

## $n$ -particle quantum statistics on graphs

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TexAMP – 11/14

# Quantum statistics

Single particle space configuration space  $X$ .

## Two particle statistics - alternative approaches:

- Quantize  $X^{\times 2}$  and restrict Hilbert space to the symmetric or anti-symmetric subspace.

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Bose-Einstein/Fermi-Dirac statistics.

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Bose-Einstein/Fermi-Dirac statistics.

- (Leinaas and Myrheim '77)  
Treat particles as indistinguishable,  $\psi(x_1, x_2) \equiv \psi(x_2, x_1)$ .  
Quantize two particle configuration space.

## Definition

Configuration space of  $n$  indistinguishable particles in  $X$ ,

$$C_n(X) = (X^{\times n} - \Delta_n) / S_n$$

where  $\Delta_n = \{x_1, \dots, x_n \mid x_i = x_j \text{ for some } i \neq j\}$ .

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1st homology groups of  $C_n(\mathbb{R}^d)$ :

- $H_1(C_n(\mathbb{R}^d)) = \mathbb{Z}_2$  for  $d \geq 3$ .  
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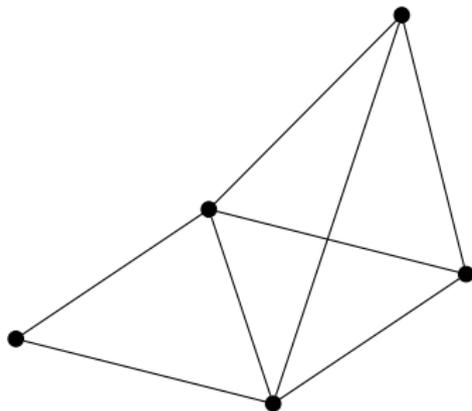
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- $H_1(C_n(\mathbb{R})) = 1$   
 particles cannot be exchanged.

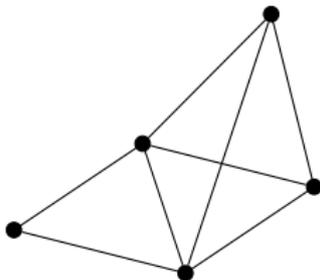
What happens on a graph where the underlying space has arbitrarily complex topology?



## Graph connectivity

- Given a connected graph  $\Gamma$  a *k-cut* is a set of  $k$  vertices whose removal makes  $\Gamma$  disconnected.
- $\Gamma$  is *k-connected* if the minimal cut is size  $k$ .
- **Theorem** (Menger) For a  $k$ -connected graph there exist at least  $k$  independent paths between every pair of vertices.

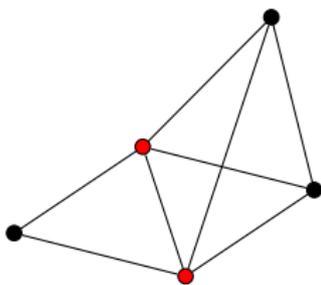
Example:



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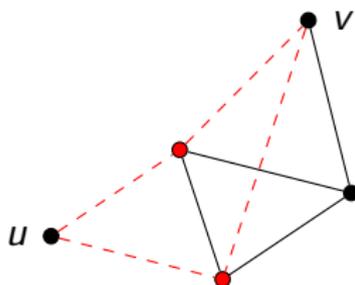


Two cut

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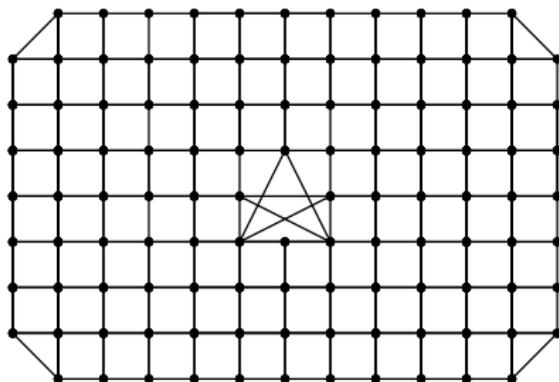
Two independent paths joining  $u$  and  $v$ .

## Features of graph statistics

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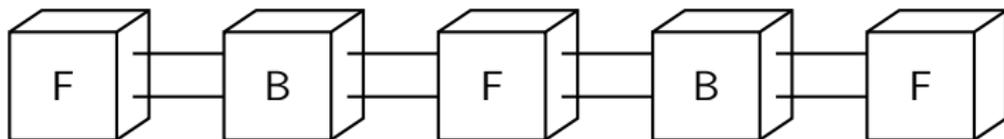
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A two dimensional lattice with a small section that is non-planar is locally planar but has Bose/Fermi statistics.

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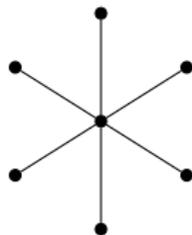


For example, one could construct a chain of 3-connected non-planar components where particles behave with alternating Bose/Fermi statistics.

On **1-connected graphs** the statistics *depends* on the no. of particles  $n$ .

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Example, star with  $E$  edges.



no. of anyon phases

$$\binom{n + E - 2}{E - 1} (E - 2) - \binom{n + E - 2}{E - 2} + 1 .$$

# 1st homology group of graph

By the structure theorem for finitely generated modules  
 (for a suitably subdivided graph  $\Gamma$ )

$$H_1(C_n(\Gamma)) = \mathbb{Z}^k \oplus \mathbb{Z}_{n_1} \oplus \dots \oplus \mathbb{Z}_{n_l}, \quad (2)$$

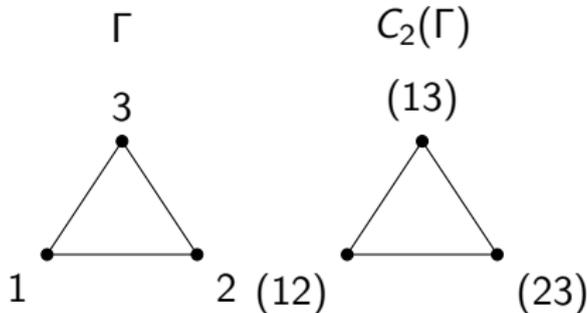
where  $n_i | n_{i+1}$ .

So  $H_1(C_n(\Gamma))$  is determined by  $k$  free (anyon) phases  $\{\phi_1, \dots, \phi_k\}$   
 and  $l$  discrete phases  $\{\psi_1, \dots, \psi_l\}$  such that for each  $i \in \{1, \dots, l\}$

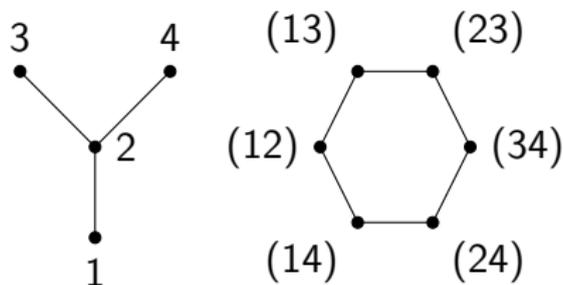
$$n_i \psi_i = 0 \pmod{2\pi}, \quad n_i \in \mathbb{N} \quad \text{and} \quad n_i | n_{i+1}. \quad (3)$$

# Basic cases

For 2 particles.



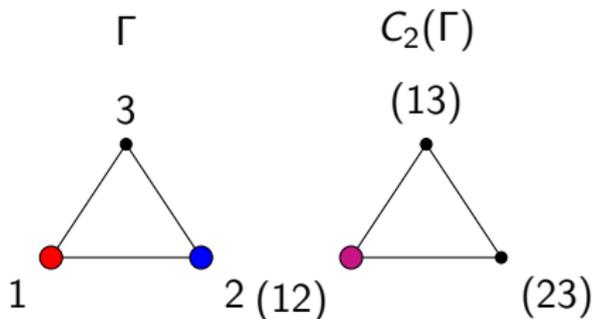
Exchange of 2 particles around loop  $c$ ; one free phase  $\phi_{c2}$ .



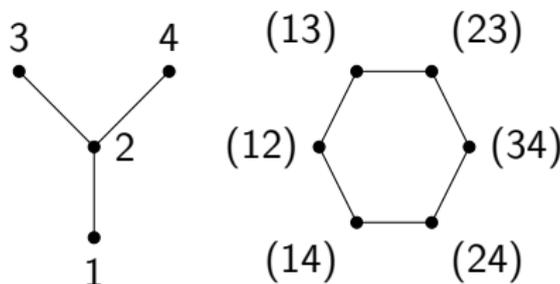
Exchange of 2 particles at Y-junction; one free phase  $\phi_Y$ .

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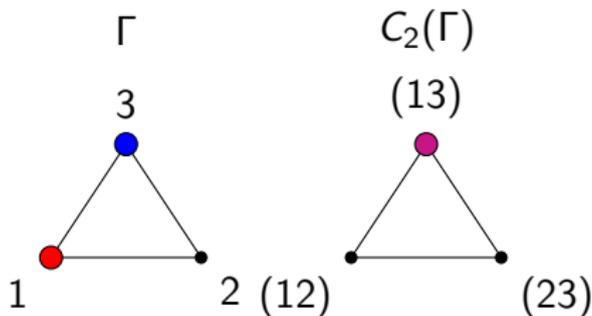
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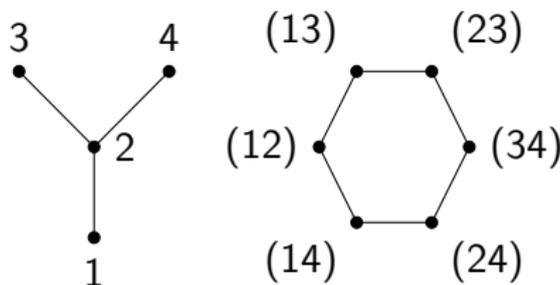
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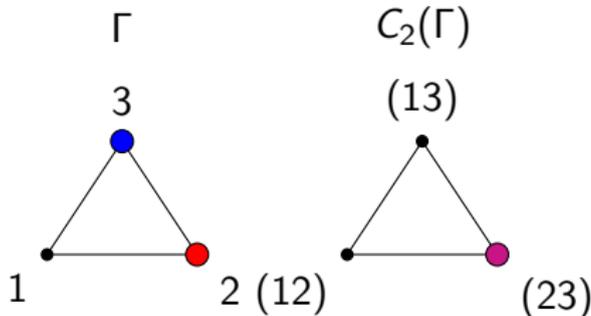
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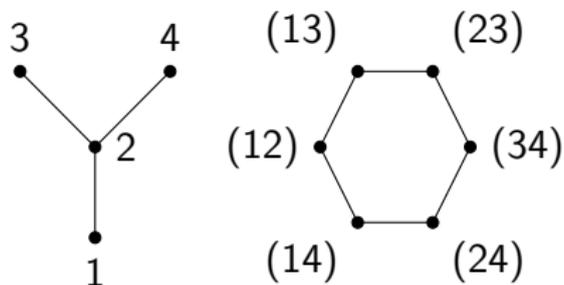
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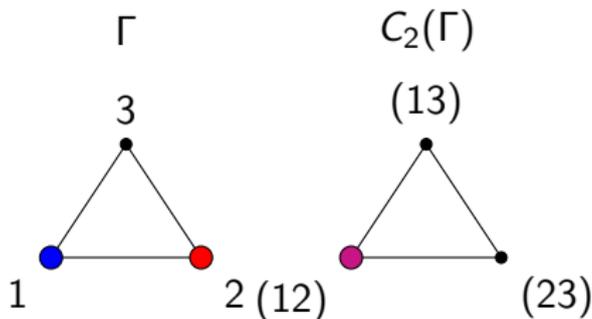
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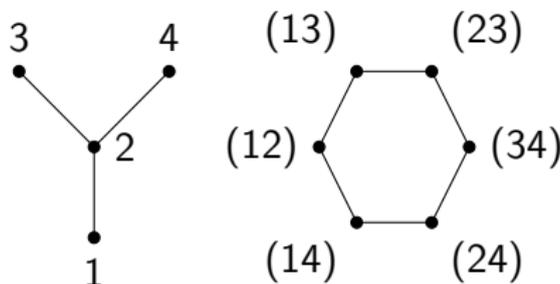
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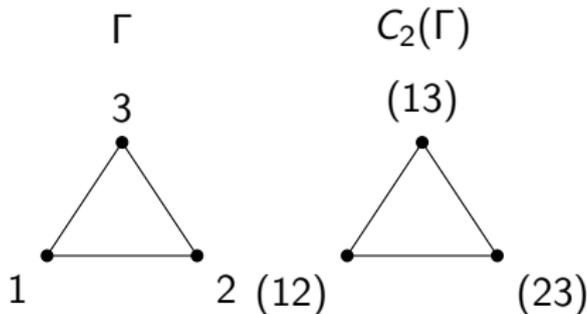
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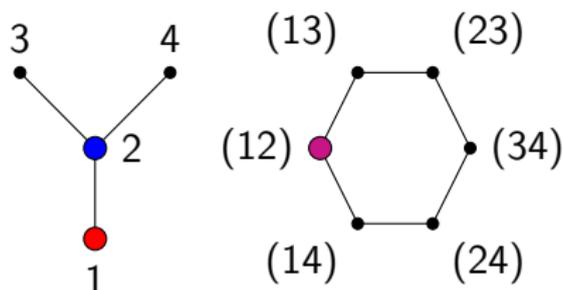
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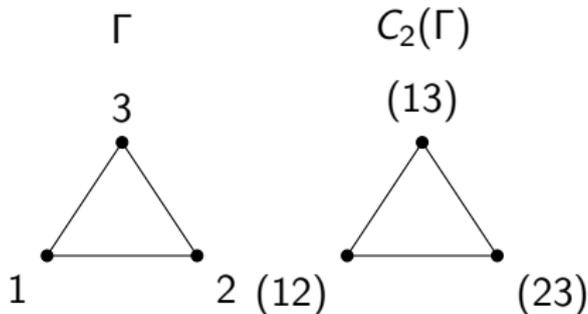
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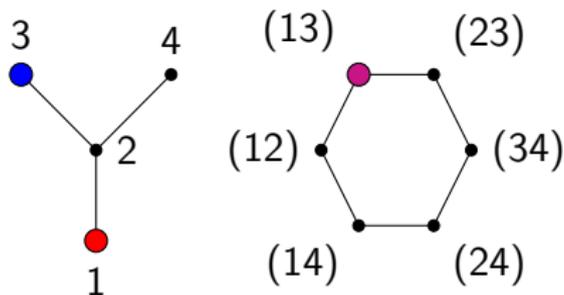
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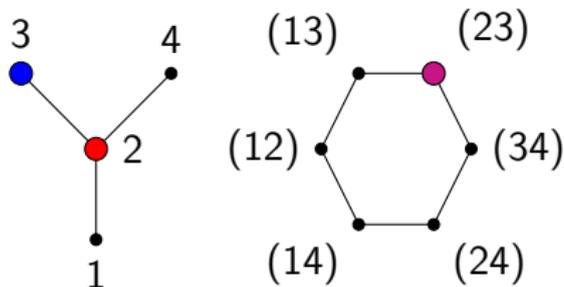
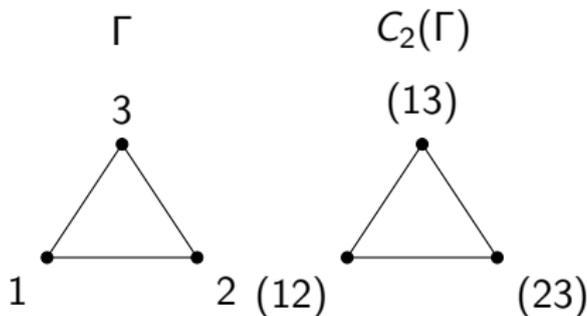
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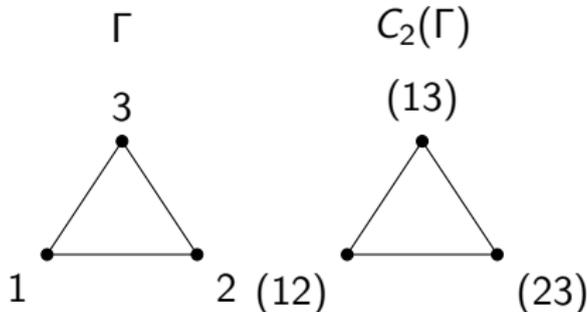
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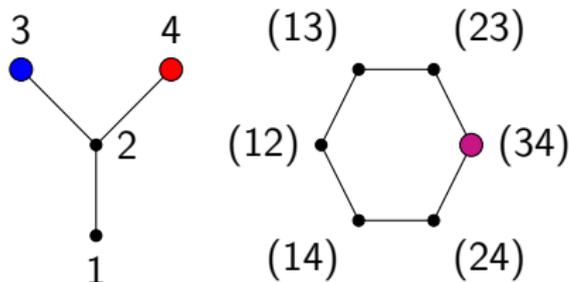


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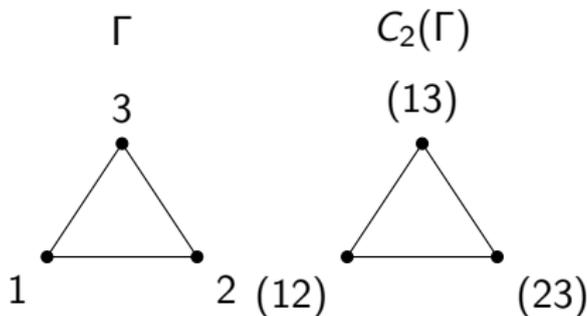
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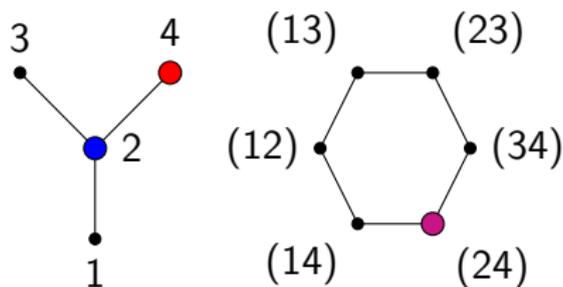
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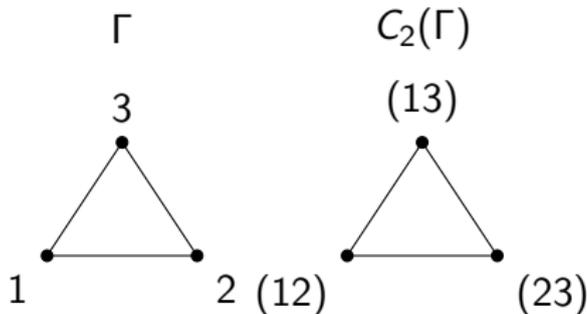
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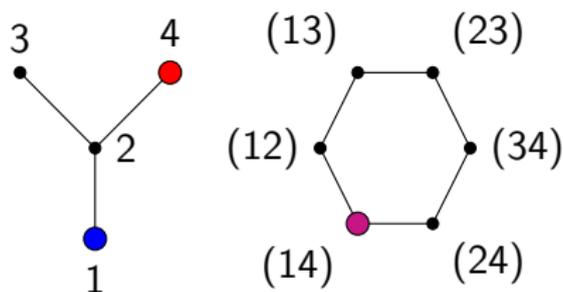
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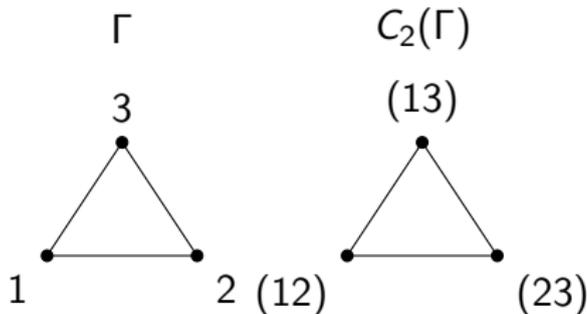
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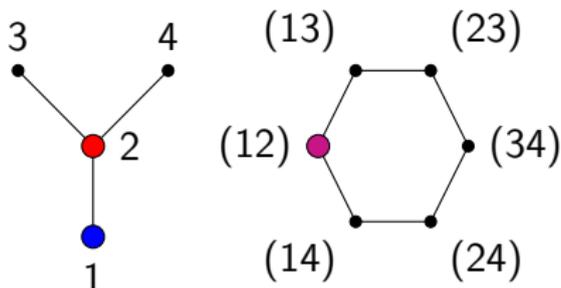
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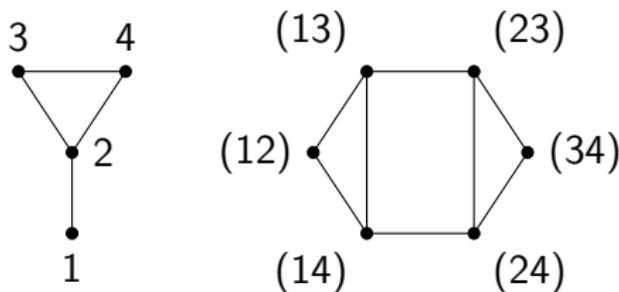


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# Lasso graph

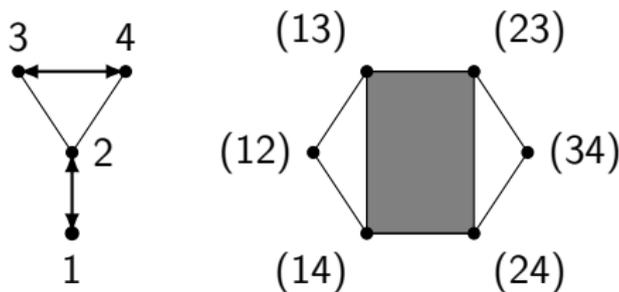


Identify three 2-particle cycles:

- (i) Rotate both particles around loop  $c$ ; phase  $\phi_{c,2}$ .
- (ii) Exchange particles on Y-subgraph; phase  $\phi_Y$ .
- (iii) Rotate one particle around loop  $c$  other particle at vertex 1;  
 $(1, 2) \rightarrow (1, 3) \rightarrow (1, 4) \rightarrow (1, 2)$ , phase  $\phi_{c,1}^1$ .

Relation from contactable 2-cell  $\phi_{c,2} = \phi_{c,1}^1 + \phi_Y$ .

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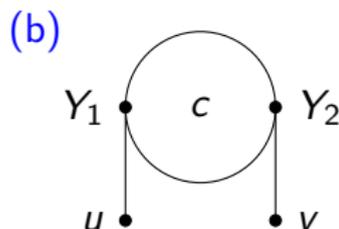
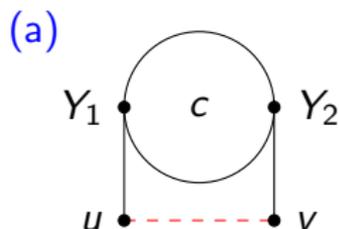
(a)  $u$  and  $v$  joined by path disjoint with  $c$ .

$\phi_{c,1}^u = \phi_{c,1}^v$  as exchange cycles homotopy equivalent.

(b)  $u$  and  $v$  *only* joined by paths through  $c$ .

Two lasso graphs so  $\phi_{c,2} = \phi_{c,1}^u + \phi_{Y_1}$  &  $\phi_{c,2} = \phi_{c,1}^v + \phi_{Y_2}$ .

Hence  $\phi_{c,1}^u - \phi_{c,1}^v = \phi_{Y_2} - \phi_{Y_1}$ .



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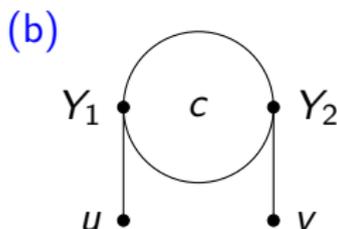
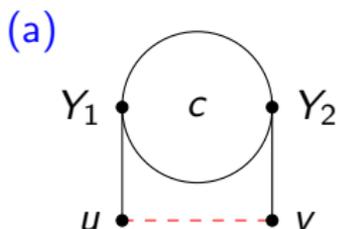
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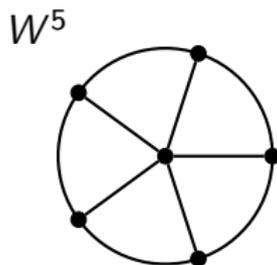
- Relations between phases involving  $c$  encoded in phases  $\phi_Y$ .

$H_1(C_2(\Gamma)) = \mathbb{Z}^{\beta_1(\Gamma)} \oplus A$ , where  $A$  determined by  $Y$ -cycles.

- In (a) we have a  $\mathcal{B}$  subgraph & using (b) also  $\phi_{Y_1} = \phi_{Y_2}$ .

## 3-connected graphs

The prototypical 3-connected graph is a *wheel*  $W^k$ .



### Theorem (Wheel theorem)

Let  $\Gamma$  be a simple 3-connected graph different from a wheel. Then for some edge  $e \in \Gamma$  either  $\Gamma \setminus e$  or  $\Gamma / e$  is simple and 3-connected.

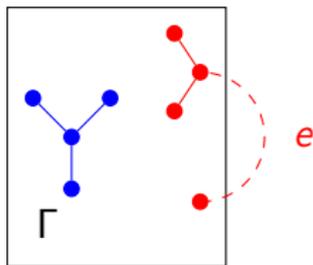
- $\Gamma \setminus e$  is  $\Gamma$  with the edge  $e$  removed.
- $\Gamma / e$  is  $\Gamma$  with  $e$  contracted to identify its vertices.

## Lemma

*For 3-connected simple graphs all phases  $\phi_\gamma$  are equal up to a sign.*

**Sketch proof.** The lemma holds on  $K_4$  (minimal wheel). By wheel thm we only need to show that adding an edge or expanding a vertex any new phases  $\phi_\gamma$  are the same as the original phase.

*Adding an edge:*  $\Gamma \cup e$

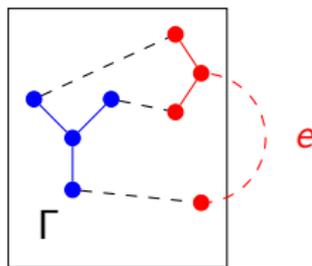


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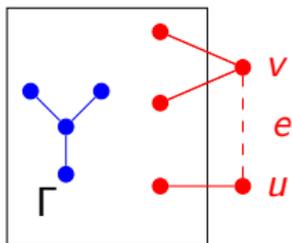
Using 3-connectedness identify independent paths in  $\Gamma$  to make  $\mathcal{B}$ .  
 Then  $\phi_{\Gamma \cup e} = \phi_\Gamma$ .

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*Vertex expansion:* Split vertex of degree  $> 3$  into two vertices  $u$  and  $v$  joined by a new edge  $e$ .

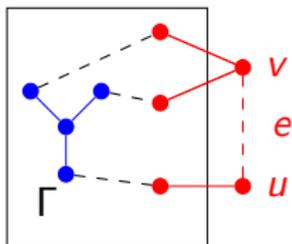


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Using 3-connectedness identify independent paths in  $\Gamma$  to make  $\beta$ .

Then  $\phi_Y = \phi_{\beta}$ .

## Theorem

*For a 3-connected simple graph,  $H_1(C_2(\Gamma)) = \mathbb{Z}^{\beta_1(\Gamma)} \oplus A$ , where  $A = \mathbb{Z}_2$  for non-planar graphs and  $A = \mathbb{Z}$  for planar graphs.*

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## Proof.

- For  $K_5$  and  $K_{3,3}$  every phase  $\phi_Y = 0$  or  $\pi$ . By Kuratowski's theorem a non-planar graph contains a subgraph which is isomorphic to  $K_5$  or  $K_{3,3}$ .

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## Proof.

- For  $K_5$  and  $K_{3,3}$  every phase  $\phi_Y = 0$  or  $\pi$ . By Kuratowski's theorem a non-planar graph contains a subgraph which is isomorphic to  $K_5$  or  $K_{3,3}$ .
- For planar graphs the anyon phase can be introduced by drawing the graph in the plane and integrating the anyon vector potential  $\frac{\alpha}{2\pi} \hat{z} \times \frac{r_1 - r_2}{|r_1 - r_2|^2}$  along the edges of the two-particle graph, where  $r_1$  and  $r_2$  are the positions of the particles.

## Summary

- Full classification of abelian quantum statistics on graphs by decomposing graph in 1-, 2- and 3-connected components.
- Physical insight into dependence of statistics on graph connectivity.
- Interesting new features of graph statistics.
- Statistics incorporated in gauge potential.



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