

Strong traces for solutions to multidimensional scalar conservation laws

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Abstract

In this paper we consider multidimensional scalar conservation laws without BV estimates defined in a subset $\Omega \subset \mathbb{R}^+ \times \mathbb{R}^d$. We show that, with a non-degeneracy hypothesis on the flux, we can define a strong notion of trace at the boundary of Ω reached by L^1 convergence.

Key words. conservation law – trace theorem – kinetic formulation – boundary value problem – averaging lemma

1. Introduction

In this article we consider an open subset $\Omega \subset \mathbb{R}^+ \times \mathbb{R}^d$ and functions $u \in L^\infty(\Omega)$ solutions to the scalar conservation law:

$$\partial_t u + \operatorname{div}_x A(u) = 0, \quad (t, x) \in \Omega. \quad (1)$$

The flux function A is assumed to be regular from \mathbb{R} to \mathbb{R}^d . As is usual, we deal only with entropy solutions, namely which fulfill in the sense of distribution

$$\partial_t \phi(u) + \operatorname{div}_x H(u) \leq 0, \quad (t, x) \in \Omega, \quad (2)$$

for every convex function ϕ and related entropy flux defined by $H' = A' \phi'$. The problem is to define in a strong way the value of u at the boundary of Ω . Following [3] we consider domains Ω which have a “regular deformable Lipschitz boundary” namely which verify:

(i): For each $\hat{z} = (\hat{t}, \hat{x}) \in \partial\Omega$, there exists $r_{\hat{z}} > 0$, a Lipschitz mapping

$\gamma_{\hat{z}} : \mathbb{R}^d \rightarrow \mathbb{R}$ and an isometry for the euclidean norm $\mathcal{R}_{\hat{z}} : \mathbb{R}^{d+1} \rightarrow \mathbb{R}^{d+1}$ such that:

$$\begin{aligned} \mathcal{R}_{\hat{z}}(\hat{z}) &= 0, \\ \mathcal{R}_{\hat{z}}(\Omega) \cap]-r_{\hat{z}}, r_{\hat{z}}[^{d+1} &= \{y \in]-r_{\hat{z}}, r_{\hat{z}}[^{d+1} \mid y_0 > \gamma_{\hat{z}}(y_1, \dots, y_d)\}. \end{aligned}$$

(ii): There exists at least one $\partial\Omega$ -regular Lipschitz deformation,

where, for K open subset of $\partial\Omega$, we call K -regular Lipschitz deformation every function $\psi : [0, 1] \times K \rightarrow \bar{\Omega}$ homeomorphism bi-Lipschitz over its image which verifies:

(iii): $\psi(0, \cdot) \equiv I_K$, where I_K is the identity map over K .

(iv):

$$\lim_{s \rightarrow 0^+} \nabla \psi(s, \cdot) \circ \tilde{\gamma}_{\hat{z}} = \nabla \tilde{\gamma}_{\hat{z}} \text{ in } L^1(]-r_{\hat{z}}, r_{\hat{z}}[^d \cap \mathcal{R}_{\hat{z}}(K)),$$

where for every $\hat{z} \in \partial\Omega$ $\tilde{\gamma}_{\hat{z}}$ is the restriction to $]-r_{\hat{z}}, r_{\hat{z}}[^d \cap \mathcal{R}_{\hat{z}}(K)$ of the map $(y_1, \dots, y_d) \rightarrow \mathcal{R}_{\hat{z}}^{-1}(\gamma_{\hat{z}}(y_1, \dots, y_d), y_1, \dots, y_d)$. Notice that the set of Ω with regular Lipschitz boundary includes star-shaped domains, the domains whose boundaries satisfy the cone property, and domains with \mathcal{C}^2 boundaries ([3]).

More precisely, we want to study the limit of $u(\psi(s, \cdot))$ when s goes to zero for such $\partial\Omega$ -regular Lipschitz deformation.

It is possible to define weak traces using Young measures (as in [17] for instance). Notice that since the limit is weak, even if we can define the trace of u itself (denoted u^τ), we did not know if the trace of $G(u)$ is equal to $G(u^\tau)$ for a nonlinear function G . The aim of this paper is to show that, with an hypothesis of non-degeneracy of A , in fact the trace is reached by L^1_{loc} convergence.

Theorem 1. *Let $\Omega \subset \mathbb{R}^{d+1}$ have a regular deformable Lipschitz boundary, and assume that $A \in [C^3(\mathbb{R})]^d$ and verifies for every $(\tau, \zeta) \in \mathbb{R} \times \mathbb{R}^d$, $(\tau, \zeta) \neq (0, 0)$:*

$$\mathcal{L}(\{\xi \mid \tau + \zeta \cdot A'(\xi) = 0\}) = 0, \quad (3)$$

where \mathcal{L} is the Lebesgue measure. Then for every function $u \in L^\infty(\Omega)$ which verifies (1)(2) in Ω , there exists $u^\tau \in L^\infty(\partial\Omega)$ such that, for every $\partial\Omega$ -regular Lipschitz deformation ψ and every compact set $K \subset\subset \partial\Omega$:

$$\text{esslim}_{s \rightarrow 0} \int_K |u(\psi(s, \hat{z})) - u^\tau(\hat{z})| d\mathcal{H}^d(\hat{z}) = 0, \quad (4)$$

where \mathcal{H}^d is the d -dimensional Hausdorff measure. Especially, for every smooth function G , we have

$$[G(u)]^\tau = G(u^\tau). \quad (5)$$

This can be seen as a regularization effect at the boundary induced by the non-degeneracy of A (see [5, 14] for other results of this kind). Our method is based on the kinetic formulation of scalar conservation laws introduced by Lions Perthame Tadmor [11] which allows to use the so-called averaging lemmas [1, 9, 7, 16] with the assumption (3). This can be seen as a non-degeneracy property since it avoids flux functions whose restriction to an open subset is linear. Hence we use a localization method first introduced in [18]. Let us mention two consequences to motivate this result.

Initial layers

Here, we fix the open set $\Omega =]0, +\infty[\times \mathbb{R}^d$. Let us consider an initial value $u_0 \in L^\infty(\mathbb{R}^d)$. Kruzkov has shown in [10] the uniqueness of the solution of (1)(2) verifying for any $R > 0$:

$$\frac{1}{T} \int_0^T \int_{|x| \leq R} |u(t, x) - u_0(x)| dx dt \longrightarrow 0, \quad (6)$$

when T tends to 0. Now, if we don't have this assumption but only the weak form:

$$\frac{1}{T} \int_0^T \int_{\mathbb{R}^d} (u(t, x) - u_0(x)) \phi(x) dx dt \longrightarrow 0, \text{ when } T \rightarrow 0$$

for every test function ϕ , the question is: Have we got still uniqueness? This question was first mentioned by Natalini and arises naturally in relaxation limit problems. Chen and Rascle have shown in [4] that the answer is yes in the monodimensional case with an assumption of non-degeneracy on the flux. With the hypothesis (3), Theorem 1 extends this result to the multi-dimensional case. Indeed, it shows that every uniformly bounded solution of (1)(2) in $]0, +\infty[\times \mathbb{R}^d$ verifies (6), and so, is the unique Kruzkov solution. Notice that this problem is still open for a general flux A .

Initial-boundary problem

Now we consider a regular open subset $\omega \in \mathbb{R}^d$, we denote ν the normal at the boundary and we take $\Omega =]0, +\infty[\times \omega$. We introduce an initial value $u_0 \in L^\infty(\mathbb{R}^d)$ and a boundary value $u^b \in L^\infty(]0, +\infty[\times \partial\omega)$. The following strong boundary condition has been introduced by Bardos, Leroux and Nedelec in [2]:

$$\text{sign}(u(t, \hat{x}) - u^b(t, \hat{x}))(A(u(t, \hat{x})) \cdot \nu(\hat{x}) - A(k) \cdot \nu(\hat{x})) \geq 0 \quad (7)$$

$$\text{for all } k \text{ in the interval with endpoints } u(t, \hat{x}) \text{ and } u^b(t, \hat{x}) \quad (8)$$

for almost every boundary point $(t, \hat{x}) \in]0, +\infty[\times \partial\omega$. They have shown that if u_0 and u^b are regular, then there exists a unique solution $u \in C(\mathbb{R}^+, L^1_{\text{loc}}(\omega))$ solution of (1)(2) in Ω , satisfying (7)(8) on $\partial\omega$ and $u(t = 0) = u_0$. Moreover this is the physical relevant one since it can be obtained by the viscosity method. Some equivalent formulations have been proposed

aftwards (see for example [8] [12]). But all those formulations need to restrict themselves to regular boundary values in order to have BV solutions which allows them to have strong traces reached in L^1 . This was needed in order to give a meaning at (7). A more general formulation has been introduced by Szepessy in [17] which can be defined without strong traces, using DiPerna tools from [6]. But he shows the existence and uniqueness of the solution only in the framework of regular boundary values. Finally, Otto has introduced a quite complicated weak formulation [13] which allows him to show the existence and uniqueness of solution for general boundary values.

Now, if we consider a flux A verifying (3), then thanks to Theorem 1, the condition (7) is well defined whatever the solution $u \in L^\infty(\Omega)$ of (1)(2) is. Then Otto's condition is equivalent to (7) and we can restrict ourselves to this strong formulation even for non smooth boundary values.

2. Reformulation of the problem

We have:

$$\partial\Omega \subset \bigcup_{z \in \partial\Omega} \mathcal{R}_z^{-1}(\cdot - r_z, r_z)^{[d+1]},$$

where r_z and \mathcal{R}_z are defined by property (i) of regular deformable Lipschitz boundary. Since $\partial\Omega \cap B(0, n)$ is a compact subset for every integer n , there exists a finite set I_n such that:

$$(\partial\Omega \cap B(0, n)) \subset \bigcup_{\alpha \in I_n} \mathcal{R}_{\hat{z}_\alpha}^{-1}(\cdot - r_{\hat{z}_\alpha}, r_{\hat{z}_\alpha})^{[d+1]}.$$

So, $I = \cup I_n$ is a countable set such that:

$$\partial\Omega = \bigcup_{\alpha \in I} \Gamma_\alpha,$$

where Γ_α is defined (using the property (i) of regular deformable Lipschitz boundary) by:

$$\Gamma_\alpha = \mathcal{R}_\alpha^{-1}(\{y \in \cdot - r_\alpha, r_\alpha\}^{[d+1]} \mid y_0 = \gamma_\alpha(y_1, \dots, y_d)\}).$$

In order to simplify the notation we write α instead of \hat{z}_α in the indices, and we denote in the same way

$$\Omega_\alpha = \{y \in \cdot - r_\alpha, r_\alpha\}^{[d+1]} \mid y_0 > \gamma_\alpha(y_1, \dots, y_d)\}.$$

From now on we work in Ω_α and in the new y coordinates. We denote $A_\alpha(\xi) = \mathcal{R}_\alpha(\xi, A^1(\xi), \dots, A^d(\xi))$ and $H_\alpha(\xi) = \mathcal{R}_\alpha(\phi(\xi), H^1(\xi), \dots, H^d(\xi))$. We define $u_\alpha : \Omega_\alpha \rightarrow \mathbb{R}$ by $u_\alpha(y) = u(\mathcal{R}_\alpha^{-1}(y))$. In the y coordinates, (1)(2) correspond in Ω_α to:

$$\operatorname{div}_y A_\alpha(u_\alpha) = 0 \tag{9}$$

$$\operatorname{div}_y H_\alpha(u_\alpha) \leq 0. \quad (10)$$

We use for the proof the kinetic formulation due to Lions, Perthame and Tadmor. In order to do so we set $L = \|u\|_{L^\infty(\Omega)}$ and introduce a new variable $\xi \in]-L, L[$ and denote for every $v \in]-L, L[$:

$$\begin{aligned} \chi(v, \xi) &= \mathbf{1}_{\{0 \leq \xi \leq v\}} \text{ if } v \geq 0 \\ &= -\mathbf{1}_{\{v \leq \xi \leq 0\}} \text{ if } v < 0. \end{aligned}$$

Then we introduce new functions called microscopic functions which depend on ξ and on a variable z which can lie on Ω_α , on Γ_α or on a local space as we will see later. We will consider especially the following ones:

Definition 1. Let N be an integer, \mathcal{O} be an open set of \mathbb{R}^N , and the microscopic function $f \in L^\infty(\mathcal{O} \times]-L, L[)$ be such that $0 \leq \operatorname{sgn}(\xi)f(z, \xi) \leq 1$ for almost every (z, ξ) . We say that f is a χ -function if there exists $u \in L^\infty(\mathcal{O})$ such that for almost every $z \in \mathcal{O}$:

$$f(z, \cdot) = \chi(u(z), \cdot).$$

Notice that in this case $u(z) = \int_{-L}^L f(z, \xi) d\xi$.

Lions Perthame and Tadmor have shown in [11] the following theorem:

Theorem 2. A function $u \in L^\infty(\Omega_\alpha)$ with $|u| \leq L$ is solution of (9)(10) in Ω_α if and only if there exists a nonnegative measure $m \in \mathcal{M}^+(\Omega_\alpha \times]-L, L[)$ such that the related χ -function f defined by $f(y, \cdot) = \chi(u(y), \cdot)$ for almost every $y \in \Omega_\alpha$ verifies:

$$a(\xi) \cdot \nabla_y f = \partial_\xi m \quad (11)$$

in $\Omega_\alpha \times]-L, L[$ with $a(\xi) = A_\alpha'(\xi)$.

Notice that the hypothesis (3) is equivalent to:

$$\mathcal{L}(\{\xi \mid \zeta \cdot a(\xi) = 0\}) = 0, \quad (12)$$

for every $\zeta \in \mathbb{R}^{d+1} \setminus \{0\}$.

The χ -functions have the following properties which are linked to Young measures.

Lemma 1. Let \mathcal{O} be an open set of \mathbb{R}^N and $f_n \in L^\infty(\mathcal{O} \times]-L, L[)$ be a sequence of χ -functions converging weakly to $f \in L^\infty(\mathcal{O} \times]-L, L[)$. We denote $u_n(\cdot) = \int_{-L}^L f_n(\cdot, \xi) d\xi$ and $u(\cdot) = \int_{-L}^L f(\cdot, \xi) d\xi$. Then the three following propositions are equivalent:

- f_n converges strongly to f in $L^1_{\text{loc}}(\mathcal{O} \times]-L, L[)$,
- u_n converges strongly to u in $L^1_{\text{loc}}(\mathcal{O})$,
- f is a χ -function.

Proof of Lemma 1: Notice that if we denote $\mu_n(\cdot, \xi)$ the Young measure related to $u_n(\cdot) = \int_{-L}^L f_n(\cdot, \xi) d\xi$ then, since f_n is a χ -function we have

$$\mu_n = \delta_0(\xi) - \partial_\xi f_n. \quad (13)$$

If f_n converges strongly, the same holds for u_n . If u_n converges strongly then passing to the limit in (13) we find that $\delta_0(\xi) - \partial_\xi f$ is a Dirac function with respect to ξ providing that f is a χ -function. Finally, if f is a χ -function, in particular $\text{sgn}(\xi)f = f^2$ so $\|f_n\|_{L^2_{\text{loc}}(\mathcal{O} \times]-L, L[)}$ converges to $\|f\|_{L^2_{\text{loc}}(\mathcal{O} \times]-L, L[)}$. This provides the strong convergence of f_n in $L^2_{\text{loc}}(\mathcal{O} \times]-L, L[)}$ and then in $L^1_{\text{loc}}(\mathcal{O} \times]-L, L[)}$. \square

In the following we will use the notation:

$$a(\xi) = (a^0(\xi), \hat{a}(\xi)) \text{ with } \hat{a}(\xi) = (a^1(\xi), \dots, a^d(\xi)) \quad (14)$$

$$y = (y_0, \hat{y}) \text{ with } \hat{y} = (y_1, \dots, y_d). \quad (15)$$

3. A criterion for the existence of strong traces

In this section we consider Ω_α for a fixed α and the χ -function f associated to u_α . Since Γ_α is parametrized by \hat{y} , for every Γ_α -regular Lipschitz deformation ψ and every $\hat{y} \in]-r_\alpha, r_\alpha[^d$ we set:

$$\begin{aligned} \tilde{\psi}(s, \hat{y}) &= \psi(s, \mathcal{R}_\alpha^{-1}(\gamma_\alpha(\hat{y}), \hat{y})), \\ f_\psi(s, \hat{y}, \xi) &= f(\tilde{\psi}(s, \hat{y}), \xi). \end{aligned}$$

Proposition 1. *Let f be a solution of (11) in $\Omega_\alpha \times]-L, L[$ with $a(\xi)$ verifying (12). Then there exists a unique $f^\tau \in L^\infty(]-r_\alpha, r_\alpha[^d \times]-L, L[)$ such that for all ψ Γ_α -regular Lipschitz deformation:*

$$\text{esslim}_{s \rightarrow 0} f_\psi(s, \cdot, \cdot) = f^\tau \text{ in } H^{-1}(]-r_\alpha, r_\alpha[^d \times]-L, L[).$$

This shows the existence of a weak trace on Γ_α which does not depend on the way chosen to reach the boundary. This result is closely related to [13] and a closed version was used in [17]. We give here a proof in our framework of kinetic formulation.

Proof of Proposition 1: Since $\|f_\psi(s, \cdot, \cdot)\|_{L^\infty} \leq 1$, by Weak compactness and Sobolev imbedding, for every regular Lipschitz deformation ψ and every sequence s^n which tends to 0 there exists a subsequence n_p and a function $g_\psi^\tau \in L^\infty(]-r_\alpha, r_\alpha[^d \times]-L, L[)$ such that

$$f_\psi(s^{n_p}, \cdot, \cdot) \xrightarrow{H^{-1} \cap L^\infty W^*} g_\psi^\tau \text{ when } n_p \rightarrow +\infty. \quad (16)$$

Let us now show that g_ψ^τ does not depend on ψ , on the sequence s^n and s^{n_p} . In order to do so, let us first consider the entropy flux associated to entropy ϕ :

$$H_\phi(y) = \int_{-L}^L a(\xi)\phi'(\xi)f(y, \xi) d\xi. \quad (17)$$

Multiplying (11) by $\phi'(\xi)$ and integrating it with respect to ξ we find:

$$\operatorname{div}_y H_\phi = - \int_{-L}^L \phi''(\xi)m(y, d\xi) \in \mathcal{M}(]-r_\alpha, r_\alpha[^{d+1}).$$

We can now use the following Theorem proved by Chen and Grid in [3]:

Theorem 3. *Let Ω be an open set with regular deformable Lipschitz boundary and $F \in [L^\infty(\Omega)]^{d+1}$ be such that $\operatorname{div}_y F$ is a bounded measure. Then there exists $F.\nu \in L^\infty(\partial\Omega)$ such that for every ψ $\partial\Omega$ -regular Lipschitz deformation:*

$$\operatorname{esslim}_{s \rightarrow 0} F(\psi(s, \cdot)).\nu_s(\cdot) = F.\nu \text{ in } L^\infty(\partial\Omega) \text{ w*},$$

where ν_s is a unit outward normal field of $\psi(\{s\} \times \partial\Omega)$.

So, since Γ_α is a regular deformable Lipschitz boundary, this Theorem insures that there exists $H_\phi^\tau.\nu \in L^\infty(]-r_\alpha, r_\alpha[^d)$ which does not depend on ψ such that

$$H_\phi(\tilde{\psi}(s, \cdot)).\nu_s(\cdot) \xrightarrow[s \rightarrow 0]{\mathcal{D}'(]-r_\alpha, r_\alpha[^d)} H_\phi^\tau.\nu, \quad (18)$$

for every regular Lipschitz deformation ψ . Thanks to property (iv) of regular Lipschitz deformation, ν_s converges strongly in $L^1(]-r_\alpha, r_\alpha[^d)$ to ν , unit outward normal field of Γ_α . So, using (17) and (16), (18) leads to:

$$\int_{]-r_\alpha, r_\alpha[^d} \int_{-L}^L \varphi(\hat{y})\phi'(\xi)a(\xi).\nu(\hat{y})g_\psi^\tau(\hat{y}, \xi) d\xi d\hat{y} = \int_{]-r_\alpha, r_\alpha[^d} H_\phi^\tau.\nu(\hat{y})\varphi(\hat{y}) d\hat{y}$$

for every test functions $(\phi, \varphi) \in \mathcal{D}(]-L, L]) \times \mathcal{D}(]-r_\alpha, r_\alpha[^d)$. The right-hand side of this equation is independent of ψ , sequence s^n and subsequence s^{n_p} so $a(\xi).\nu(\hat{y})g_\psi^\tau(\hat{y}, \xi)$ does not depend on those quantities too. Since (thanks to (12)):

$$\mathcal{L}(\{\xi \mid a(\xi).\nu(\hat{y}) = 0\}) = 0,$$

g_ψ^τ does not depend on those quantities too. We denote it f^τ and by uniqueness of the limit, $f_\psi(s, \cdot, \cdot)$ converges in $H^{-1}(]-r_\alpha, r_\alpha[^d \times]-L, L])$ to $f^\tau(\cdot, \cdot)$ for every regular Lipschitz deformation ψ . \square

The following description of strong traces, namely which are reached strongly by L^1 convergence, is a straightforward consequence of Lemma 1 with $N = d$, $\mathcal{O} =]-r_\alpha, r_\alpha[^d$, $f_n = f_\psi(s_n, \cdot, \cdot)$ and $f = f^\tau$ where $s_n \rightarrow 0$.

Proposition 2. *The function f^τ is a χ -function if and only if f^τ is a strong trace, namely for every ψ regular Lipschitz deformation:*

$$\operatorname{esslim}_{s \rightarrow 0} f_\psi(s, \cdot, \cdot) = f^\tau \text{ in } L^1(]-r_\alpha, r_\alpha[^d \times]-L, L[).$$

4. The localization method

In this section Ω_α is still fixed and we prove that $f^\tau(\hat{y}, \cdot)$ is a χ -function for almost every $\hat{y} \in]-r_\alpha, r_\alpha[^d$. We fix the particular Γ_α -Lipschitz deformation on Ω_α defined by: $\tilde{\psi}(s, \hat{y}) = (s + \gamma_\alpha(\hat{y}), \hat{y})$. Now we identify the variable $s = y_0$ and we denote

$$\tilde{f}(y, \xi) = f_\psi(y_0, \hat{y}, \xi) = f(\tilde{\psi}(y), \xi).$$

Notice that $\tilde{\psi}(y) \in \Omega_\alpha$ if and only if $y \in]0, r_\alpha[\times]-r_\alpha, r_\alpha[^d$. This will be convenient for the localization method. From (11) we find that \tilde{f} is solution of:

$$\tilde{a}^0(\hat{y}, \xi) \partial_{y_0} \tilde{f} + \hat{a}(\xi) \cdot \nabla_{\hat{y}} \tilde{f} = \partial_\xi \tilde{m}, \quad (19)$$

where

$$\tilde{a}^0(\hat{y}, \xi) = a^0(\xi) - \nabla \gamma_\alpha(\hat{y}) \cdot \hat{a}(\xi) \quad (20)$$

$$= \lambda(\hat{y}) a(\xi) \cdot \nu(\hat{y}) \quad (21)$$

with a $\lambda(\hat{y}) \neq 0$ and $\tilde{m}(y, \xi) = \tilde{m}(\tilde{\psi}(y), \xi)$.

Before introducing the notion of rescaled solution, let us state two lemmas.

Lemma 2. *There exists a sequence ϵ_n which converges to 0 and a set $\mathcal{E} \subset]-r_\alpha, r_\alpha[^d$ with $\mathcal{L}(]-r_\alpha, r_\alpha[^d \setminus \mathcal{E}) = 0$ such that for every $\hat{y} \in \mathcal{E}$ and every $R > 0$*

$$\lim_{n \rightarrow \infty} \frac{1}{\epsilon_n^d} \tilde{m}(]0, R\epsilon_n[\times (\hat{y} +) - R\epsilon_n, R\epsilon_n[^d \times]-L, L[) = 0.$$

Proof of Lemma 2: For every integer N , we denote

$$M_\epsilon^N(\hat{y}) = \frac{1}{\epsilon^d} \tilde{m}(]0, N\epsilon[\times (\hat{y} +) - N\epsilon, N\epsilon[^d \times]-L, L[).$$

Since M_ϵ^N is nonnegative, the L^1 norm of M_ϵ^N is :

$$\begin{aligned} & \int_{]-r_\alpha, r_\alpha[^d} M_\epsilon^N(\hat{y}) d\hat{y} \\ &= \int_{]-r_\alpha, r_\alpha[^d} \frac{1}{\epsilon^d} \int_0^{N\epsilon} \int_{]-N\epsilon, N\epsilon[^d} \int_{-L}^L \tilde{m}(z_0, \hat{y} + \hat{z}, \xi) d\xi d\hat{z} dz_0 d\hat{y} \\ &\leq \frac{1}{\epsilon^d} \int_{]-N\epsilon, N\epsilon[^d} \int_0^{N\epsilon} \int_{-L}^L \int_{]-r_\alpha - N\epsilon, r_\alpha + N\epsilon[^d} \tilde{m}(z_0, \hat{y}, \xi) d\hat{y} d\xi dz_0 d\hat{z}. \end{aligned}$$

We denote abusively $\tilde{m}(dz_0, d\hat{z}, d\xi) = \tilde{m}(z_0, \hat{z}, \xi) dz_0 d\hat{z} d\xi$ in this computation as if it was a function. This calculation is still correct since we just use the Fubini Theorem and a linear change of variable which are valid for measures. The last inequality can be written as:

$$\begin{aligned} & \int_{]-r_\alpha, r_\alpha[^d} M_\epsilon^N(\hat{y}) d\hat{y} \\ & \leq \frac{1}{\epsilon^d} \int_{]-N\epsilon, N\epsilon[^d} \tilde{m}([0, N\epsilon[\times] - r_\alpha - N\epsilon, r_\alpha + N\epsilon[^d \times] - L, L]) d\hat{z} \\ & \leq N^d \tilde{m}([0, N\epsilon[\times] - r_\alpha - N\epsilon, r_\alpha + N\epsilon[^d \times] - L, L]). \end{aligned}$$

By monotone convergence, since $\bigcap_{\epsilon > 0}]0, N\epsilon[= \emptyset$, this converges to 0 when ϵ converges to 0. Finally the L^1 norm of M_ϵ^N converges to 0 so there exists a subsequence ϵ_n and a set $\mathcal{E}_N \subset]-r_\alpha, r_\alpha[^d$ with $\mathcal{L}([-r_\alpha, r_\alpha[^d \setminus \mathcal{E}_N) = 0$ such that for every $\hat{y} \in \mathcal{E}_N$ $M_{\epsilon_n}^N(\hat{y})$ converges to 0 when ϵ_n goes to 0. By diagonal extraction, we can choose ϵ_n such that for every integer N and every $\hat{y} \in \mathcal{E}_N$, $M_{\epsilon_n}^N(\hat{y})$ converges to 0. This sequence ϵ_n with subset $\mathcal{E} = \bigcap_N \mathcal{E}_N$ verifies the required condition. \square

Lemma 3. *There exists a subsequence still denoted ϵ_n and a subset \mathcal{E}' of $]-r_\alpha, r_\alpha[^d$ such that $\mathcal{E}' \subset \mathcal{E}$, $\mathcal{L}([-r_\alpha, r_\alpha[^d \setminus \mathcal{E}') = 0$, and for every $\hat{y} \in \mathcal{E}'$ and every $R > 0$:*

$$\lim_{\epsilon_n \rightarrow 0} \int_{-L}^L \int_{]-R, R[^d} |f^\tau(\hat{y}, \xi) - f^\tau(\hat{y} + \epsilon_n \hat{y}, \xi)| d\underline{\hat{y}} d\xi = 0,$$

$$\lim_{\epsilon_n \rightarrow 0} \int_{-L}^L \int_{]-R, R[^d} |\tilde{a}^0(\hat{y}, \xi) - \tilde{a}^0(\hat{y} + \epsilon_n \hat{y}, \xi)| d\underline{\hat{y}} d\xi = 0.$$

Proof of Lemma 3: For every integer N we denote:

$$F_{\epsilon_n}^N(\hat{y}) = \int_{-L}^L \int_{]-N, N[^d} |f^\tau(\hat{y}, \xi) - f^\tau(\hat{y} + \epsilon_n \hat{y}, \xi)| d\underline{\hat{y}} d\xi.$$

Since $f^\tau \in L^\infty([-r_\alpha, r_\alpha[^d \times] - L, L])$, the L^1 norm of this function goes to zero as n tends to ∞ so there exists a subsequence still denoted ϵ_n and a subset $\mathcal{E}'_N \subset \mathcal{E}$ with $\mathcal{L}([-r_\alpha, r_\alpha[^d \setminus \mathcal{E}'_N) = 0$ such that for every $\hat{y} \in \mathcal{E}'_N$, $F_{\epsilon_n}^N(\hat{y})$ converges to 0 when n tends to infinity. By diagonal extraction we can find a subsequence such that this holds true for every N . Then this subsequence and $\mathcal{E}' = \bigcap_N \mathcal{E}'_N$ fulfill the required condition for the first limit. We consider in the same way the term with \tilde{a}^0 noticing that

$$\tilde{a}^0(\hat{y}, \xi) - \tilde{a}^0(\hat{y} + \epsilon_n \hat{y}, \xi) = \hat{a}(\xi) \cdot [\nabla \gamma_\alpha(\hat{y}) - \nabla \gamma_\alpha(\hat{y} + \epsilon_n \hat{y})]$$

with $\nabla \gamma_\alpha \in L^\infty([-r_\alpha, r_\alpha[^d)$. \square

We are now able to introduce the localization method. We denote

$$\Omega_\alpha^\epsilon =]0, r_\alpha/\epsilon[\times] - r_\alpha/\epsilon, r_\alpha/\epsilon[^d.$$

We want to show that for every $\hat{y} \in \mathcal{E}'$, $f^\tau(\hat{y}, \cdot)$ is a χ -function. From now on we fix such a $\hat{y} \in \mathcal{E}'$, and we denote $y = (0, \hat{y})$ the associated point on Γ_α . Then, we rescale the \tilde{f} function by introducing a new function \tilde{f}_ϵ which depends on a new variable $\underline{y} \in \Omega_\alpha^\epsilon$ and which is defined by:

$$\tilde{f}_\epsilon(\underline{y}, \xi) = \tilde{f}(y + \epsilon \underline{y}, \xi). \quad (22)$$

This function depends obviously on y but, since it is fixed all along this section, we skip it in the notation. Function \tilde{f}_ϵ is still a χ -function and we can notice that:

$$\tilde{f}_\epsilon(0, \underline{y}, \xi) = f^\tau(\hat{y} + \epsilon \underline{y}, \xi). \quad (23)$$

Hence we expect to gain some knowledge on $f^\tau(\hat{y}, \cdot)$ itself by studying the limit of \tilde{f}_ϵ when $\epsilon \rightarrow 0$. We define in the same way:

$$\tilde{a}_\epsilon^0(\underline{y}, \xi) = \tilde{a}^0(\hat{y} + \epsilon \underline{y}, \xi),$$

and we get from (19):

$$\tilde{a}_\epsilon^0(\underline{y}, \xi) \partial_{\underline{y}_0} \tilde{f}_\epsilon + \hat{a}(\xi) \cdot \nabla_{\underline{y}} \tilde{f}_\epsilon = \partial_\xi \tilde{m}_\epsilon, \quad (24)$$

where \tilde{m}_ϵ is the nonnegative measure defined for every real $R_1^i < R_2^i, L_1 < L_2$ by:

$$\tilde{m}_\epsilon \left(\prod_{0 \leq i \leq d} [R_1^i, R_2^i] \times [L_1, L_2] \right) = \frac{1}{\epsilon^d} \tilde{m} \left(\prod_{0 \leq i \leq d} [y_i + \epsilon R_1^i, y_i + \epsilon R_2^i] \times [L_1, L_2] \right). \quad (25)$$

We now pass to the limit when ϵ goes to 0 in the scaling.

Proposition 3. *There exist a sequence ϵ_n which goes to 0, and a χ -function $\tilde{f}_\infty \in L^\infty(\mathbb{R}^+ \times \mathbb{R}^d \times] - L, L[)$ such that \tilde{f}_{ϵ_n} converges strongly to \tilde{f}_∞ in $L^1_{\text{loc}}(\mathbb{R}^+ \times \mathbb{R}^d \times] - L, L[)$ and:*

$$\tilde{a}^0(\hat{y}, \xi) \partial_{\underline{y}_0} \tilde{f}_\infty + \hat{a}(\xi) \cdot \nabla_{\underline{y}} \tilde{f}_\infty = 0. \quad (26)$$

Notice that $\tilde{a}^0(\hat{y}, \xi)$ does not depend on the local variable \underline{y} . In fact, we have $\tilde{a}^0(\hat{y}, \xi) = \lambda(\hat{y}) a(\xi) \cdot \nu(\hat{y})$ where \hat{y} is the fixed point of the localization.

Proof of Proposition 3: We consider the sequence ϵ_n of Lemma 3. By weak compactness, there exists a function $\tilde{f}_\infty \in L^\infty(\mathbb{R}^+ \times \mathbb{R}^d \times] - L, L[)$ such that, up to extraction, \tilde{f}_{ϵ_n} converges weakly in L^∞ to \tilde{f}_∞ . Thanks to Lemma 2, \tilde{m}_{ϵ_n} converges to 0 in the sense of measure. Thanks to Lemma 3, $\tilde{a}_{\epsilon_n}^0$ converges strongly in $L^1_{\text{loc}}(\mathbb{R}^d \times] - L, L[)$ to $\tilde{a}^0(\hat{y}, \cdot)$, so passing to the limit in (24) gives (26). The strong convergence is an application of averaging lemmas. Here we use the following one which is a particular case of the version of Perthame and Souganidis (see [16]):

Theorem 4. *Let N be an integer, f_n be bounded in $L^\infty(\mathbb{R}^N)$ and $\{h_n^1, h_n^2\}$ be relatively compact in $[L^p(\mathbb{R}^N)]^{2N}$ with $1 < p < +\infty$ solutions of the transport equation:*

$$a(\xi) \cdot \nabla_y f_n = \partial_\xi(\nabla_y \cdot h_n^1) + \nabla_y \cdot h_n^2,$$

where $a \in [C^2(\mathbb{R})]^N$ verifies the non-degeneracy condition (12). Let $\phi \in \mathcal{D}(\mathbb{R})$, then the average $u_n^\phi(y) = \int_{\mathbb{R}} \phi(\xi) f_n(y, \xi) d\xi$ is relatively compact in $L^p(\mathbb{R}^N)$.

In order to do so we use the strategy of [15]. First we localize in \underline{y}, ξ . We introduce Φ_1, Φ_2 with values in $[0, 1]$ such that $\Phi_1 \in \mathcal{D}(\mathbb{R}^+ \times \mathbb{R}^d)$, $\Phi_2 \in \mathcal{D}(\mathbb{R})$, and $\text{Supp}(\Phi_1) \subset]1/(2R), 2R[\times]-2R, 2R[^d$, $\text{Supp}(\Phi_2) \subset]-2L, 2L[$. Moreover $\phi_1(\underline{y}) = 1$ for $\underline{y} \in]1/R, R[\times]-R, R[^d$ and $\Phi_2(\xi) = 1$ for $\xi \in]-L, L[$. Hence for $\epsilon < r_\alpha/(2R)$, we can define on $\mathbb{R} \times \mathbb{R}^d \times \mathbb{R}$:

$$\tilde{f}_\epsilon^R = \Phi_1 \Phi_2 \tilde{f}_\epsilon,$$

(where $\tilde{f}_\epsilon^R = 0$ if \tilde{f}_ϵ is not defined). On $]1/R, R[\times]-R, R[^d \times]-L, L[$ we have $\tilde{f}_\epsilon^R = \tilde{f}_\epsilon$. So, if we denote $a_{\hat{y}}(\xi) = (\tilde{a}^0(\hat{y}, \xi), \hat{a}(\xi))$ (which depends only on ξ since \hat{y} is fixed), from (24) we get:

$$a_{\hat{y}}(\xi) \cdot \nabla_{\underline{y}} \tilde{f}_\epsilon^R = \partial_\xi(\Phi_1 \Phi_2 \tilde{m}_\epsilon) - \Phi_1 \Phi_2' \tilde{m}_\epsilon + a_{\hat{y}}(\xi) \cdot \nabla_{\underline{y}} \phi_1 \phi_2 \tilde{f}_\epsilon \quad (27)$$

$$+ \partial y_0 [(\tilde{a}^0(\hat{y}, \xi) - \tilde{a}_\epsilon^0(\hat{y}, \xi)) \tilde{f}_\epsilon^R] \quad (28)$$

$$= \partial_\xi \mu_{1,\epsilon} + \mu_{2,\epsilon} + \partial y_0 [(\tilde{a}^0(\hat{y}, \xi) - \tilde{a}_\epsilon^0(\hat{y}, \xi)) \tilde{f}_\epsilon^R], \quad (29)$$

where μ_{1,ϵ_n} and μ_{2,ϵ_n} are measures uniformly bounded with respect to n . Moreover thanks to Lemma 3 $a_{\hat{y}}^0 - \tilde{a}_{\epsilon_n}^0(\hat{y}, \xi)$ converges to 0 in $L^1_{\text{loc}}(\mathbb{R}^d \times]-L, L[)$. So it converges to 0 in L^p_{loc} for every $1 \leq p < \infty$ since this functions are bounded in L^∞ . Since the measures are compactly imbedding in $W^{-1,p}$ for $1 \leq p < (d+2)/(d+1)$, we can apply Theorem 4 with $N = d+1$, $f_n = \tilde{f}_{\epsilon_n}^R$, $\phi(\xi) = \phi_2(\xi)$ and $a(\xi) = a_{\hat{y}}(\xi)$. It follows that $\int \tilde{f}_{\epsilon_n}^R \phi_2(\xi) d\xi$ is compact in L^p for $1 \leq p < (d+2)/(d+1)$. And so by uniqueness of the limit, $\int \tilde{f}_{\epsilon_n}(\cdot, \xi) d\xi$ converges strongly to $\int \tilde{f}_\infty(\cdot, \xi) d\xi$ in $L^1_{\text{loc}}(\mathbb{R}^{d+1})$. Lemma 1 ensures us that \tilde{f}_{ϵ_n} converges strongly to \tilde{f}_∞ in $L^1_{\text{loc}}(\mathbb{R}^{d+1} \times]-L, L[)$ and that \tilde{f}_∞ is a χ -function. \square

We now turn to characterize the limit function \tilde{f}_∞ . First notice that from (26) $\tilde{a}^0(\hat{y}, \xi) \tilde{f}_\infty$ lies in $C^0(\mathbb{R}^+, W^{-1,\infty}(\mathbb{R}^d \times]-L, L[))$. From (12) $\tilde{a}^0(\hat{y}, \xi) \neq 0$ for almost every ξ so we can define $\tilde{f}_\infty(0, \cdot, \cdot)$ from it.

Proposition 4. *For every $\hat{y} \in \mathcal{E}'$ we have the following equality:*

$$\tilde{f}_\infty(0, \hat{y}, \xi) = f^r(\hat{y}, \xi)$$

for almost every $(\hat{y}, \xi) \in \mathbb{R}^d \times]-L, L[$.

Notice that this implies that $\tilde{f}_\infty(0, \cdot, \cdot)$ does not depend on \underline{y} .

Proof of Proposition 4: Let us introduce

$$h_\Phi^\epsilon(\underline{y}_0) = \int_{-L}^L \int_{\mathbb{R}^d} [\tilde{a}_\epsilon^0(\underline{y}, \xi) \tilde{f}_\epsilon(\underline{y}_0, \underline{y}, \xi) - \tilde{a}^0(\underline{y}, \xi) \tilde{f}_\infty(\underline{y}_0, \underline{y}, \xi)] \Phi(\underline{y}, \xi) d\underline{y} d\xi$$

for every test function $\Phi \in C_0^\infty(\mathbb{R}^d \times]-L, L[)$. We have:

$$|h_\Phi^\epsilon(0)| \leq C(\|h_\Phi^\epsilon\|_{L^1(]0,1[)} + \|\partial_{\underline{y}_0} h_\Phi^\epsilon\|_{\mathcal{M}(]0,1[)}),$$

and from (24) and (26) h_Φ^ϵ is a BV function and:

$$\begin{aligned} \|\partial_{\underline{y}_0} h_\Phi^\epsilon\|_{\mathcal{M}} &\leq C_\Phi(\|\tilde{f}_\epsilon - \tilde{f}_\infty\|_{L^1_{\text{loc}}} + \|\tilde{m}_\epsilon\|_{\mathcal{M}}) \\ \|h_\Phi^\epsilon\|_{L^1(]0,1[)} &\leq C_\Phi(\|\tilde{f}_\epsilon - \tilde{f}_\infty\|_{L^1_{\text{loc}}} + \|\tilde{a}_\epsilon^0 - \tilde{a}^0\|_{L^1_{\text{loc}}}). \end{aligned}$$

Lemma 2 and definition (25) ensure that $\|\tilde{m}_{\epsilon_n}\|_{\mathcal{M}}$ converges to 0, Lemma 3 that $\|\tilde{a}_{\epsilon_n}^0 - \tilde{a}^0\|_{L^1_{\text{loc}}}$ converges to 0 and Proposition 3 that $\|\tilde{f}_{\epsilon_n} - \tilde{f}_\infty\|_{L^1_{\text{loc}}}$ converges to 0. So $h_\Phi^{\epsilon_n}(0)$ converges to 0 when n converges to $+\infty$. Remembering (23), thanks to Lemma 3, $\tilde{f}_{\epsilon_n}(0, \cdot, \cdot)$ converges strongly to f^τ , and so, $\tilde{a}_{\epsilon_n}^0(\cdot, \cdot) \tilde{f}_{\epsilon_n}(0, \cdot, \cdot)$ converges strongly to $\tilde{a}^0(\underline{y}, \cdot) f^\tau$. Then

$$\tilde{a}^0(\underline{y}, \xi) \tilde{f}_\infty(0, \underline{y}, \xi) = \tilde{a}^0(\underline{y}, \xi) f^\tau(\underline{y}, \xi)$$

for almost every $(\underline{y}, \xi) \in \mathbb{R}^d \times]-L, L[$ which leads to the desired result noticing that $\tilde{a}^0(\underline{y}, \xi) \neq 0$ for almost every ξ . \square

From (26) we deduce that:

$$\tilde{f}_\infty(\tilde{a}^0(\underline{y}, \xi) \underline{y}_0, \underline{y} + \hat{a}(\xi) \underline{y}_0, \xi) = f^\tau(\underline{y}, \xi),$$

for almost every $\underline{y}_0 > 0$. But $\tilde{a}^0(\underline{y}, \xi) \neq 0$ for almost every ξ so

$$\tilde{f}_\infty(\underline{y}, \xi) = f^\tau(\underline{y}, \xi)$$

for almost every $(\underline{y}, \xi) \in \mathbb{R}^{d+1} \times]-L, L[$, which is constant with respect to \underline{y} . Finally since \tilde{f}_∞ is a χ -function for almost every \underline{y} , we deduce that:

Proposition 5. *For every $\hat{y} \in \mathcal{E}'$ function $f^\tau(\hat{y}, \cdot)$ is a χ -function.*

5. Proof of Theorem 1

For every α and every ψ $\partial\Omega$ -regular Lipschitz deformation, the restriction of ψ to Γ_α is a Γ_α -regular Lipschitz deformation. From Proposition 5 and Proposition 2 it follows:

$$\operatorname{esslim}_{y_0 \rightarrow 0} \int_{-r_\alpha}^{r_\alpha} \int_{-L}^L |f_\psi(y_0, \hat{y}, \xi) - f^\tau(\hat{y}, \xi)| d\xi d\hat{y} = 0. \quad (30)$$

We define $u^\tau \in L^\infty(\partial\Omega)$ by:

$$u^\tau(\hat{z}) = \int_{-L}^L f^\tau(\hat{y}, \xi) d\xi$$

if $(\gamma_\alpha(\hat{y}), \hat{y}) = \mathcal{R}_\alpha(\hat{z})$. Notice that it is uniquely well defined thanks to the uniqueness of f^τ in Proposition 1. Finally, since $\|\nabla\gamma_\alpha\|_{L^\infty([-r_\alpha, r_\alpha])} \leq C$, from (30) and remembering that $u(\psi(s, \hat{z})) = \int_{-L}^L f_\psi(s, \hat{y}, \xi) d\xi$ if $\mathcal{R}_\alpha(\hat{z}) = (\gamma_\alpha(\hat{y}), \hat{y})$, we deduce:

Proposition 6. *For every α and every ψ $\partial\Omega$ -regular Lipschitz deformation,*

$$\operatorname{esslim}_{s \rightarrow 0} \int_{\Gamma_\alpha} |u(\psi(s, \hat{z})) - u^\tau(\hat{z})| d\mathcal{H}^d(\hat{z}) = 0.$$

Finally for every compact set $K \subset\subset \partial\Omega$, $\{\Gamma_\alpha\}$ is a covering of K with open sets of $\partial\Omega$ so there exists I_0 finite set such that $K \subset \bigcup_{\alpha \in I_0} \Gamma_\alpha$ and so

$$\int_K |u(\psi(s, \hat{z})) - u^\tau(\hat{z})| d\mathcal{H}^d(\hat{z}) \leq \sum_{\alpha \in I_0} \int_{\Gamma_\alpha} |u(\psi(s, \hat{z})) - u^\tau(\hat{z})| d\mathcal{H}^d(\hat{z}).$$

Then Theorem 1 follows from Proposition 6. \square

Acknowledgements. I am grateful to P.Marcati who brought this problem to my knowledge.

References

1. V. I. Agoshkov. Spaces of functions with differential-difference characteristics and the smoothness of solutions of the transport equation. *Dokl. Akad. Nauk SSSR*, 276(6):1289–1293, 1984.
2. C. Bardos, A. Y. le Roux, and J.-C. Nédélec. First order quasilinear equations with boundary conditions. *Comm. Partial Differential Equations*, 4(9):1017–1034, 1979.
3. Gui-Qiang Chen and Hermano Frid. Divergence-measure fields and hyperbolic conservation laws. *Arch. Ration. Mech. Anal.*, 147(2):89–118, 1999.
4. Gui-Qiang Chen and Michel Rascle. Initial layers and uniqueness of weak entropy solutions to hyperbolic conservation laws. *Arch. Ration. Mech. Anal.*, to appear.

5. C. Cheverry. Effet régularisant pour une loi de conservation scalaire multidimensionnelle. In *Seminaire: Équations aux Dérivées Partielles, 1998-1999*, pages Exp. No. XXIV, 15. École Polytech., Palaiseau, 1999.
6. R. J. DiPerna. Convergence of approximate solutions to conservation laws. *Arch. Rational Mech. Anal.*, 82(1):27-70, 1983.
7. R. J. DiPerna, P.-L. Lions, and Y. Meyer. L^p regularity of velocity averages. *Ann. Inst. H. Poincaré Anal. Non Linéaire*, 8(3-4):271-287, 1991.
8. François Dubois and Philippe Le Floch. Boundary conditions for nonlinear hyperbolic systems of conservation laws. *J. Differential Equations*, 71(1):93-122, 1988.
9. François Golse, Benoît Perthame, and Rémi Sentis. Un résultat de compacité pour les équations de transport et application au calcul de la limite de la valeur propre principale d'un opérateur de transport. *C. R. Acad. Sci. Paris Sér. I Math.*, 301(7):341-344, 1985.
10. S. N. Kružkov. First order quasilinear equations with several independent variables. *Mat. Sb. (N.S.)*, 81 (123):228-255, 1970.
11. P.-L. Lions, B. Perthame, and E. Tadmor. A kinetic formulation of multidimensional scalar conservation laws and related equations. *J. Amer. Math. Soc.*, 7(1):169-191, 1994.
12. A. Nouri, A. Omrane, and J. P. Vila. Boundary conditions for scalar conservation laws from a kinetic point of view. *J. Statist. Phys.*, 94(5-6):779-804, 1999.
13. Felix Otto. Initial-boundary value problem for a scalar conservation law. *C. R. Acad. Sci. Paris Sér. I Math.*, 322(8):729-734, 1996.
14. Felix Otto. A regularizing effect of nonlinear transport equations. *Quart. Appl. Math.*, 56(2):355-375, 1998.
15. B. Perthame. Lecture notes on kinetic formulation of conservation laws. *Studies in Advanced Mathematics*, 15:111-139, 2000.
16. B. Perthame and P. E. Souganidis. A limiting case for velocity averaging. *Ann. Sci. École Norm. Sup. (4)*, 31(4):591-598, 1998.
17. Anders Szepessy. Measure-valued solutions of scalar conservation laws with boundary conditions. *Arch. Rational Mech. Anal.*, 107(2):181-193, 1989.
18. A. Vasseur. Time regularity for the system of isentropic gas dynamics with $\gamma = 3$. *Comm. Partial Differential Equations*, 24(11-12):1987-1997, 1999.

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