

Homework 3

- 1) Consider the equations of motion for a single rigid body subject to a constant gravitational force \mathbf{f}^g and a concentrated elastic force \mathbf{f}^c at \mathbf{x} with energy U . In terms of the usual components we have

$$\dot{\mathbf{q}} = \mathbf{v}, \quad \dot{\mathbf{Q}} = \mathbf{Q}[\boldsymbol{\omega} \times], \quad M\dot{\mathbf{v}} = \mathbf{f}^c + \mathbf{f}^g, \quad \Gamma\dot{\boldsymbol{\omega}} = (\Gamma\boldsymbol{\omega}) \times \boldsymbol{\omega} + \boldsymbol{\tau}^c, \quad (1)$$

where $\mathbf{f}^c = -\frac{\partial U}{\partial \mathbf{x}}$, $\boldsymbol{\tau}^c = \mathbf{Q}^T[(\mathbf{x} - \mathbf{q}) \times \mathbf{f}^c]$ and $\mathbf{x} = \mathbf{q} + \mathbf{Q}\boldsymbol{\xi}$ for some fixed body point $\boldsymbol{\xi}$.

- (a) Define the total kinetic energy of the body by $\Psi = \frac{1}{2}\mathbf{v} \cdot M\mathbf{v} + \frac{1}{2}\boldsymbol{\omega} \cdot \Gamma\boldsymbol{\omega} \geq 0$. Show that the rate of change of Ψ along any solution of (1) is given by

$$\dot{\Psi} = (\mathbf{f}^c + \mathbf{f}^g) \cdot \mathbf{v} + \boldsymbol{\tau}^c \cdot \boldsymbol{\omega}.$$

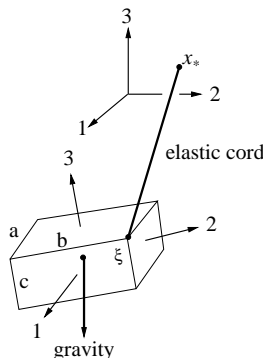
- (b) Define the total potential energy of the body by $\Phi = U(\mathbf{x}) - \mathbf{f}^g \cdot \mathbf{q}$, where $\mathbf{x} = \mathbf{q} + \mathbf{Q}\boldsymbol{\xi}$. Show that the rate of change of Φ along any solution of (1) is given by

$$\dot{\Phi} = -(\mathbf{f}^c + \mathbf{f}^g) \cdot \mathbf{v} - \boldsymbol{\tau}^c \cdot \boldsymbol{\omega}.$$

- (c) Use (a) and (b) to deduce that the total energy $E = \Psi + \Phi$ must remain constant along any solution of (1).

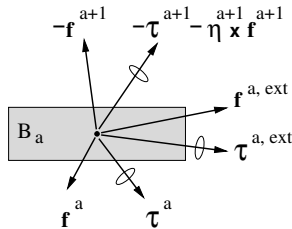
- (d) Consider the case $\mathbf{f}^g = 0$ and let $\mathbf{Q} = \varphi(\boldsymbol{\theta})$ be a coordinate map for SO_3 . Let $(\mathbf{q}_*, \boldsymbol{\theta}_*)$ be a local minimum of the elastic energy U so that $(\mathbf{q}_*, \boldsymbol{\theta}_*, 0, 0)$ is a steady state for the body. For simplicity assume there is a constant $\alpha > 0$ such that $U(\mathbf{q}, \boldsymbol{\theta}) - U(\mathbf{q}_*, \boldsymbol{\theta}_*) \geq \alpha(|\mathbf{q} - \mathbf{q}_*|^2 + |\boldsymbol{\theta} - \boldsymbol{\theta}_*|^2)$ and $\Psi(\mathbf{v}, \boldsymbol{\omega}) \geq \alpha(|\mathbf{v}|^2 + |\boldsymbol{\omega}|^2)$ for all $(\mathbf{q}, \boldsymbol{\theta}, \mathbf{v}, \boldsymbol{\omega})$. (Thus U and Ψ are “convex” with their global minimum at $(\mathbf{q}_*, \boldsymbol{\theta}_*)$ and $(0, 0)$ respectively). Use the result in (c) to show that the steady state $(\mathbf{q}_*, \boldsymbol{\theta}_*, 0, 0)$ is neutrally stable.

- 2) Consider a uniform rigid block of mass M and dimensions a , b and c subject to a gravitational force with components $\mathbf{f}^g = (0, 0, -Mg)$ and a concentrated Hookean spring force with components $\mathbf{f}^c = K(|\boldsymbol{\eta}| - \ell)\boldsymbol{\eta}/|\boldsymbol{\eta}|$ where $\boldsymbol{\eta} = \mathbf{x}_* - \mathbf{x}$ and $\mathbf{x} = \mathbf{q} + \mathbf{Q}\boldsymbol{\xi}$. Here K is the spring stiffness, ℓ is the relaxed length, \mathbf{x}_* are the coordinates of the the spring anchor point (fixed frame) and $\boldsymbol{\xi}$ are the coordinates of the spring attachment point (body frame). Recall that $\Gamma = \text{diag}(\Gamma_1, \Gamma_2, \Gamma_3)$ is a diagonal matrix of rotational inertia coefficients where $\Gamma_1 = \frac{M}{12}(b^2 + c^2)$, $\Gamma_2 = \frac{M}{12}(a^2 + c^2)$, $\Gamma_3 = \frac{M}{12}(a^2 + b^2)$.



- (a) In kg-m-s units consider the parameters $M = 1$, $a = 0.16$, $b = 0.24$, $c = 0.10$, $g = 10$, $K = 20$, $\ell = 0.25$, $\boldsymbol{\xi} = (\frac{a}{2}, \frac{b}{2}, \frac{c}{2})$ and $\mathbf{x}_* = (\frac{a}{2}, \frac{b}{2}, 3c)$. Use the explicit midpoint method to solve (1) on the interval $[0, 10]$ with various different time steps, say $\Delta t = 0.01, 0.02, 0.025$. Use the initial conditions $(\mathbf{q}_0, \mathbf{Q}_0, \mathbf{v}_0, \mathbf{w}_0) = (0, I, 0, 0)$. For each different time step make and superimpose plots of the kinetic energy Ψ , potential energy Φ and the total energy E versus t . (See the Matlab file *myprogram3A.m* at the course webpage for help getting started. To use this file you will also need to download the updated version of *plotBox.m*).
- (b) Are the numerical results consistent with your analysis in Problem 1abc? In particular, do the slopes of the kinetic and potential energy graphs differ by a minus sign? Does the total energy remain constant? Based on your results, do you think that $\Delta t = 0.025$ is an appropriate time step for simulating this system? What happens for time steps bigger than $\Delta t = 0.025$?

- 3) Consider a chain of rigid bodies where the net elastic and external loads on a typical body B_a are as follows:



$$\begin{aligned} \mathbf{f}^a &= \sum_i f_i^a \mathbf{d}_i^{a-1} \\ \boldsymbol{\tau}^a &= \sum_i \tau_i^a \mathbf{d}_i^a \\ \boldsymbol{\eta}^a &= \sum_i \eta_i^a \mathbf{d}_i^{a-1} \\ \mathbf{f}^{a, \text{ext}} &= \sum_i f_i^{a, \text{ext}} \mathbf{e}_i \\ \boldsymbol{\tau}^{a, \text{ext}} &= \sum_i \tau_i^{a, \text{ext}} \mathbf{d}_i^a. \end{aligned}$$

Show that the equations of motion for the usual configuration and velocity variables $(q^a, Q^a, v^a, \omega^a)$ for B_a take the form

$$\begin{aligned} \dot{q}^a &= v^a \\ \dot{Q}^a &= Q^a [\omega^a \times] \\ M^a \dot{v}^a &= Q^{a-1} f^a - Q^a f^{a+1} + f^{a, \text{ext}} \\ \Gamma^a \dot{\omega}^a &= (\Gamma^a \omega^a) \times \omega^a + \tau^a - \Lambda^{a+1} \tau^{a+1} - \eta^{a+1} \times f^{a+1} + \tau^{a, \text{ext}}, \end{aligned}$$

where Λ^{a+1} is the relative rotation matrix and η^{a+1} is the relative displacement between bodies B_a and B_{a+1} , that is, $\mathbf{d}_j^{a+1} = \sum_i \Lambda_{ij}^{a+1} \mathbf{d}_i^a$ and $\mathbf{q}^{a+1} = \mathbf{q}^a + \sum_i \eta_i^{a+1} \mathbf{d}_i^a$.

- 4) Consider a rigid body chain with bodies B_a ($a = 1, \dots, n$) and quadratic elastic interaction energies

$$U^a(\eta^a, u^a) = \frac{1}{2} \begin{Bmatrix} \eta^a - \hat{\eta}^a \\ u^a - \hat{u}^a \end{Bmatrix}^T \begin{pmatrix} G^a & C^a \\ [C^a]^T & K^a \end{pmatrix} \begin{Bmatrix} \eta^a - \hat{\eta}^a \\ u^a - \hat{u}^a \end{Bmatrix},$$

where $G^a \in \mathbf{R}^{3 \times 3}$ is the translational stiffness matrix, $K^a \in \mathbf{R}^{3 \times 3}$ is the rotational stiffness matrix, $C^a \in \mathbf{R}^{3 \times 3}$ is the coupling matrix and $\hat{\eta}^a, \hat{u}^a \in \mathbf{R}^3$ are the relaxed configuration parameters associated with the junction between bodies B_{a-1} and B_a (see figure below). For simplicity we suppose the chain is uniform, so that the energy parameters are independent of the index a .

- (a) Consider a chain composed of $n = 20$ uniform blocks, each of mass $M = 0.1$ and dimensions $0.16 \times 0.24 \times 0.10$ along the 1-2-3 body axes. Illustrate the relaxed shape of the chain for each of the following cases:

$$(\hat{\eta}, \hat{u}) = (0, 0, .5, 0, 0, 0), \quad (0, 0, .5, 0, 0, .17), \quad (.01, 0, .5, .01, .17).$$

In each case describe whether the chain appears curved/straight and twisted/untwisted.

- (b) Consider a chain as in (a) with $n = 10$, relaxed shape $(\hat{\eta}, \hat{u}) = (0, 0, .5, 0, 0, 0)$, translational stiffness $G = \text{diag}(20, 20, 20)$ and rotational stiffness $K = \text{diag}(3, 3, 3)$. Suppose the last body B_n is subject to a net external force with components $f^{n, \text{ext}} = (0, 0, t)$ for $t \in [0, 7]$. Use the explicit midpoint method to solve the chain equations of motion for each of the following values of the coupling matrix:

$$C = \text{diag}(0, 0, 0), \quad \text{diag}(0, 0, 1), \quad \text{diag}(0, 0, -1).$$

In each case use a time step of $\Delta t = 0.01$ and take the relaxed configuration with zero velocities as the initial condition. Illustrate the initial and final shape of the chain in each case. Briefly describe how the coupling matrix C affects the results. (See the Matlab file *myprogram3B.m* at the course webpage for help getting started. This file will also be useful for part (a).)

