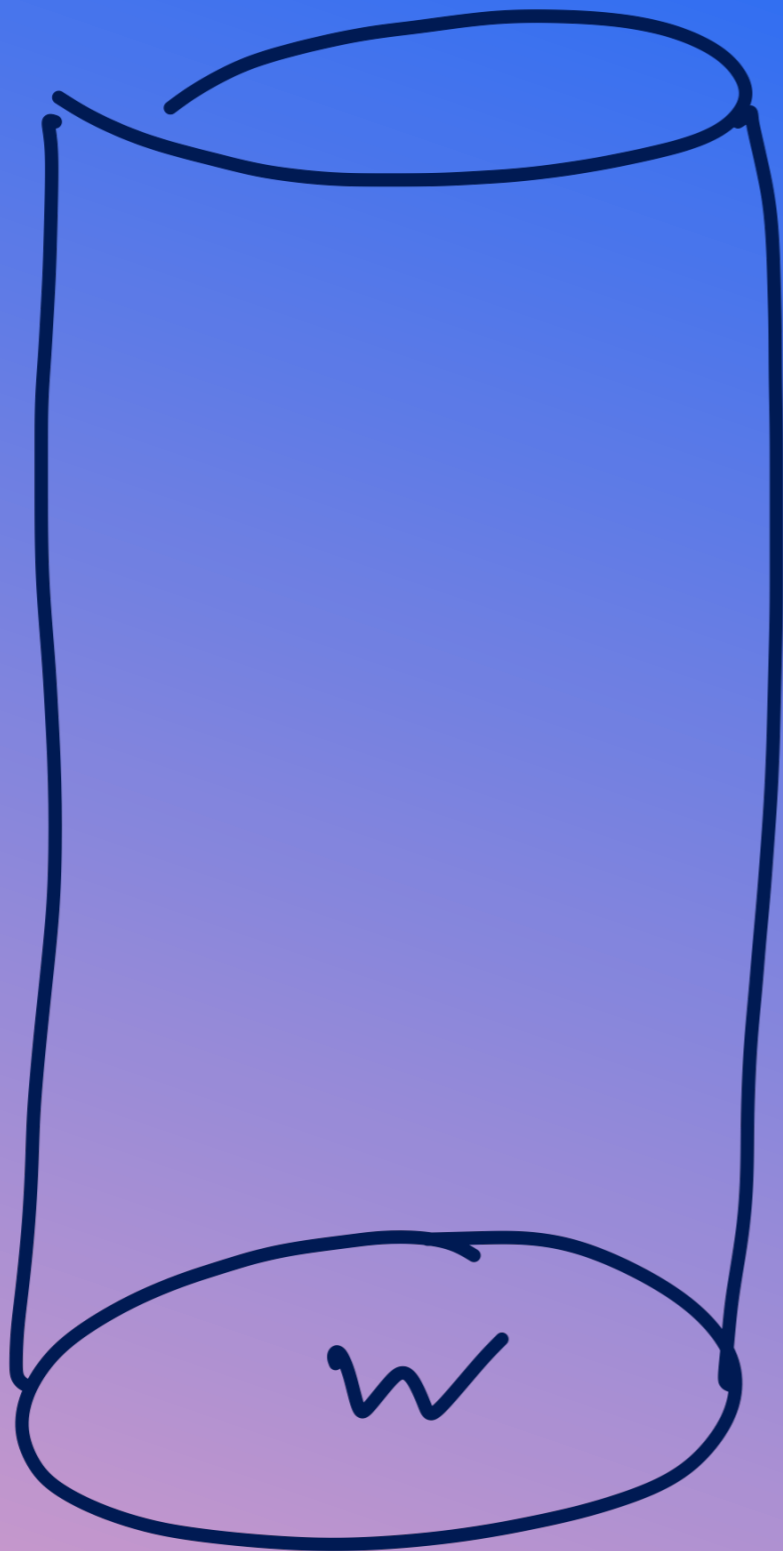
Remaining  
eq'n:

$$d\phi = \phi \wedge \phi$$

$\rightsquigarrow \phi = \sum t_a dx^a$  where

$$[t_a, t_b] = \epsilon_{abc} t_c$$

Altogether



Flat  $SL(n, \mathbb{C})$   
Connection  
(element of  
Hitchin  
moduli  
space).

Nahm condition.



The Nahm boundary condition  
in general:



Choose principle embedding

$$\rho: \mathfrak{su}(2) \rightarrow \mathfrak{g}$$

i.e.  $t_1, t_2, t_3 \in \mathfrak{g}$

$$[t_a, t_b] = \epsilon_{abc} t_c$$

$$\phi_\rho: TW \hookrightarrow \text{ad } E|_W \quad (\text{isom. when } G = \text{SU}_2)$$

$$\phi_\rho = \sum t_i dx^i$$

$$A = A_0 + \mathcal{O}(y^\epsilon), \quad \phi = \frac{1}{y} \phi_\rho + \mathcal{O}(y^{-1+\epsilon})$$

# Regularity (M-Witten)

$$A \sim A_0 + \kappa A_1 + \dots$$

$$\phi \sim \frac{1}{y} \mathcal{P}_g + \sum \mathcal{P}_j y^{\mathcal{A}_j} \quad (\mathcal{A}_j \rightarrow \infty)$$

$A_0 \cong$  Levi-Civita connection

(via intertwining

$$\mathcal{P}_g : TW \hookrightarrow \text{ad}(E)|_W)$$

This relies on analysis  
of linearization:

$$\mathcal{L}_{KW} \sim B_y \partial_y + \sum B_{x_j} \partial_{x_j} + \frac{1}{y} B_0$$

(singular potential)

Also,

$$\text{Ind}(\mathcal{L}_{K^W}) =$$

$$- (\dim \mathfrak{g}) \chi(X)$$

$$\left( = 0 \text{ when } X = \mathbb{R}^+ \times W \right).$$

For  $KW_t$ , there are  
'tilted' Nahm pole conditions

Replace  $t$  in  $KW_t$  by

$$t = \tan\left(\frac{\pi}{4} - \frac{3}{2}\beta\right) \quad \left(\begin{array}{l} t=1 \\ \beta=0 \end{array}\right)$$

$$A_1 - \tan\beta\phi_1 = 0$$

$X_1 =$  distinguished  
direction in boundary

Model  $(G = SU(2))$

$$A_z = \frac{1}{\Delta y} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \sin \beta$$

$$A_{\bar{z}} = \frac{1}{\Delta y} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \sin \beta$$

$$\phi_z = \frac{1}{\Delta y} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \cos \beta$$

$$\phi_{\bar{z}} = \frac{1}{\Delta y} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \cos \beta$$

$$\phi_1 = \frac{i}{2y} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \cos \beta$$

Note

$A, \phi, \phi_1$  all  $\sim \frac{1}{g}$

leading parts of  $A, \phi$

upper/lower triangular

Program:

Study sol<sup>'ns</sup> of  $KW_1 = 0$

(or  $iKW_t = 0$ ) on

$W \times \mathbb{R}^+$  with these

boundary conditions.

New invariants of  $W$ ?

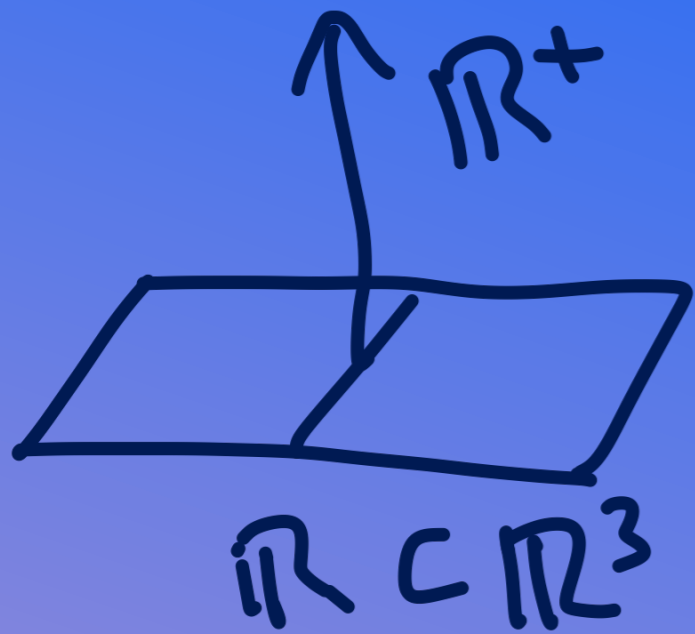
(Gaiotto - Witten)



More generally, consider  
knot  $K \subset W$  at  $y=0$

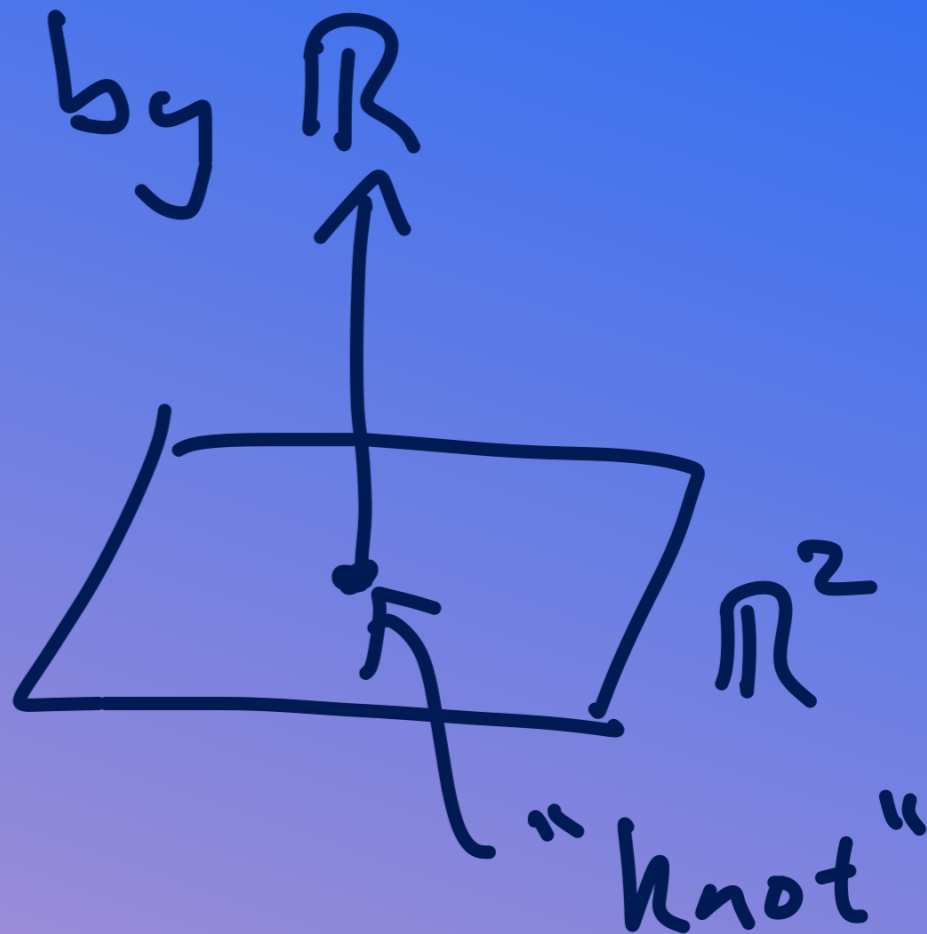


Introduce new  
boundary conditions  
with extra  
singularity at  $K$



model  
knot

mod out by  $\mathbb{R}$



$(\rho, \theta, \psi)$

spherical coords.

$$A = - (n+1) \sin^n \psi \frac{F_n(\psi)}{F_{n+1}(\psi)} d\theta \begin{pmatrix} i/2 & 0 \\ 0 & -i/2 \end{pmatrix}$$

$$f = (\phi_2 - i\phi_3) = \frac{2(n+1)}{f} \frac{\sin^n \psi e^{i\theta}}{F_{n+1}(\psi)} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$$

$$\phi_1 = \frac{n+1}{f} \frac{G_{n+1}(\psi)}{F_{n+1}(\psi)} \begin{pmatrix} i/2 & 0 \\ 0 & -i/2 \end{pmatrix}$$

$$F_\ell(\psi) = (1 + \cos \psi)^\ell - (1 - \cos \psi)^\ell$$

$$G_\ell(\psi) = \quad " \quad + \quad "$$

For other  $G$

model. knot solution

due to Mikhaylov.

---

$t \neq 1$  ?

Similar regularity theory  
 for solutions of  $KW(A, \phi) = 0$   
 with knot singularities.



$$A, \phi \sim$$

$$\sum \rho^{\alpha_j} S^{\mu_l} F_{je}(\theta, x_l)$$

$\rho$  = distance to knot

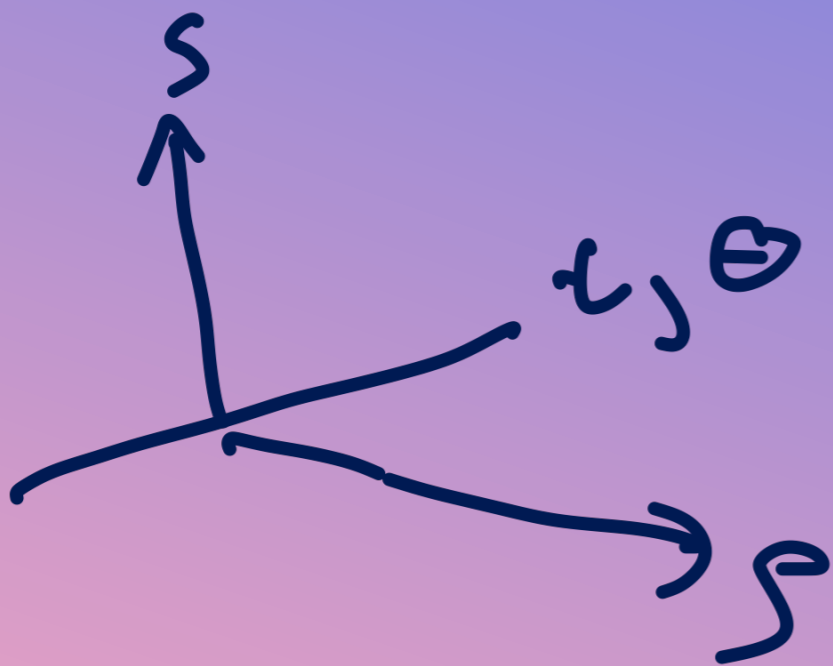
$S$  = spherical coord

$\sim$  dist. to lower boundary

Based on analysis of

$$\mathcal{L}_{KW} \sim \mathcal{B}_f \partial_f + \mathcal{B}_t \partial_t$$

$$\dagger \frac{1}{f} \left( \mathcal{B}_s \partial_s + \mathcal{B}_\Theta \partial_\Theta + \frac{1}{s} \mathcal{B}_D \right)$$



Bigger program (Gaiotto-Witten)

Count of solns  $\rightsquigarrow$

coefficients of Jones  
polynomial of  $K$ .

---

Haydys-Witten equations

$\rightsquigarrow$  categorification

§ Khovanov homology

# Steps to accomplish all of this

- 1) Local properties of moduli space (M-Witten, He)
- 2) Tilted case? ( $t \neq 1$ )
- 3) Existence of solutions with or without knots } today
- 4) Compactness (Taubes)



How to

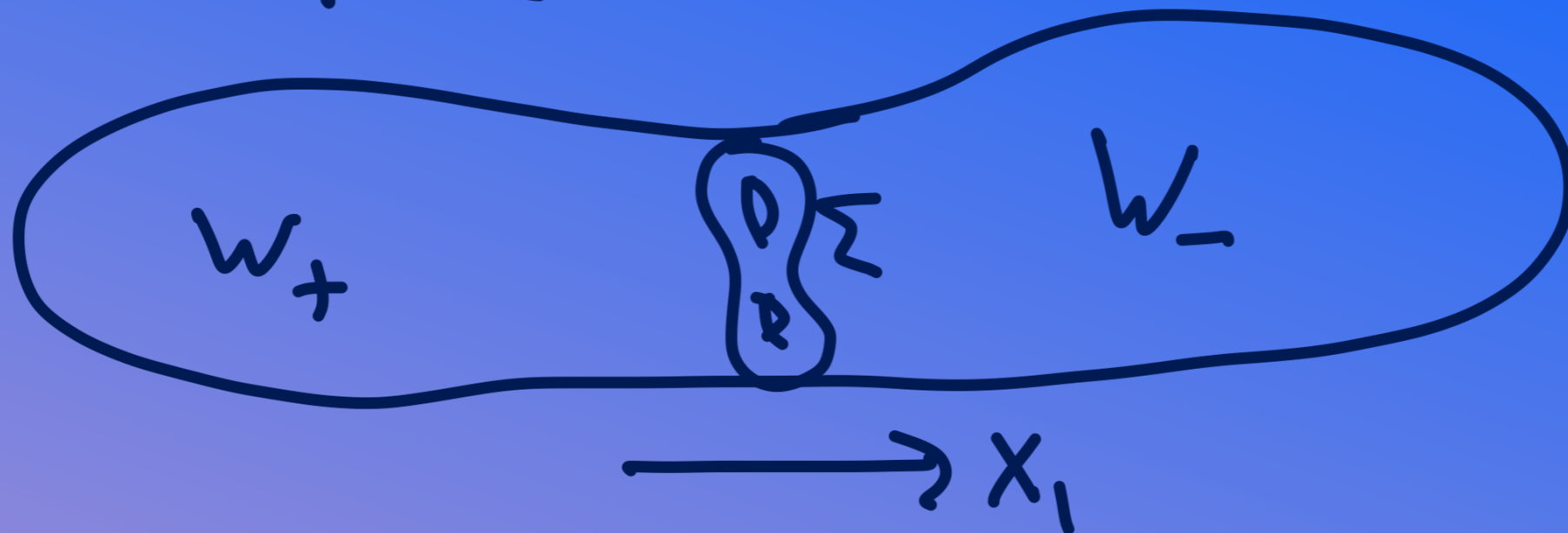
construct solutions?

- dimensionally-reduced cases

- gluing

- ?

$$W = W_+ U_\Sigma W_-$$



chongate along  $\Sigma$

$\rightarrow$  Solns on  $\mathbb{R}_{x_1} \times \Sigma_{x_2, x_3} \times \mathbb{R}_y^+$

Look for  $x_1$ -indep sol'ns

This leads to the  
Extended Bogomolny (EB)  
equations on  $\Sigma \times \mathbb{R}^+$

$$\text{Fields: } \begin{cases} A = A_\Sigma \\ \phi = \phi_\Sigma \\ \phi_1 \end{cases}$$

$$F_A - \phi \wedge \phi = * d_A \phi_1$$

$$d_A \phi = * [\phi, \phi_1]$$

$$d_A^* \phi = 0$$

} EB

Note:  $\left. \begin{array}{l} \phi_1 = 0 \\ y\text{-indep.} \end{array} \right\} \Rightarrow$  Hitchin's eqns

$\phi = 0$  Bogomolny

$x$ -independent,  $\phi_1 = 0 \Rightarrow$  Nahm

Study KW on

$S' \times \Sigma \times \mathbb{R}^+$  (or  $S' \times \Sigma \times \mathbb{R}$ )

Look for  $S'$ -independent  
solutions

Solutions on  $S^1_{x_1} \times \Sigma \times \mathbb{R}_y$

which are independent of

$x_1$  and  $y$  are

flat  $SL(n, \mathbb{C})$  connections on  $\Sigma$

i.e. elements of the

Hitchin moduli space.

# Theorems (He - M.)

Based on conjectures/predictions  
of Gaiotto - Witten

- 1) Solutions with Naim pole  
singularity at  $y=0$  on  $\Sigma$   
(no knots)  $\longleftrightarrow$  Hitchin  
component of  
flat  $SL(n, \mathbb{R})$  connections

1) (cont.)

In other words

$$\{ \text{sols. to EBE} \} / \sim = \mathcal{M}_{\text{EBE}}$$

$$\longleftrightarrow \mathcal{M}_{\text{Hit}}(\text{SL}_n(\mathbb{R}))$$

diff. isomorphism



2) Solutions with knot singularities  
at  $y=0$  on  $\Sigma$   $\longleftrightarrow$

$\{ (\mathcal{E}, \mathcal{F}, L) : (\mathcal{E}, \mathcal{F}) \text{ stable}$

Higgs pair (element of  
Hitchin moduli space)

and  $L \subset \mathcal{E}$  line subbundle }  
knot  $K = \text{divisor on } \Sigma$

determined by tangency of  $L, \mathcal{F}(L), \mathcal{F}^2(L),$   
etc.

3) All solutions on  $S' \times \Sigma \times \mathbb{R}^+$   
with these bdry conditions  
are  $S'$  invariant.

4) Solutions of EBE on  $\Sigma \times \mathbb{R}^+$   
with tilted Nahm pole at  $y=0$

←→ twisted oper

moduli space

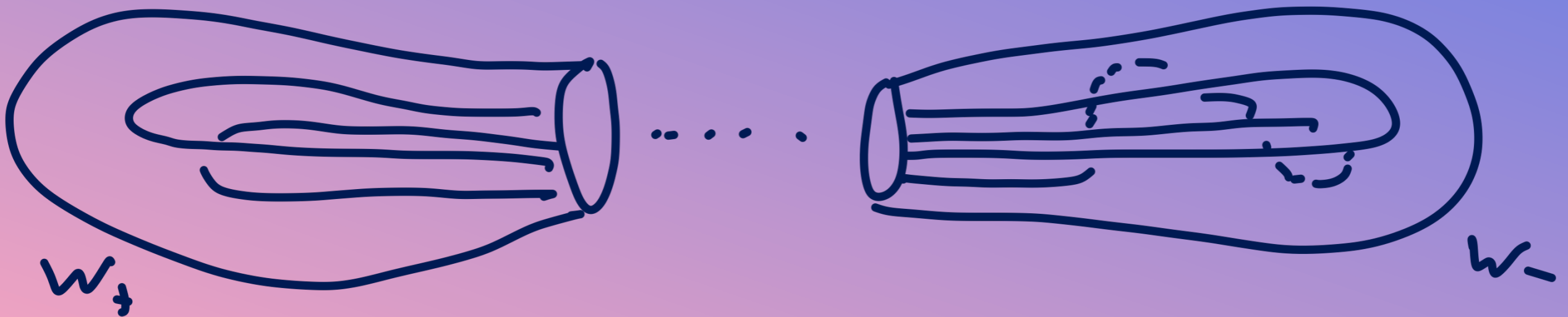
(inside Hitchin moduli  
space).

What's missing?

- tilted Nahm boundary conditions  
with knot singularities  
(in progress).

- more general solutions

$$W \neq S^1 \times \Sigma$$



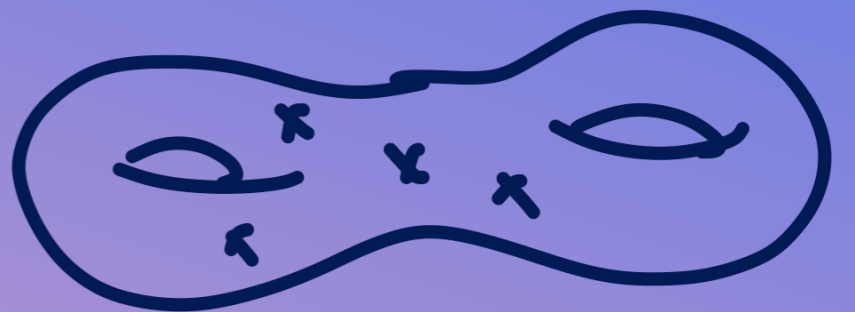
We discuss 2) when  
rank  $E = 2$

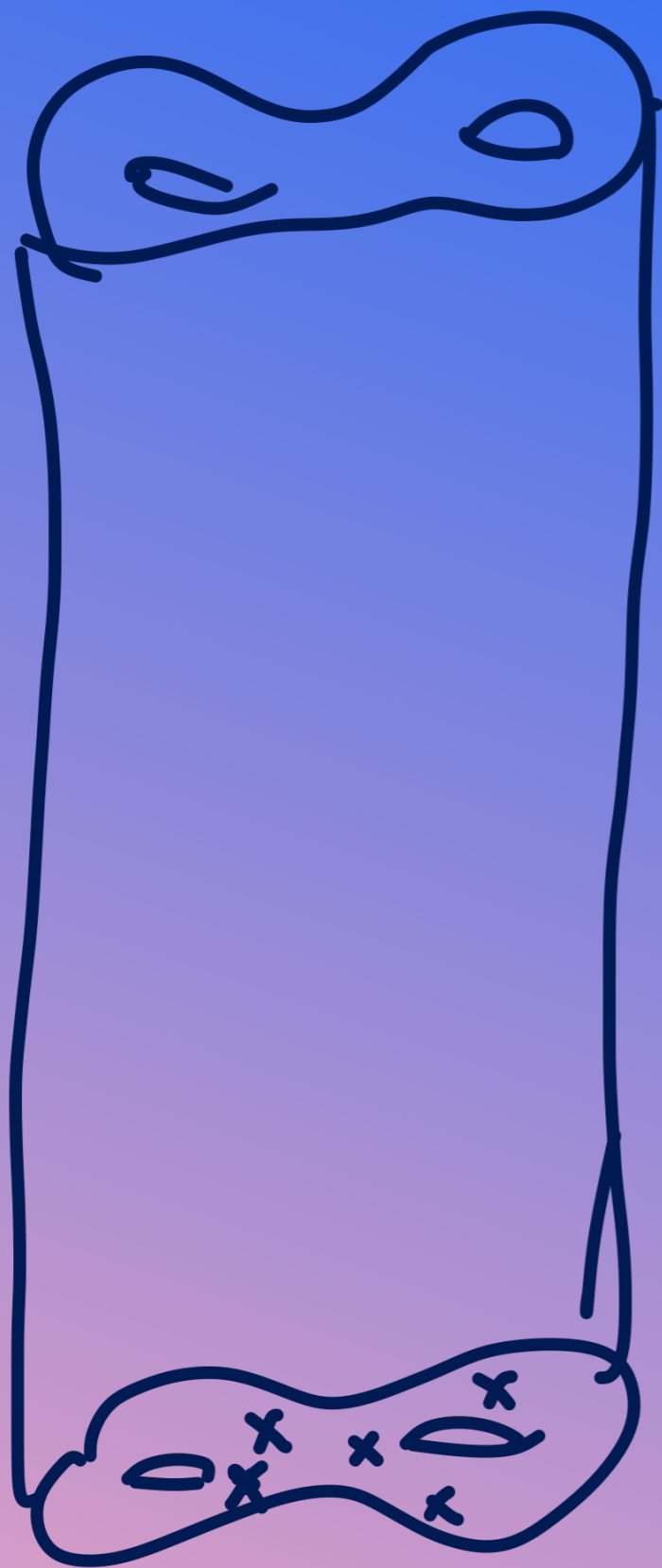
$(E, \varphi) \in \mathcal{M}_{\text{Hitchin}}$ ,  $L \subset E$   
line subbundle

$L \wedge \varphi(L)$  vanishes at

$$\sum n_j p_j$$

The  $p_j$  are the knots





$(E, g, L)$   
at  $y = \infty$

Soln. of  
Extended Bogomolny

knots at  $y = 0$

Formalism:

$$\nabla_A = \nabla_2 dx_2 + \nabla_3 dx_3 + \nabla_y dy$$

$$\phi = \frac{1}{2} (\phi_z dz + \phi_{\bar{z}} d\bar{z})$$

$$\mathcal{D}_1 = (\nabla_2 + i\nabla_3) d\bar{z}$$

$$\mathcal{D}_2 = \text{ad}(\phi_2 - i\phi_3) = [\phi_2 - i\phi_3, \cdot]$$

$$\mathcal{D}_3 = \nabla_y - i\phi_1$$

EBE  $\Leftrightarrow$

$$[D_i, D_j] = 0 \quad i, j = 1, 2, 3$$

and

$$\frac{i}{2} \wedge ([D_1, D_1^\dagger] + [D_2, D_2^\dagger] + [D_3, D_3^\dagger]) = 0$$

$\underbrace{\hspace{15em}}_{\Omega_H}$

# Hermitian Yang-Mills Formalism

Donaldson, Uhlenbeck-Yau

$$[\mathcal{D}_i, \mathcal{D}_j] = 0 \quad \exists c\text{-invariant}$$

remaining equation is  
real moment map



$$H = H_0 e^s \Rightarrow$$

$$\Omega_H = \Omega_{H_0} + r(s) \mathcal{L}_{H_0} s + Q(s)$$

$$r(s) = (e^{ads} - 1) / ads$$

$$\mathcal{L}_{H_0} s = \frac{i}{2} \Lambda (D, D_1^{H_0} + D_2 D_2^{H_0} + D_3 D_3^{H_0})$$

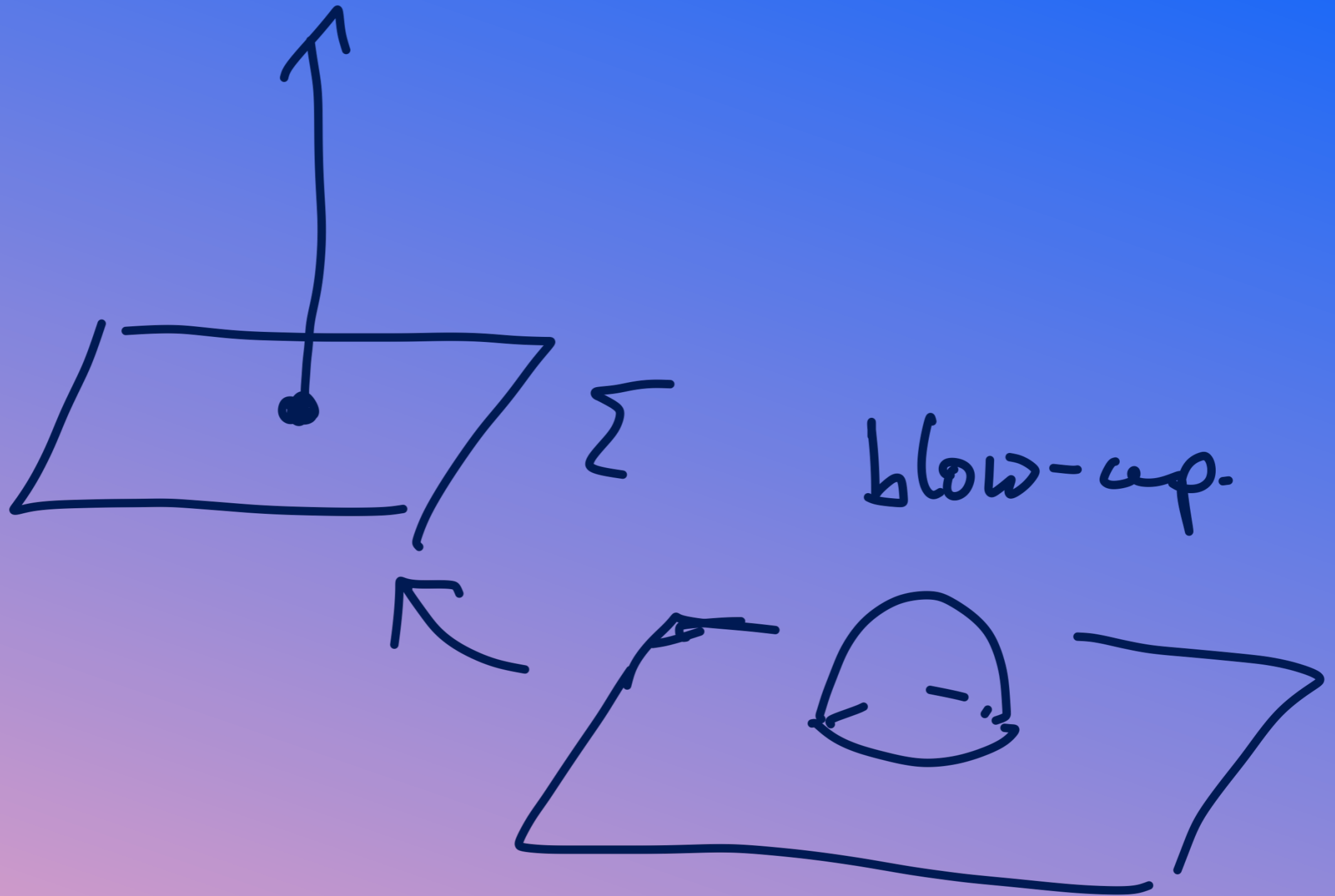
(Laplace-type)

- Pick extremely good initial approximate sol'n.



Solve in Taylor series at these two dimension one boundaries

(and at  $\infty$ )



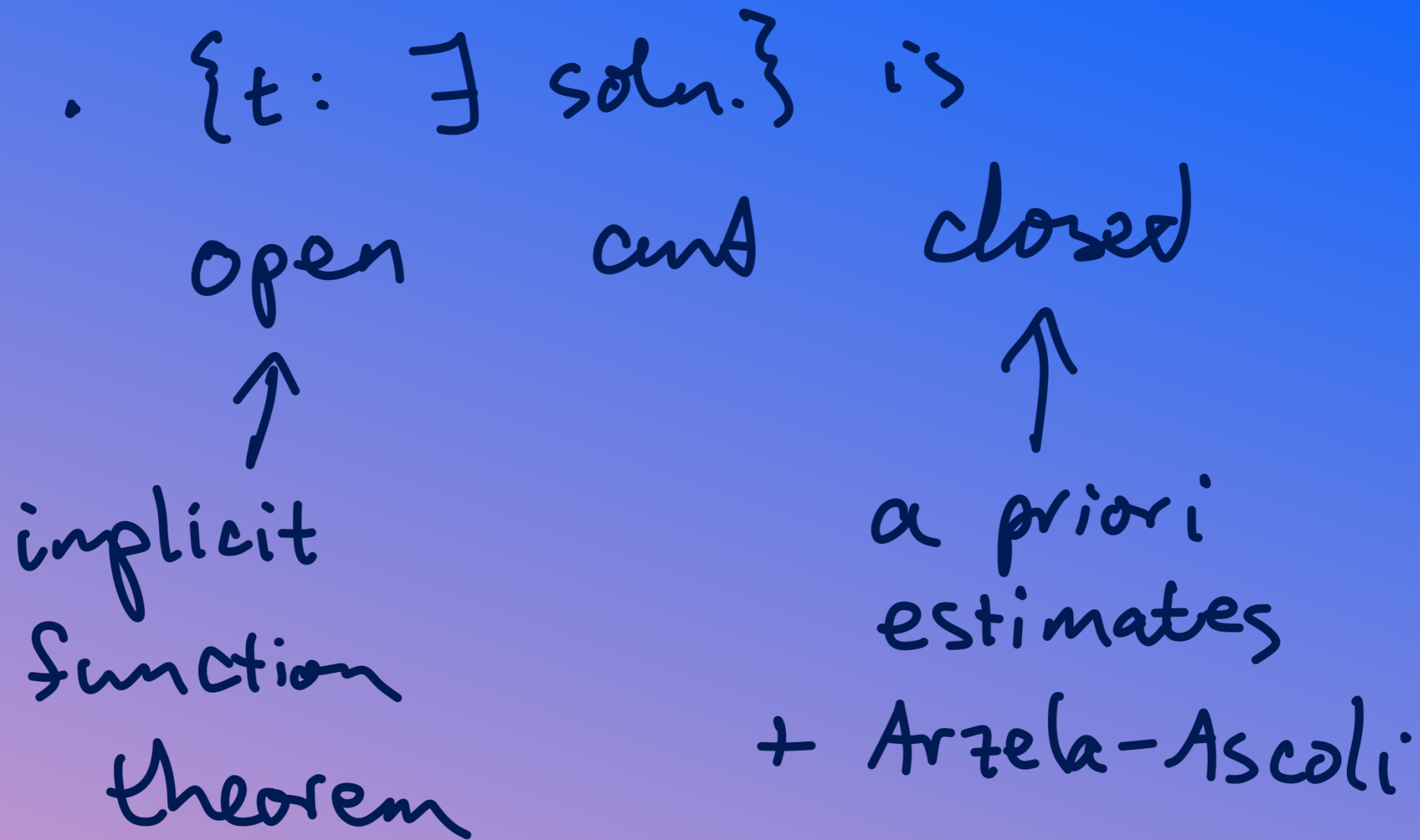
- Then obtain a priori estimates on derivatives of  $s$

- Deform equation

$$\Omega_H^t = 0 \quad (\text{family of equations})$$

$$\Omega_H^0 = 0 \quad \text{has soln. } H = H_0$$

$$\Omega_H^1 = 0 \quad \text{is what we want.}$$



(all the work is buried here)

Back to original problem:

Counting solutions of KW

(or just flat  $SL(2, \mathbb{C})$  connections)

on  $X^4$  or  $W^3 \times \mathbb{R}^+$

Taubes:

If  $(A_j, \phi_j)$  is

a sequence of solutions,

then ...

Either  $(A_j, \phi_j)$  converges

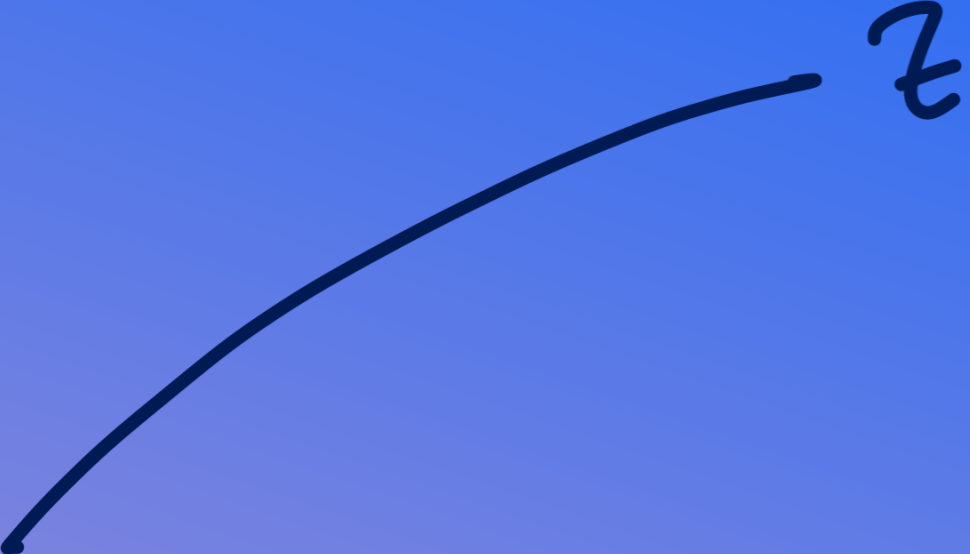
up to unitary gauge

OR a rescaled sequence of

the  $\phi_j$  converge to a

" $\mathbb{Z}_2$ -harmonic spinor"



  $Z$   $L = \mathbb{R}$ -line bundle  
on  $X \setminus Z$  with  $\mathbb{Z}_2$   
holonomy.

$\omega$   $L$ -valued 1-form

$$(d + d^*)\omega = 0$$

$$|\omega| \leq C \text{dist}(\cdot, Z)^{1/2}.$$

Very Overdetermined.



In fact, given any (smooth)  
 $Z$ , (codim 2 in  $X^n$ )  
and  $L \rightarrow X \setminus Z$ , can find

$\omega$  such that  $(d+d^*)\omega = 0$

$$\omega \sim \sum_{j=0}^{\infty} r^{(j-1)/2} \omega_j$$

$$r = \text{dist}(\cdot, Z)$$

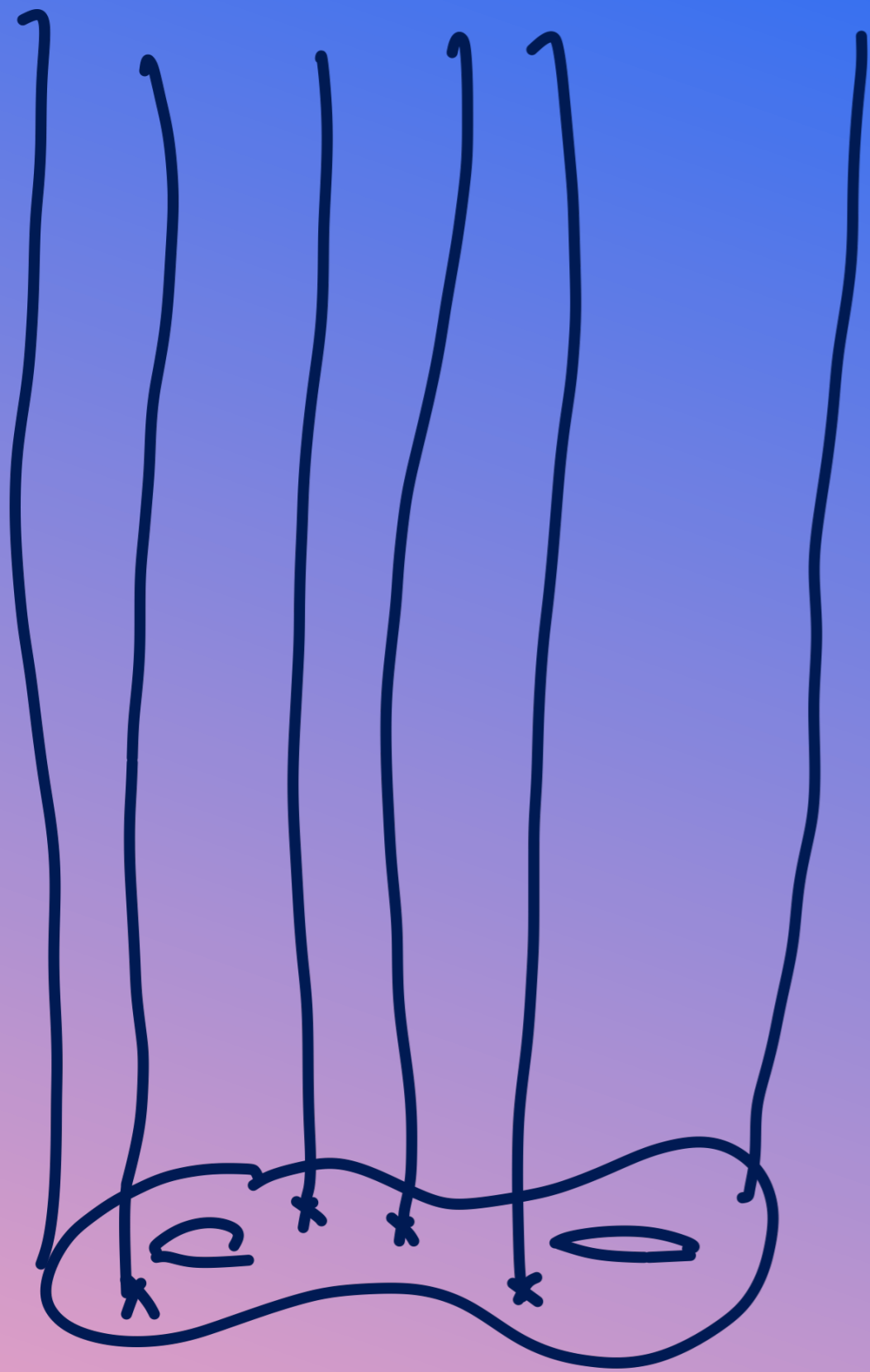
We require  $\omega_0 = 0$  (prescribed  
Cauchy data)

B. Zhang (rectifiability)

Doan - Walpuski (abstract  
existence theorem)

He - M. (concrete model  
of this degeneration).





diverging sequence  
in  $\mathcal{M}_{\text{Hitchin}}$

(well-understood)

$\rightsquigarrow$  singularities at  
 $p_j \times \mathbb{R}^+$

$(p_1, \dots, p_{4g-4})$

zeros of holom.

quadratic differential.

What is the geometry  
of the degeneration  
locus in general  $X$ ?

These correspond to  
the virtual KW solns  
at  $\infty$  ( $KW_t$  comes into play.)

Happy

Birthday

Dan !!!