

Coupling to fermions

$$\begin{array}{cc} \text{Spin}(\mathbb{R}^4) & \text{Spin}(\mathbb{R}^4) \\ \downarrow & \downarrow \end{array}$$

Recall the spin representation of $\text{Cliff}(\mathbb{R}^4)$, $S \simeq S^+ \oplus S^-$.

$\text{Spin}(\mathbb{R}^4) \simeq \text{SU}(2)_+ \times \text{SU}(2)_-$ and in fact S^+ is fund^l rep. of $\text{SU}(2)_+$
 S^- is " " " $\text{SU}(2)_-$

Now if X is a spin 4-manifold (w/ principal $\text{Spin}(4)$ bundle B) then we get assoc. spin bundles

S_B^\pm . Dirac operator $\not{D}: T(S_B^\pm) \rightarrow T(S_B^\mp)$.

Now we can define a new theory by $\mathcal{L} = \begin{cases} \psi^+ \in T(\pi S_B^+) \\ \psi^- \in T(\pi S_B^-) \end{cases}$

$$S = \frac{1}{2} \int_X \langle \psi^+, \not{D} \psi^- \rangle \quad \langle, \rangle = \text{Hermitian inner product on } S^+$$

That's free (quadratic).

Can also couple to gauge fields: $\mathcal{L} = \begin{cases} (P, \nabla) \\ \psi^+ \in \pi T(S_B^+ \otimes V_P) \\ \psi^- \in \pi T(S_B^- \otimes V_P) \end{cases}$

$$S = \frac{1}{2} \int_X \langle \psi^+, \not{D} \psi^- \rangle \quad \not{D} = \text{covariant Dirac op:} \\ \sum_i p(e_i) D_{e_i} \quad (e_i \text{ orthonormal basis for } TX)$$

[Our rules for path integration over these fermions will be by analytic continuation from signature (3,1).]

It's a general principle of QFT that fermions usually "have $\frac{1}{2}$ -integer spin"
i.e. they are associated with reps of $\text{so}(\mathbb{R}^n)$ that don't integrate to $\text{SO}(\mathbb{R}^n)$
but only to $\text{Spin}(\mathbb{R}^n)$. "Spin-statistics theorem" (but has loopholes...)

Roughly a consequence of unitarity.