

Preliminaries and general tools

LP decomposition: Let $\chi \in C_0^\infty(B_1(0))$ be such that $\chi|_{\frac{B_1(0)}{B_{1/2}(0)}} = 1$ and define $\varphi(\xi) := \chi(\xi/2) - \chi(\xi)$ and denote $\lambda_q := 2^q$. We obtain a partition of unity:

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and denote $\lambda_q := 2^q$. We obtain a partition of unity:

$$\chi(\xi) + \sum_{q \geq 0} \varphi(\lambda_q^{-1} \xi) = 1.$$

Furthermore, $|p-q| \geq 2 \Rightarrow \text{supp } \varphi(\lambda_q^{-1} \cdot) \cap \text{supp } \varphi(\lambda_p^{-1} \cdot) = \emptyset$.

Let the FT be def. as $\mathcal{F}u := \int e^{ikx} u(x) dx$ and introduce the LP projections

$$\Delta_q u := \begin{cases} \mathcal{F}^{-1}(\varphi(\lambda_q^{-1} \cdot) \mathcal{F}u) = \lambda_q^3 h(\lambda_q \cdot) * u & , q \geq 0 \\ \mathcal{F}^{-1}(\chi \mathcal{F}u) = h^0 * u & , q = -1 \end{cases}$$

where we define $h := \mathcal{F}^{-1} \varphi$, $h^0 := \mathcal{F}^{-1} \chi$. The LP projection up to the Q^{th} shell is given by

$$S_Q u := \sum_{q=-1}^Q \Delta_q u = \mathcal{F}^{-1}(\chi(\lambda_{Q+1}^{-1} \cdot) \mathcal{F}u) = \lambda_{Q+1}^3 h^0(\lambda_{Q+1} \cdot) * u.$$

Note that we have

$$u = \lim_{Q \rightarrow \infty} S_Q u \text{ a.e.}$$

Def. Besov space: Let $s \in \mathbb{R}$, $1 \leq p, r \leq \infty$. Then we define the inhomogeneous Besov space

$$B_{p,r}^s := \left\{ u \in \mathcal{F}' \mid \|u\|_{B_{p,r}^s} := \left\| \left(\lambda_q^s \|\Delta_q u\|_{L^p} \right)_{q \geq 0} \right\|_{\ell^r} \right\}$$

also fractional derivatives:

Useful inequalities: (i) $b \geq a \geq 1, \alpha \in \mathbb{N}_0^3$: $\|D^\alpha \Delta_q u\|_b \leq \lambda_q^{|\alpha| + d(\frac{1}{a} - \frac{1}{b})} \|\Delta_q u\|_a$ (Bernstein ineq. (BI))

(ii) $\|u\|_p \sim \left\| \left(\|\Delta_q u\|_{L^p} \right)_{q \geq -1} \right\|_{\ell^2}$ (Littlewood-Paley ineq. (LPI))

$$\sim \|D^\alpha \Delta_q u\|_a$$

$$(iii) \int_{S_2} \left(\int_{S_1} |F(x,y)|^q \mu_1(dx) \right)^{p/q} \mu_2(dy) \leq \int_{S_1} \left(\int_{S_2} |F(x,y)|^p \mu_2(dy) \right)^{q/p} \mu_1(dx)$$

In addition, we have the following continuous embeddings for $b \geq a \geq 1$
 in the sense of norm-estimates

$$B_{a,r}^s \subseteq B_{b,r}^{s-d(\frac{1}{a}-\frac{1}{b})}$$

$$B_{a,2}^0 \subseteq L^a \quad \text{for } a \geq 2$$

and furthermore

$$H^{5/6}(\mathbb{R}^3) \subseteq B_{\frac{4}{3},2}^{\frac{2}{3}}(\mathbb{R}^3) \subseteq B_{\frac{15}{2},2}^{\frac{1}{2}}(\mathbb{R}^3) \subseteq B_{3,2}^{\frac{1}{3}}(\mathbb{R}^3).$$

Def. (Weak solution)

u is called a weak solution of the incompressible Euler equation w/ initial data $u_0 \in L^2(\mathbb{R}^3)$ iff

- $u \in C_w([0,T]; L^2(\mathbb{R}^3))$

- $\forall \varphi \in C^1([0,T]; \mathcal{S}(\mathbb{R}^3))$ w/ $\nabla_x \cdot \varphi = 0 \quad \forall 0 \leq t \leq T$

$$(u(t), \varphi(t))_2 - (u(0), \varphi(0))_2 - \int_0^t (u(s), \partial_s \varphi(s))_2 ds = \int_0^t \int_{\mathbb{R}^3} [(u(s) \cdot \nabla) \varphi(s)] \cdot u(s) dx ds$$

- $\nabla_x \cdot u(t) = 0$ in the distributional sense for all $0 \leq t \leq T$.

Estimates on energy flux

Define

$$\mathcal{T}_Q(t) := \frac{1}{2} \frac{d}{dt} \|S_Q u(t)\|_2^2 \stackrel{\text{Assume we can interchange}}{=} \int \text{Tr} [S_Q(u \otimes u) \cdot \nabla S_Q u] dx,$$

$$\widehat{f}^{-1} \varphi(\omega) := (2\pi)^{-3} \int_{\mathbb{R}^3} e^{-ikx} \varphi(k) dk \quad (\text{FT})$$

Then we have

$$\mathcal{T}_Q(t) = \int \text{Tr} [r_Q(u, u) \cdot \nabla S_Q u] dx + \int \text{Tr} [(u - S_Q u) \otimes (u - S_Q u) \cdot \nabla S_Q u] dx$$

w/ $r_Q(u, u) = \int h_Q^\nu(y) (u(x-y) - u(x)) \otimes (u(x-y) - u(x)) dy$

and $h_Q^\nu(y) := \lambda_{Q+1}^3 h^\nu(\lambda_{Q+1} y)$ Used $\int h_Q^\nu(y) dy = 1$

Note that we used

$$S_Q(u \otimes u) = r_Q(u, u) + (u - S_Q u) \otimes (u - S_Q u) + (S_Q u)^{\otimes 2}$$

together w/

$$\int \text{Tr} [S_Q u \otimes S_Q u \cdot \nabla S_Q u] dx = \int S_{Qj} u_j S_{Qj} u_j \partial_j S_{Qi} u_i = -\frac{1}{2} \int |S_Q u|^2 S_Q(\nabla \cdot u) = 0$$

where for $S_Q v \in L^2$ if $\widehat{S_Q v} \in L^2$, we conclude

$$\int S_Q \partial_j v_j S_Q u_i = \int \widehat{S_Q \partial_j v_j} \widehat{S_Q u_i} = - \int \partial_j v_j S_Q \partial_i u_i$$

We continue w/

$$|\mathcal{T}_Q| \leq \|r_Q(u, u)\|_{3/2} \|\nabla S_Q u\|_3 + \|u - S_Q u\|_3^2 \|\nabla S_Q u\|_3$$

$$\leq \left[\int h_Q^\nu(y) \|u(\cdot - y) - u\|_3^2 dy + \|u - S_Q u\|_3^2 \right] \|\nabla S_Q u\|_3$$

Have

$$\|u(\cdot - y) - u\|_3^2 \stackrel{L^1}{\leq} \left\| \left(\sum_q |\Delta_q u(\cdot - y) - u|^2 \right)^{1/2} \right\|_3^2 \stackrel{M1}{\leq} |y|^2 \sum_{q \leq Q} \|\nabla \Delta_q u\|_3^2 + \sum_{q > Q} \|\Delta_q u\|_3^2$$

$$\stackrel{BL}{\leq} |y|^2 \sum_{q \leq Q} \lambda_q^2 \|\Delta_q u\|_3^2 + \sum_{q > Q} \|\Delta_q u\|_3^2$$

Furthermore, we have

$$\begin{aligned} \|\nabla_{S_Q} u\|_3^2 &\lesssim \left\| \left(\sum_{k \leq Q} |\Delta_k S_Q \nabla u|^2 \right)^{1/2} \right\|_3^2 \lesssim \sum_{k \leq Q+1} \left\| \sum_{q=k-1}^{k+1} \nabla \Delta_k \Delta_q u \right\|_3^2 \\ &\stackrel{SI}{\lesssim} \sum_{q \leq Q} \sum_{q'-q-2}^{(q+2) \wedge Q} \lambda_q \|\Delta_q u\|_3 \lambda_{q'} \|\Delta_{q'} u\|_3 \\ &\leq \sum_{q \leq Q} \lambda_q^2 \|\Delta_q u\|_3^2 \end{aligned}$$

$\leq \sum_{q-1 \leq k \leq q+1} \sum_{q-2 \leq q' \leq q+2}$
 3 cont. term.

In addition,

$$\|u - S_Q u\|_3^2 = \left\| \sum_{q > Q} \Delta_q u \right\|_3^2 \lesssim \sum_{q > Q} \|\Delta_q u\|_3^2$$

(anal. to above)

Using $\int |y|^{2\alpha} |y| dy \approx \lambda_Q^{-2}$, we end up w/

$$\begin{aligned} |\Pi_Q| &\lesssim \left(\lambda_Q^{-2} \sum_{q \leq Q} \lambda_q^2 \|\Delta_q u\|_3^2 + \sum_{q > Q} \|\Delta_q u\|_3^2 \right) \left(\sum_{q \leq Q} \lambda_q^2 \|\Delta_q u\|_3^2 \right)^{1/2} \\ &= \left(\lambda_Q^{-2} \sum_{q \leq Q} \lambda_q^{4/3} d_q^2 + \sum_{q > Q} \lambda_q^{-2/3} d_q^2 \right) \left(\sum_{q \leq Q} \lambda_q^{4/3} d_q^2 \right)^{1/2} \\ &= \left(\lambda_Q^{-4/3} \sum_{q \leq Q} \lambda_q^{4/3} d_q^2 + \lambda_Q^{2/3} \sum_{q > Q} \lambda_q^{-2/3} d_q^2 \right) \left(\sum_{q \leq Q} \lambda_q^{4/3} d_q^2 \lambda_Q^{-4/3} \right)^{1/2} \end{aligned}$$

where we introduced $d_q := \lambda_q^{1/3} \|\Delta_q u\|_3$.

In here we used

$$\sum_{q > Q} \lambda_q^{-2/3} d_q^2 \leq \sup_{q > Q} d_q^2 \sum_{q > Q} \lambda_q^{-2/3} = \sup_{q > Q} d_q^2 \frac{\lambda_{Q+1}^{-2/3}}{1 - 2^{-2/3}} \lesssim \lambda_Q^{-2/3} \sup_{q > Q} d_q^2$$

Lemma: Let $b > 1$, $(a_q)_{q \geq -1} \in C_0$. Then one has

$$b^{-Q} \sum_{q=-1}^Q b^q a_q \xrightarrow{Q \rightarrow \infty} 0.$$

Proof: Let $\varepsilon > 0$. Choose $Q_0 \in \mathbb{N}$ s.t. $\forall q \geq Q_0$ $|a_q| \leq \frac{\varepsilon(b-1)}{b}$. Then

$$\begin{aligned} \left| b^{-Q} \sum_{q=-1}^Q b^q a_q \right| &\leq b^{-Q} \left[\sum_{q=-1}^{Q_0-1} b^q |a_q| + \sum_{q=Q_0}^Q b^q |a_q| \right] \leq b^{-Q} \left[\underbrace{\sum_{q=-1}^{Q_0-1} b^q |a_q|}_{=: C_{Q_0}} + \frac{\varepsilon(b-1)}{b} \frac{b^{Q+1} - b^{Q_0}}{b-1} \right] \\ &\leq C_{Q_0} b^{-Q} + \varepsilon \end{aligned}$$

$$\Rightarrow \limsup_{Q \rightarrow \infty} \left| b^{-Q} \sum_{q=-1}^Q b^q a_q \right| \leq \varepsilon.$$

□

As a consequence, if we assume $d \in C_0$ ($\Leftrightarrow u \in B_{3, \infty}^{1/2}$), we obtain

$$\lim_{Q \rightarrow \infty} |\Pi_Q| \leq \lim_{Q \rightarrow \infty} \left(\lambda_Q^{-\frac{4}{3}} \sum_{q \leq Q} \lambda_q^{2/3} d_q^2 + \lambda_Q^{-\frac{2}{3}} \lambda_Q \sup_{q > Q} d_q^2 \right) \left(\sum_{q \leq Q} \lambda_q^{4/3} d_q^2 \lambda_Q^{-\frac{4}{3}} \right)^{1/2} = 0.$$

General comment on sufficed Besov regularity:

Note that we find for any multi-index $\alpha \in \mathbb{N}_0^3$

$$\begin{aligned} \left| \Pi_Q \left[\underbrace{S_Q(u \otimes u)}_{\text{not precisely stated}} \cdot \underbrace{D_Q^\alpha u}_{\text{is fine above}} \right] \right| &\leq \left(\lambda_Q^{-2} \sum_{q \leq Q} \lambda_q^2 \|\Delta_q u\|_3^2 + \sum_{q > Q} \|\Delta_q u\|_3^2 \right) \left(\sum_{q \leq Q} \lambda_q^{2|\alpha|} \|\Delta_q u\|_3^2 \right)^{1/2} \\ &= \left(\lambda_Q^{-2+|\alpha|-\beta} \sum_{q \leq Q} \lambda_q^{2-2\beta} \eta_q^{2|\alpha|} + \sum_{q > Q} \lambda_q^{-2\beta} \eta_q^{2|\alpha|} \right) \left(\lambda_Q^{2\beta-2|\alpha|} \sum_{q \leq Q} \lambda_q^{2|\alpha|-2\beta} \eta_q^2 \right)^{1/2} \end{aligned}$$

where $\beta \in \mathbb{R}$ is to be determined and $\eta_q := \lambda_q^\beta \|\Delta_q u\|_3$.

In order to apply the above lemma, we impose

$$\begin{cases} 2-|\alpha|+\beta = 2-2\beta \\ -|\alpha|+\beta = -2\beta \end{cases} \Leftrightarrow \beta = \frac{|\alpha|}{3}.$$

This justifies the above calculation.

Alternative auxiliary result:

Lemma 2: Let $H \in \ell^1(\mathbb{Z})$, $(a_q)_{q \geq 1} \in \ell^{\infty}$ be an arbitrary sequence. Then we have

$$\limsup_{Q \rightarrow \infty} |H * a(Q)| \leq \|H\|_1 \cdot \limsup_{Q \rightarrow \infty} |a(Q)|$$

Proof: Let $\varepsilon > 0$. Fix $M \in \mathbb{N}$ s.t.

$$\sum_{|q| \geq M} |H(q)| < \frac{\varepsilon}{2(\|H\|_1 + 1)}$$

$$\forall q \geq M: |a_q| < \limsup_{k \rightarrow \infty} |a_k| + \frac{\varepsilon}{2(\|H\|_1 + 1)} \quad \left. \vphantom{\sum_{|q| \geq M} |H(q)|} \right\} (*)$$

Let $Q \geq 2M$. We obtain

$$|H * a(Q)| \leq \sum_{q \leq M} |H(Q-q)| |a_q| + \sum_{q > M} |H(Q-q)| |a_q|$$

$$\leq \|H\|_1 \cdot \limsup_{k \rightarrow \infty} |a_k| + \varepsilon.$$

$\varepsilon > 0$ yields the result. □

In particular, if we define

$$H(q) := \begin{cases} \lambda_q^{-4/3} & , q \geq 0 \\ \lambda_q^{2/3} & , q < 0 \end{cases}$$

we obtain by the above estimates

$$|\Pi_Q| \leq (H * d_q^2(Q))^{3/2}$$

w/ $H \in \ell^1(\mathbb{Z})$. Thus the above lemma is applicable and yields

$$\limsup_{Q \rightarrow \infty} |\Pi_Q| \leq \limsup_{Q \rightarrow \infty} d_q^3 = 0.$$

Energy flux between shells

Define

$$\Pi_{Q_0, Q_1} := \frac{1}{2} \frac{d}{dt} \|S_{Q_0, Q_1} u\|_2^2 = \int \text{Tr} [S_{Q_0, Q_1} (u \otimes u) \cdot \nabla S_{Q_0, Q_1} u] dx$$

w/ $S_{Q_0, Q_1} := S_{Q_1} - S_{Q_0}$

Formal computation:

$$\begin{aligned} S_{Q_0, Q_1}^2 &= (S_{Q_1} - S_{Q_0})^2 = S_{Q_1}^2 + S_{Q_0}^2 - 2 S_{Q_0} S_{Q_1} \text{ since } \Delta_{Q_0} \Delta_{Q_1} = 0 \text{ if } |Q_0 - Q_1| \geq 1 \\ &= S_{Q_1}^2 - S_{Q_0}^2 - 2 S_{Q_0} (S_{Q_1} - S_{Q_0}) = S_{Q_1}^2 - S_{Q_0}^2 - 2 \Delta_{Q_0} \Delta_{Q_1} \end{aligned}$$

projections commute

This leads to

$$\Pi_{Q_0, Q_1} = \Pi_{Q_1} - \Pi_{Q_0} - 2 \int_{\mathbb{R}^3} \text{Tr} [\Delta_{Q_0} (u \otimes u) \cdot \nabla \Delta_{Q_0} u] dx$$

w/

$$\Delta_{Q_0} v(x) := \int \bar{h}_{Q_0}(y) v(x-y) dy, \quad \bar{h}_{Q_0} := \mathcal{F}^{-1} [\varphi(\lambda_{Q_0}^{-1}) \varphi(\lambda_{Q_0}^{-1})]$$

As before, we decompose the last term on which we do not already have an upper bound using

$$\Delta_{Q_0} (u \otimes u) \stackrel{H^1}{=} \underbrace{\int \bar{h}_{Q_0}(y) (u(x-y) \otimes u(x-y)) dy}_{=: \bar{r}_{Q_0}(u \otimes u)} + \Delta_{Q_0} u \otimes u(x) + u(x) \otimes \Delta_{Q_0} u(x) - u(x) \otimes u(x)$$

This leads to

$$\begin{aligned} \int \text{Tr} [\Delta_{Q_0} (u \otimes u) \cdot \nabla \Delta_{Q_0} u] dx &= \int \text{Tr} [\bar{r}_{Q_0}(u \otimes u) \cdot \nabla \Delta_{Q_0} u] dx + \int \text{Tr} [\Delta_{Q_0} u \otimes u \otimes u \cdot \nabla \Delta_{Q_0} u] dx + \int \text{Tr} [u \otimes \Delta_{Q_0} u \cdot \nabla \Delta_{Q_0} u] dx \\ &\quad - \int \text{Tr} [u \otimes u \cdot \nabla \Delta_{Q_0} u] dx \end{aligned}$$

see below

Have $|I| \leq \|\bar{r}_{Q_0}(u \otimes u)\|_{3/2} \|\nabla \Delta_{Q_0} u\|_3$

where as above

$$\begin{aligned} \|\bar{r}_{Q_0}(u \otimes u)\|_{3/2} &\leq \int |\bar{h}_{Q_0}(y)| \|u(x-y) \otimes u(x-y)\|_3^2 dy \lesssim \int |\bar{h}_{Q_0}(y)| |y|^2 dy \sum_{q \leq Q_0} \lambda_q^2 \|u\|_3^2 + \sum_{q > Q_0} \|u\|_3^2 \\ &\lesssim \lambda_{Q_0}^2 \sum_{q \leq Q_0} \|u\|_3^2 + \sum_{q > Q_0} \|u\|_3^2 \\ \|\nabla \Delta_{Q_0} u\|_3^2 &\lesssim \sum_{k=Q_0-1}^{Q_0} \|\Delta_k \nabla \Delta_{Q_0} u\|_3^2 \lesssim \lambda_{Q_0-1}^2 \|\Delta_{Q_0-1} u\|_3^2 + \lambda_{Q_0}^2 \|\Delta_{Q_0} u\|_3^2 \end{aligned}$$

$$\Rightarrow |I| \leq \left(\lambda_{Q_0}^{-\frac{4}{3}} \sum_{q \in Q_0} \lambda_q^{\frac{4}{3}} d_q^2 + \lambda_{Q_0}^{-\frac{2}{3}} \sum_{q > Q_0} \lambda_q^{-\frac{2}{3}} d_q^2 \right) \left(d_{Q_{-1}}^2 + d_{Q_0}^2 \right)^{\frac{1}{2}}$$

$\leq \lambda_{Q_0}^{-\frac{4}{3}} \sum_{q \in Q_0} \lambda_q^{\frac{4}{3}} d_q^2$ due to $\lambda_{Q_{-1}-q}^{-\frac{4}{3}} \leq \lambda_{Q_0-q}^{-\frac{2}{3}}$

$$|II| \leq \|\bar{\Delta}_{Q_0} u\|_{\dot{B}_{Q_0}^{\frac{3}{2}}} \|\nabla \Delta_{Q_0} u\|_3 \leq \|\bar{\Delta}_{Q_0} u\|_3 \|\dot{S}_{Q_0} u\|_3 \|\nabla \Delta_{Q_0} u\|_3 \leq \left(d_{Q_{-1}}^2 + d_{Q_0}^2 \right) \left(\lambda_{Q_0}^{-\frac{2}{3}} \sum_{q \in Q_0} \lambda_q^{\frac{2}{3}} d_q^2 \right)$$

$$|III| \leq \|\dot{S}_{Q_0} u\|_3^2 \|\nabla \bar{\Delta}_{Q_0} u\|_3 \leq \left(\lambda_{Q_0}^{-\frac{2}{3}} \sum_{q \in Q_0} \lambda_q^{\frac{2}{3}} d_q^2 \right) \left(d_{Q_{-1}}^2 + d_{Q_0}^2 \right)^{\frac{1}{2}}$$

Using $\lambda_{Q_0-q}^{-\frac{4}{3}} \leq \lambda_{Q_0-q}^{-\frac{2}{3}}$ in the above estimates on $|\Pi_Q|$ together w/ $\bar{H}(q) := \lambda_{|q|}^{-\frac{2}{3}}$, we arrive at

$$|\Pi_{Q_0, Q_1}| \leq \left(\bar{H} * d^2 \right)^{3/2}(Q_0) + \left(\bar{H} * d^2 \right)^{3/2}(Q_1),$$

where we also defined $d^2 := (d_q^2)_{q \geq -1}$.

Estimate on helicity

The helicity is def. as

$$H := \int u \cdot \nabla \times u \, dx$$

and the LP helicity H_Q satisfies

$$\begin{aligned} h_Q &:= \frac{d}{dt} H_Q = \frac{d}{dt} \int S_Q u \cdot S_Q w \, dx = \varepsilon_{ijk} \int \partial_t (S_Q u_i S_Q \partial_j u_k) \, dx \\ &= \varepsilon_{ijk} \int \left[S_Q (\partial_t u_i) S_Q (\partial_j u_k) + S_Q u_i S_Q (\partial_j \partial_t u_k) \right] \, dx \\ &= 2 \varepsilon_{ijk} \int S_Q (\partial_t u_i) S_Q (\partial_j u_k) \, dx = -2 \varepsilon_{ijk} \int S_Q (u_i \partial_t u_j) S_Q (\partial_j u_k) \, dx \\ &= 2 \int \text{Tr} \left[S_Q (u \otimes u) \cdot \nabla S_Q w \right]. \end{aligned}$$

$i+k$ yields odd k due to

Using the above estimates on $S_Q(u \otimes u)$ from the energy estimate, we find

$$|h_Q| \leq \left(\|S_Q(u \otimes u)\|_{3/2} + \|u \otimes u\|_{3/2} \right) \| \nabla S_Q w \|_3 \leq \left(\lambda_Q^{-2} \sum_{q \leq Q} \lambda_q^2 \| \Delta_q u \|_3^2 + \sum_{q > Q} \| \Delta_q u \|_3^2 \right) \left(\sum_{q \leq Q} \| \nabla \Delta_q w \|_3^2 \right)^{1/2}$$

Using the ~~two~~ estimates above, we obtain

$$\begin{aligned} |h_Q| &\leq \left(\lambda_Q^{-2/3} \sum_{q \leq Q} \lambda_q^{2/3} b_q^2 + \lambda_Q^{4/3} \sum_{q > Q} \lambda_q^{-4/3} b_q^2 \right) \left(\lambda_Q^{2/3} \sum_{q \leq Q} \lambda_q^{2/3} b_q^2 \right)^{1/2} \\ &\leq (M_b * b^2(Q))^{3/2} \end{aligned}$$

where we defined $b_i^2 = (b_q)_{q \geq i-1}$ $M_b^{(q,k)} = \begin{cases} \lambda_q^{-2/3} & q \geq 0 \\ \lambda_q^{4/3} & q < 0 \end{cases}$

We have $M_b \in \ell^1(\mathbb{Z})$ and thus Lemma 2 yields

$$\limsup_{Q \rightarrow \infty} |h_Q| \leq \limsup_{Q \rightarrow \infty} b_Q^3 = 0$$

provided $(b_q)_{q \geq 1} \in C_0$, i.e., $u \in B_{3, \infty}^{\frac{2}{3}}$.

Estimate on eustrophy

Define

$$P_\xi^\perp := I_3 - |\xi|^{-2} \xi \otimes \xi, \quad \xi \in \mathbb{R}^3$$

and

$$\nabla^\perp u := \mathcal{F}^{-1}(P_\cdot^\perp \mathcal{F}u),$$

and also the eustrophy

$$\Omega_Q := \frac{1}{2} \frac{d}{dt} \|S_Q w\|_2^2 = \int \text{Tr} [S_Q(u \otimes u) \cdot \nabla \nabla^\perp S_Q w] dx.$$

Employing the above estimates, we immediately obtain

$$\begin{aligned} |\Omega_Q| &\leq \left(\|r_Q(u, w)\|_{3/2} + \|u - S_Q u\|_3 \right) \|\nabla \nabla^\perp S_Q w\|_3 \\ &\leq \left(\lambda_Q^{-2} \sum_{q \leq Q} \|\Delta_q w\|_3^2 + \sum_{q > Q} \|\Delta_q u\|_3^2 \right) \left(\sum_{q \leq Q} \lambda_q^{1/2} \|\Delta_q w\|_3^2 \right)^{1/2} \\ &\leq \left(\lambda_Q^{-2} \|S_Q w\|_3^2 + \sum_{q > Q} \lambda_q^{-2} C_q^2 \right) \left(\sum_{q \leq Q} \lambda_q^4 C_q^2 \right)^{1/2}, \end{aligned}$$

where we defined $C_q := \|\Delta_q w\|_3$. If we further define

$$W(q) := \begin{cases} \lambda_q^4 & , q \geq 0 \\ \lambda_q^2 & , q < 0 \end{cases},$$

we arrive at

$$|\Omega_Q| \leq \|S_Q w\|_3^2 (W * C^2)^{1/2}(Q) + (W * C^2)^{3/2}(Q).$$

analogously def. to above

Lower bound on used estimate

l^1 sufficient

Lemma: Let $(a_q)_{q \geq 1} \in l^1$, $f: \mathbb{N} \rightarrow \mathbb{R}^+$ be some ^{strict} exponential $\begin{cases} \text{decay} \\ \text{growth} \end{cases}$. Then we have for any $N \in \mathbb{Z}^{\geq -1}$ large enough,
 $p \in [1, \infty]$, $\alpha \in [0, 1]$

$$(i) \inf_{q > N} (f(q)^{1-\alpha} a_q) f(N)^\alpha \lesssim \sum_{q > N} f(q) a_q \lesssim f(N) \|(a_q)_{q > N}\|_p$$

$$(ii) \max_{q \leq N} a_q g(q) \lesssim \sum_{q \leq N} g(q) a_q \lesssim g(N) \|(a_q)_{q \geq 1}\|_p$$

Proof: First of all, we prove the upper bounds. If $p = \infty$, both estimates are almost obvious.
 For $p < \infty$ we obtain

$$\sum_{q > N} f(q) a_q \stackrel{\text{Hölder}}{\lesssim} \|(a_q)_{q > N}\|_p \left(\sum_{q > N} f(q^{p'}) \right)^{\frac{1}{p'}} \stackrel{\text{geom.}}{\lesssim} \|(a_q)_{q > N}\|_p \left(\frac{f((N+1)^{p'})}{1 - f(p')} \right)^{\frac{1}{p'}}$$

$$\lesssim \|(a_q)_{q > N}\|_p f(N)$$

as well as

$$\sum_{q \leq N} g(q) a_q \lesssim \|(a_q)_{q \geq 1}\|_p \left(\sum_{q \leq N} g(q^{p'}) \right)^{\frac{1}{p'}} = \|(a_q)_{q \geq 1}\|_p \left(\frac{g((N+1)^{p'}) - 1}{g(p') - 1} + g(p') \right)^{\frac{1}{p'}}$$

$$\lesssim \|(a_q)_{q \geq 1}\|_p g(N).$$

As the lower bound of (ii) is also obvious, we will only show the lower bound of (i):

$$\sum_{q > N} f(q) a_q \geq \inf_{q > N} (f(q)^{1-\alpha} a_q) \sum_{q > N} f(q)^\alpha = \inf_{q > N} (f(q)^{1-\alpha} a_q) \frac{f((N+1)^\alpha)}{1 - f(\alpha)}$$

$$\gtrsim \inf_{q > N} (f(q)^{1-\alpha} a_q) f(N)^\alpha$$



